# Graduate mini couse: "High Energy Astrophysics – Selected Topics"

## Assignment #1

## Due 1pm, Thu Mar 6

This assignment is due before class on Mar 5; that is, at 1:00pm **sharp**. Assignments may be handed in in person at class; may be placed at our offices 1303 (Huirong) or 1304 (Christoph). The instructors will have office hours 2-3pm on Wed.

Show your work, and good luck!

#### Question 1 - Generalized Force Term due to Turbulence [10 pts]

(10 pts) As we discussed in class, the appropriate form of the *relativistic Vlasov equation* reads as follows

$$\frac{\partial f}{\partial t} + \upsilon \mu \frac{\partial f}{\partial Z} - \varepsilon \Omega \frac{\partial f}{\partial \phi} + \frac{1}{p^2} \frac{\partial}{\partial x_\sigma} (p^2 g_{x_\sigma} f) = S(\mathbf{x}, \mathbf{p}, t), \tag{1}$$

where f = f(x, p, t) represents the phase space density of plasma particles of a given sort,  $p = \gamma m v$  is the momentum,  $x_{\sigma}$  represents the coordinate set  $x_{\sigma} = (p, \mu, \phi, X, Y, Z)$  with spherical coordinates in momentum space and the coordinates of the guiding center for gyration  $\mathbf{R} = (X, Y, Z)$ , and Z denotes the spatial variable along the locally uniform magnetic field  $\mathbf{B}_0 = B_0 \mathbf{e}_z$  implying  $\mu = p_{\parallel}/p$ . The generalized force term  $g_{\sigma}$  includes the effects of the randomly fluctuating electromagnetic fields and S denotes the sources and sinks of particles.

Compute the generalized force terms  $g_p$  and  $g_\mu$  solely as functions of our new coordinate system  $x_\sigma$  by introducing the left-handed and right-handed polarized field components,

$$\delta B_{\mathrm{L,R}} \equiv \frac{1}{\sqrt{2}} (\delta B_x \pm i \, \delta B_y), \quad \delta B_{\parallel} \equiv \delta B_z,$$
$$\delta E_{\mathrm{L,R}} \equiv \frac{1}{\sqrt{2}} (\delta E_x \pm i \, \delta E_y), \quad \delta E_{\parallel} \equiv \delta E_z.$$

#### Question 2 - Quasilinear Theory of Cosmic Rays [20 pts]

2 (a)

(15 pts) In class, we have shown that under certain conditions the relativistic Vlasov equation corresponds to the Fokker-Planck equation that reads for the gyrotropic phase space density  $f(Z, p, \mu, t)$  as follows:

$$\frac{\partial f}{\partial t} + \upsilon \mu \frac{\partial f}{\partial Z} - S_0(p,\mu,t) = \frac{\partial}{\partial \mu} \left[ D_{\mu\mu} \frac{\partial f}{\partial \mu} + D_{\mu p} \frac{\partial f}{\partial p} \right] + \frac{1}{p^2} \frac{\partial}{\partial p} p^2 \left[ D_{\mu p} \frac{\partial f}{\partial \mu} + D_{pp} \frac{\partial f}{\partial p} \right],\tag{2}$$

where we only consider isotropic source terms  $S_0(p, \mu, t)$  and the Fokker-Planck coefficients  $D_{\mu\mu}, D_{\mu p}, D_{pp}$  have been defined in class. Your task is to derive the diffusion convection equation for cosmic ray transport,

$$\frac{\partial F}{\partial t} - S_0(Z, p, t) = \frac{\partial}{\partial Z} \left[ \kappa(Z, p, t) \frac{\partial F}{\partial Z} \right] + \frac{1}{p^2} \frac{\partial}{\partial p} \left[ p^2 \Gamma(Z, p, t) \frac{\partial F}{\partial p} \right] + \frac{v}{4} \frac{\partial A_1}{\partial Z} \frac{\partial F}{\partial p} - \frac{1}{4p^2} \frac{\partial (p^2 v A_1)}{\partial p} \frac{\partial F}{\partial Z}, (3)$$

where the spatial diffusion coefficient  $\kappa$ , the momentum diffusion coefficient  $\Gamma$ , and  $A_1$  are determined by pitch angle averages of three Fokker-Planck coefficients

$$\kappa = \frac{v^2}{8} \int_{-1}^{1} d\mu \frac{(1-\mu^2)^2}{D_{\mu\mu}(\mu)},$$
  

$$\Gamma = \frac{1}{2} \int_{-1}^{1} d\mu \left[ D_{pp}(\mu) - \frac{D_{\mu p}^2(\mu)}{D_{\mu\mu}(\mu)} \right],$$
  

$$A_1 = \int_{-1}^{1} d\mu \left(1-\mu^2\right) \frac{D_{\mu p}(\mu)}{D_{\mu\mu}(\mu)}.$$
(4)

We derived this diffusion-convection equation in the frame that is comoving with the plasma. By applying a Galilean transformation to Eqn. (3), show that it reduces to the well-known form in the laboratory system:

$$\frac{\partial F}{\partial t} - S_0(Z, p^*, t) = \frac{\partial}{\partial Z} \left[ \kappa(Z, p^*, t) \frac{\partial F}{\partial Z} \right] - V \frac{\partial F}{\partial Z} + \frac{1}{p^{*2}} \frac{\partial}{\partial p^{*2}} \left[ p^{*2} \Gamma(Z, p^*, t) \frac{\partial F}{\partial p^*} \right] + \frac{p^*}{3} \frac{\partial V}{\partial Z} \frac{\partial F}{\partial p^*}, \quad (5)$$

where we introduced the plasma bulk speed *U*, and the cosmic ray bulk speed *V*,

$$V = v^* + \frac{1}{4p^{*2}} \frac{\partial (p^{*2}v^*A_1)}{\partial p^*}, \quad A_1 = \int_{-1}^1 d\mu \left(1 - \mu^2\right) \frac{D_{\mu p}(\mu)}{D_{\mu \mu}(\mu)}.$$
 (6)

Here, we use a mixed coordinate system in which the space coordinates (*X*) are measured in the lab system and the momentum coordinates ( $p^*$ ,  $\mu^*$ ) in the rest frame of the streaming plasma.

*Procedure:* Split the total density f into an isotropic part F and an anisotropic part  $\delta f$ ,

$$f(Z, p, \mu, t) = F(z, p, t) + \delta f(Z, p, \mu, t) \quad \text{where} \quad F(z, p, t) \equiv \frac{1}{2} \int_{-1}^{1} d\mu f(Z, p, \mu, t), \tag{7}$$

and use similar ideas as in class to combine the pitch-angle averaged equation with the original one. Use then the diffusive approximation which applies if the isotropic particle density is slowly evolving, i.e.

$$\frac{\partial F}{\partial t} = \mathcal{O}\left(\frac{F}{T}\right), \quad \frac{\partial F}{\partial z} = \mathcal{O}\left(\frac{F}{L}\right)$$

with typical length scales  $L \gg \lambda$  and time scales  $T \gg \tau$  much larger than the mean free path  $\lambda = v\tau$  and the pitch angle scattering relaxation time  $\tau \simeq \mathcal{O}(1/D_{\mu\mu})$ . Under these conditions, the particles can reach locally near-equilibrium which results in a small anisotropy, i.e.  $\delta f \ll F$ . If you regard  $\delta f$  of order  $\tau$  and F of order 1, and recall that  $D_{\mu p}$  and  $D_{pp}$  are of order  $\varepsilon/\tau$  and  $\varepsilon^2/\tau$ , respectively, where  $\varepsilon = v_{\rm ph}/v = v_A/v \ll 1$ , smaller than  $D_{\mu\mu} = \mathcal{O}(1/\tau)$ , you may characterize the differential operators by different timescales. To lowest order, the approximate equation should look as follows:

$$\frac{\partial}{\partial \mu} \left[ D_{\mu\mu} \frac{\partial \delta f}{\partial \mu} + D_{\mu p} \frac{\partial F}{\partial p} \right] \simeq v \mu \frac{\partial F}{\partial Z}.$$
(8)

Solve this equation for  $\delta f$ , determine the integration constants appropriately, and insert the solution into the pitch-angle averaged equation enables you to derive Eqns. (3) and (4). To perform the Galilean transformation, use a mixed coordinate system in which the space coordinates (X') are measured in the lab system and the momentum coordinates ( $p^*$ ,  $\mu^*$ ) in the rest frame of the streaming plasma. Justify and use the following system of equations for your transformation:

$$t = t', \quad X = X', \quad Y = Y', \quad Z = Z' + Ut'$$
 (9)

$$p^* = p^* \left( 1 - \frac{t'}{3} \frac{\partial U}{\partial Z'} \right), \quad \mu^* = \mu.$$
(10)

Finally, in your derivation, you might want to use the general scaling of the rate of adiabatic deceleration,

$$A_1(Z,p) = A_1^0(Z)\frac{p}{v}.$$
(11)

### 2 (b)

(5 pts) Explain in your own words the relevant steps and approximations that were necessary in deriving the diffusion convection equation for cosmic ray transport (5) from the Vlasov equation (1).