

# Constraining the Ultra-Large Scale Structure of the Universe

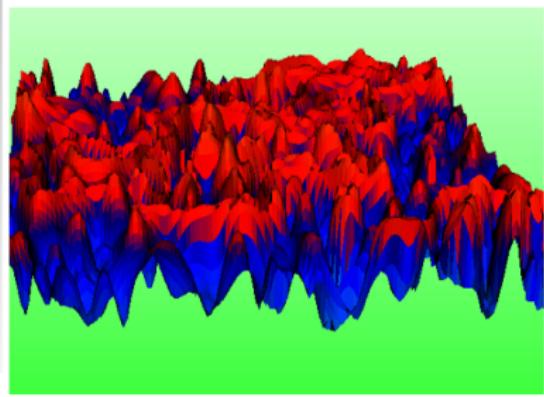
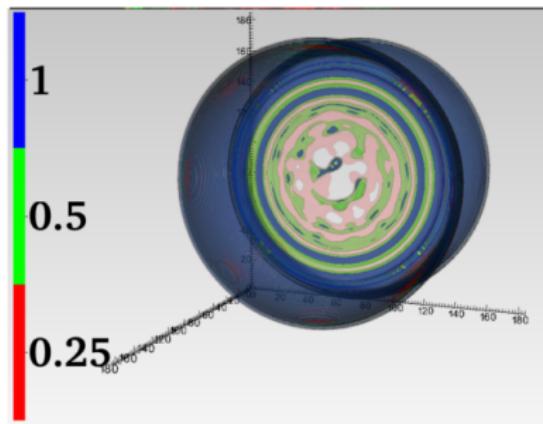
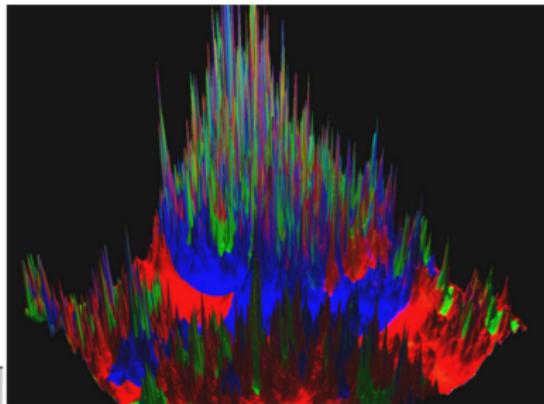
Jonathan Braden

University College London

ICG Portsmouth, May 25, 2016

w/ Hiranya Peiris, Matthew Johnson and Anthony Aguirre  
based on arXiv:1604.04001 and *in progress*

# Inhomogeneous Nonlinear Ultra-Large Scale Cosmology



## ... A Mosaic of Interesting Dynamics

- ▶ **Strongly inhomogeneous and nonlinear cosmological initial conditions** (this talk) [JB, Peiris, Johnson, Aguirre]
- ▶ Isocurvature mode conversion into intermittent density perturbations [JB, Bond, Frolov, Huang]
  - ▶ Caustic formation in chaotic long wavelength dynamics
  - ▶ Generalised form of local nonGaussianity
- ▶ First order phase transitions [JB, Bond, Mersini-Houghton]
- ▶ Entropy production in highly inhomogeneous nonlinear field theories [JB, Bond]

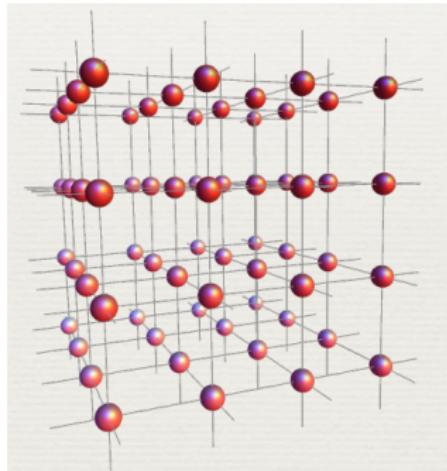
Spatially intermittent dynamics, fundamental issues in QFT, novel and poorly constrained observables

## Hybrid MPI/OpenMP Lattice Code

- ▶ Solve field equation (e.g.)

$$\ddot{\phi}_i + 3\frac{\dot{a}}{a}\dot{\phi}_i - \frac{\nabla^2 \phi_i}{a^2} + V'(\vec{\phi}) = 0$$

- ▶ 10th order Gauss-Legendre integration (general) or 8th order Yoshida (nonlinear sigma models)
- ▶ Finite-difference (fully parallel) or Pseudospectral (OpenMP)
- ▶ Optional absorbing boundaries
- ▶ Quantum fluctuations → realization of random field



- ▶ Energy conservation  $\mathcal{O}(10^{-9} - 10^{-14})$

## Hybrid MPI/OpenMP Lattice Code

- ▶ Solve field equation (e.g.)

$$\ddot{\phi}_i + 3\frac{\dot{a}}{a}\dot{\phi}_i - \frac{\nabla^2 \phi_i}{a^2} + V'(\vec{\phi}) = 0$$

- ▶ 10th order Gauss-Legendre integration (general) or 8th order Yoshida (nonlinear sigma models)
- ▶ Finite-difference (fully parallel) or Pseudospectral (OpenMP)
- ▶ Optional absorbing boundaries
- ▶ Quantum fluctuations → realization of random field



- ▶ Energy conservation  $\mathcal{O}(10^{-9} - 10^{-14})$

# Dawn of Numerical Relativity in Cosmology

- ▶ Bentivegna, **Bruni**
- ▶ Mertens, Giblin, Starkman
- ▶ Kleban, Linde, Senatore, West
- ▶ Peiris, Johnson, Feeney, Aguirre, Wainwright (symmetry reduced bubbles)
- ▶ Adamek, Daverio, Durrer, Kunz (weak field subhorizon)

# Dawn of Numerical Relativity in Cosmology

- ▶ Bentivegna, **Bruni**
- ▶ Mertens, Giblin, Starkman
- ▶ Kleban, Linde, Senatore, West
- ▶ Peiris, Johnson, Feeney, Aguirre, Wainwright (symmetry reduced bubbles)
- ▶ Adamek, Daverio, Durrer, Kunz (weak field subhorizon)

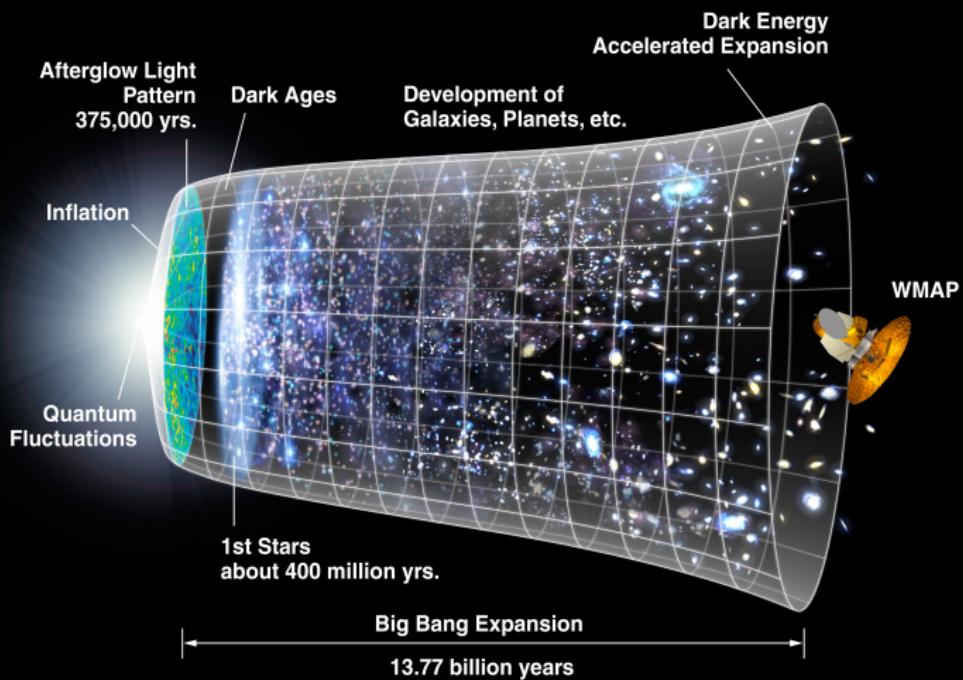
## Connection of Initial Conditions with Observables

- ▶ Requires Monte Carlo sampling of Initial Conditions
- ▶ Highly accurate integrators to achieve machine precision

# Outline

- ▶ Brief Review of Cosmology and Role of Inflation
- ▶ Inflationary Models and Initial Conditions
- ▶ Choice of Initial Conditions
- ▶ Planar Symmetric Dynamics
- ▶ Projection of Dynamics into Observations (CMB quadrupole)
- ▶ Semi-analytic Approximation (work in progress)

# History of the Universe

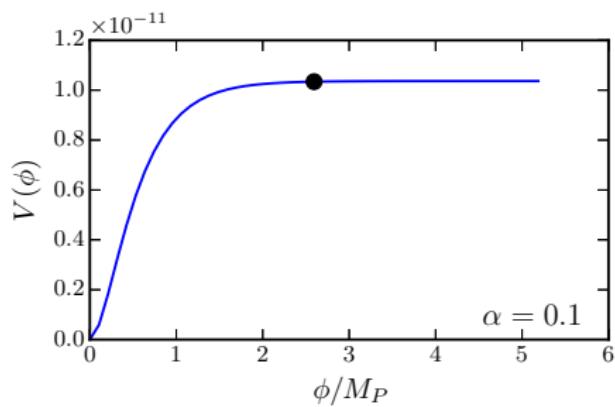


# Model Choices

$$V(\phi) = V_0 \left(1 - e^{-\sqrt{\frac{2}{3\alpha}} \frac{\phi}{M_P}}\right)^2$$

# Model Choices

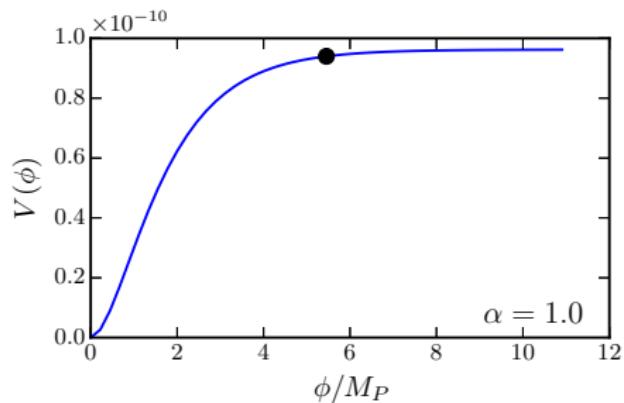
$$V(\phi) = V_0 \left( 1 - e^{-\sqrt{\frac{2}{3\alpha}} \frac{\phi}{M_P}} \right)^2$$



- $\alpha \ll 1 \rightarrow$  small-field model

## Model Choices

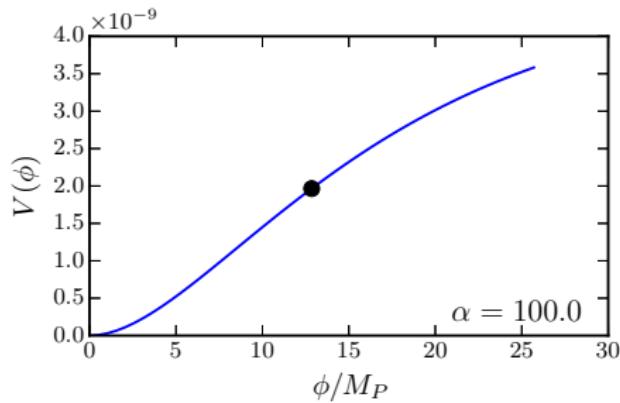
$$V(\phi) = V_0 \left( 1 - e^{-\sqrt{\frac{2}{3\alpha}} \frac{\phi}{M_P}} \right)^2$$



- ▶  $\alpha \ll 1 \rightarrow$  small-field model
- ▶  $\alpha = 1 \rightarrow$  Starobinski

## Model Choices

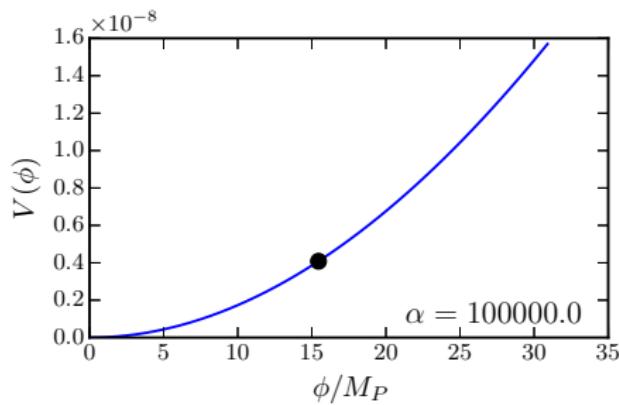
$$V(\phi) = V_0 \left( 1 - e^{-\sqrt{\frac{2}{3\alpha}} \frac{\phi}{M_P}} \right)^2$$



- ▶  $\alpha \ll 1 \rightarrow$  small-field model
- ▶  $\alpha = 1 \rightarrow$  Starobinski
- ▶  $\alpha \gg 1 \rightarrow m^2 \phi^2 / 2$

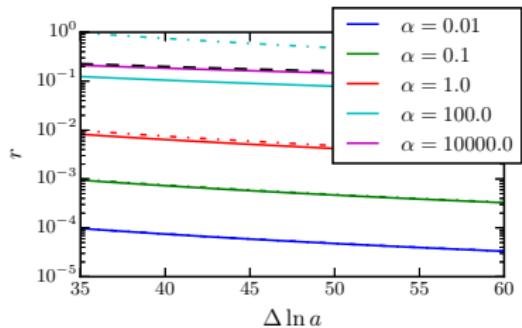
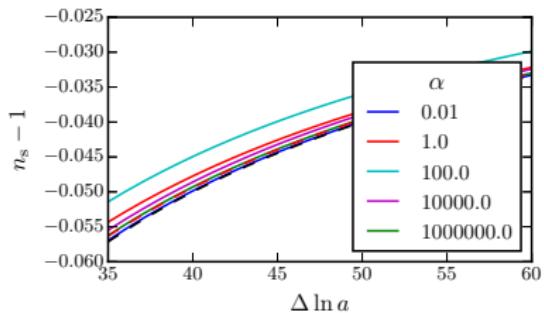
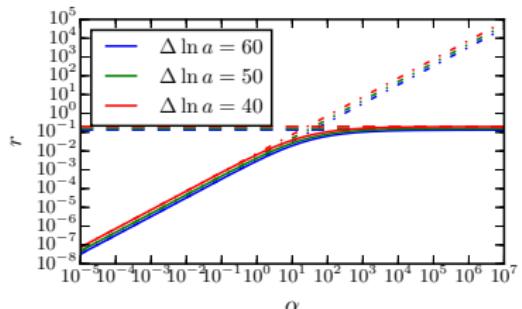
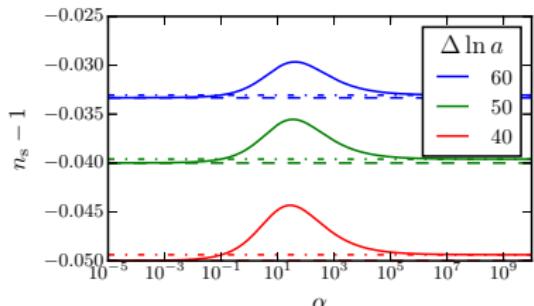
# Model Choices

$$V(\phi) = V_0 \left( 1 - e^{-\sqrt{\frac{2}{3\alpha}} \frac{\phi}{M_P}} \right)^2$$

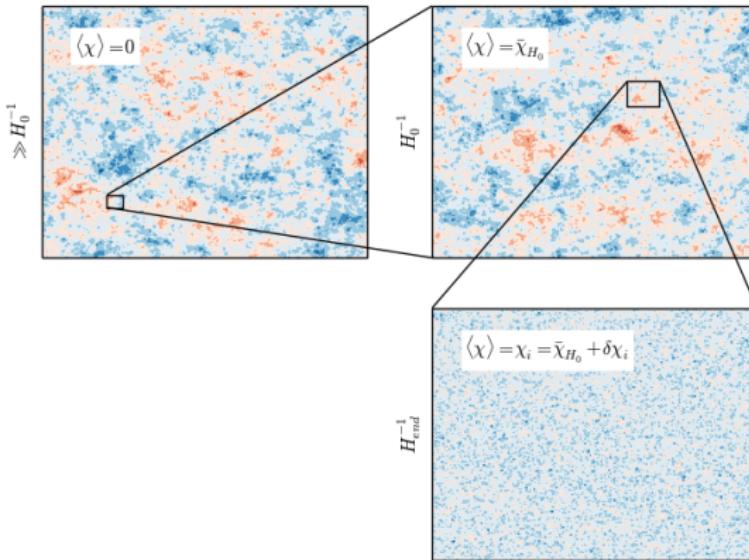


- ▶  $\alpha \ll 1 \rightarrow$  small-field model
- ▶  $\alpha = 1 \rightarrow$  Starobinski
- ▶  $\alpha \gg 1 \rightarrow m^2 \phi^2 / 2$

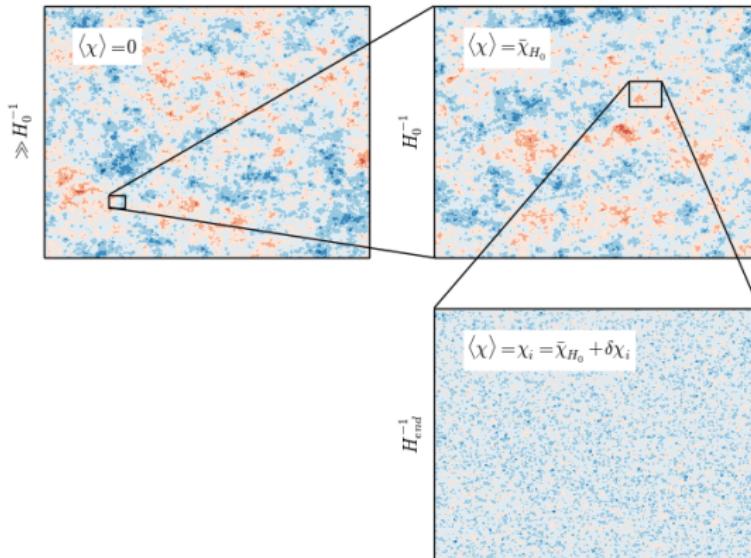
# CMB Parameter Predictions



# Ultra-Large Scale Structure

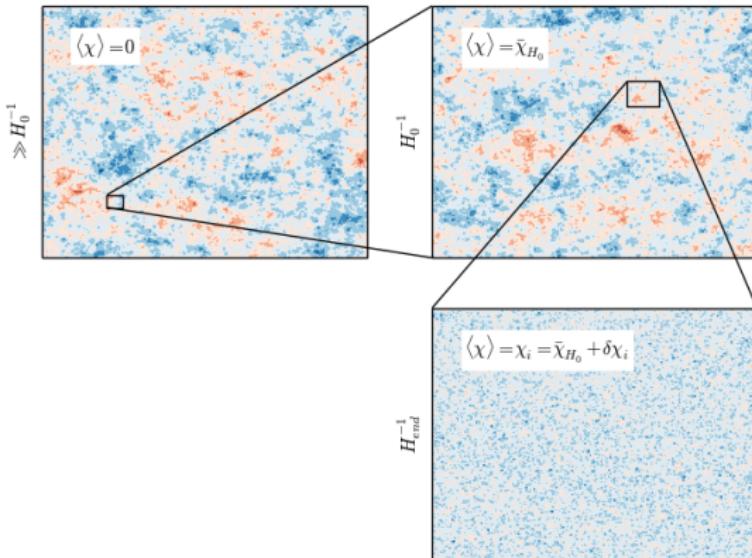


# Ultra-Large Scale Structure



Local Remnants of Ultra-Large Scale Structure?

# Ultra-Large Scale Structure



## Local Remnants of Ultra-Large Scale Structure?

- ▶ Structure present at start of inflation
- ▶ Conversion of structure during or after inflation

# Dimensional Reduction: Planar Symmetry

## Synchronous Gauge

$$ds^2 = -d\tau^2 + a_{\parallel}^2(x, \tau)dx^2 + a_{\perp}^2(x, \tau)(dy^2 + dz^2)$$

Residual gauge freedom :  $a_{\parallel}(x, \tau = 0), a_{\perp}(x, \tau = 0)$

## Isotropised Expansion

$$a \equiv \det(\gamma_{ij})^{1/6} = (a_{\parallel} a_{\perp}^2)^{1/3}$$

$$H \equiv -\frac{1}{3}\gamma^{ij}K_{ij} = -\frac{1}{3}(K^x{}_x + 2K^y{}_y)$$

## Initial Conditions: Field

Work in Spatially Flat Gauge  $a_{\parallel}(\tau = 0) = 1 = a_{\perp}(\tau = 0)$

$$\phi(x) = \bar{\phi} + \delta\hat{\phi}$$

$\bar{\phi}$  gives desired  $\mathcal{N}$  e-folds of inflation in homogeneous limit

$$3H_I^2 \equiv V(\bar{\phi})$$

## Field Fluctuations

$$\delta\hat{\phi}(x_i) = A_{\phi} \sum_{n=1} \hat{G} e^{ik_n x_i} \sqrt{P(k_n)} \quad \hat{G} = \sqrt{-2 \ln \hat{\beta}} e^{2\pi i \hat{\alpha}}$$

## Initial Conditions: Field

Work in Spatially Flat Gauge  $a_{\parallel}(\tau = 0) = 1 = a_{\perp}(\tau = 0)$

$$\phi(x) = \bar{\phi} + \delta\hat{\phi}$$

$\bar{\phi}$  gives desired  $\mathcal{N}$  e-folds of inflation in homogeneous limit

$$3H_I^2 \equiv V(\bar{\phi})$$

## Field Fluctuations

$$\delta\hat{\phi}(x_i) = A_{\phi} \sum_{n=1} \hat{G} e^{ik_n x_i} \sqrt{P(k_n)} \quad \hat{G} = \sqrt{-2 \ln \beta} e^{2\pi i \hat{a}}$$

$$P(k) = \Theta(k_{\max} - k) \quad H_I^{-1} k_{\max} = 2\pi\sqrt{3}$$

## Initial Conditions: Metric

Work in Spatially Flat Gauge  $a_{\parallel}(\tau = 0) = 1 = a_{\perp}(\tau = 0)$

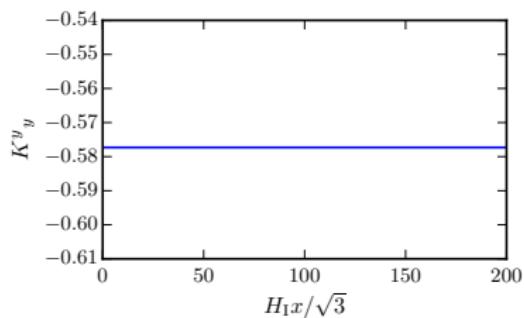
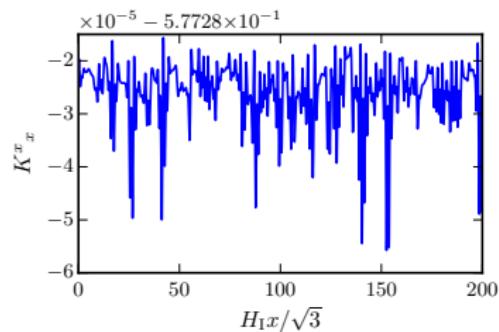
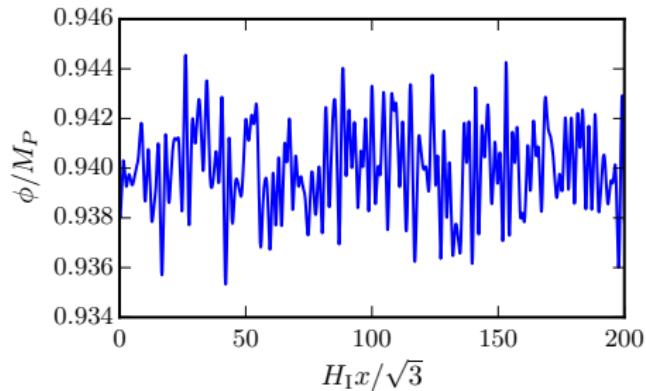
Momentum Constraint Sews Neighbouring Grid Sites Together

$$0 = \mathcal{P} = K^y_y' - \frac{a_{\perp}'}{a_{\perp}} (K^x_x - K^y_y) - \frac{\phi' \Pi_{\phi}}{2a_{\parallel} M_P^2}$$

Hamiltonian Constraint Enforces Energy Conservation

$$\begin{aligned} 0 = \mathcal{H} = & \frac{2a_{\perp}a'_{\parallel}a'_{\perp} - a_{\parallel}a'^2_{\perp} - 2a_{\parallel}a_{\perp}a''_{\perp}}{a_{\parallel}^3 a_{\perp}^2} + 2K^x_x K^y_y + K^y_y{}^2 \\ & - M_P^{-2} \left( \frac{\phi'^2 + \Pi_{\phi}^2}{2a_{\parallel}^2} + V \right) = 0 \end{aligned}$$

# Distribution of Initial Conditions



What's the response in the curvature perturbation?

# Time Evolution

$$\dot{a}_{\parallel} = -a_{\parallel} K^x_x, \quad \dot{a}_{\perp} = -a_{\perp} K^y_y,$$

$$\dot{K^x_x} = \frac{a'_{\perp}^2}{a_{\perp}^2 a_{\parallel}^2} + {K^x_x}^2 - {K^y_y}^2 + \frac{\left(\Pi_{\phi}^2 - \phi'^2\right)}{2a_{\parallel}^2 M_P^2},$$

$$\dot{K^y_y} = -\frac{a'_{\perp}^2}{2a_{\perp}^2 a_{\parallel}^2} + \frac{3}{2} {K^y_y}^2 - \frac{V(\phi)}{2M_P^2} + \frac{\left(\Pi_{\phi}^2 + \phi'^2\right)}{4a_{\parallel}^2 M_P^2},$$

$$\dot{\Pi}_{\phi} = 2K^y_y \Pi_{\phi} + \frac{1}{a_{\parallel}} \phi'' + \left( \frac{2a'_{\perp}}{a_{\parallel} a_{\perp}} - \frac{a'_{\parallel}}{a_{\perp}^2} \right) \phi' - a_{\perp} \partial_{\phi} V(\phi),$$

$$\dot{\phi} = \frac{\Pi_{\phi}}{a_{\parallel}}$$

# Numerical Approach

## Machine precision accuracy

- ▶ Gauss-Legendre time-integrator ( $\mathcal{O}(dt^{10})$ , symplectic)
- ▶ Fourier pseudospectral discretisation (exponential convergence)

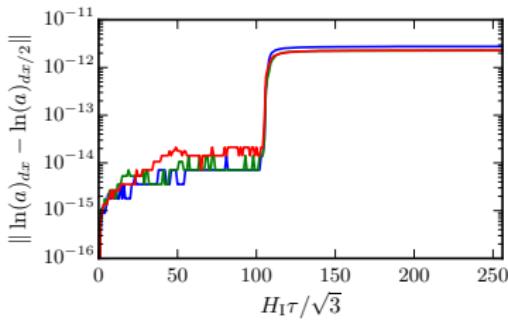
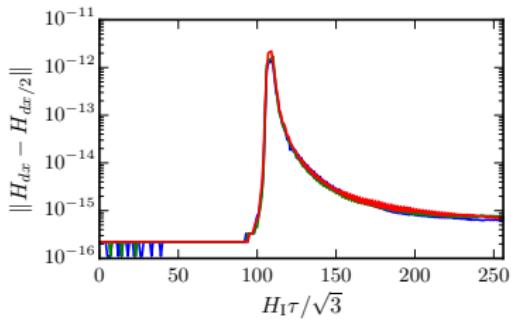
## Fast to allow sampling

- ▶ Adaptive time-stepping
- ▶ Adaptive grid spacing

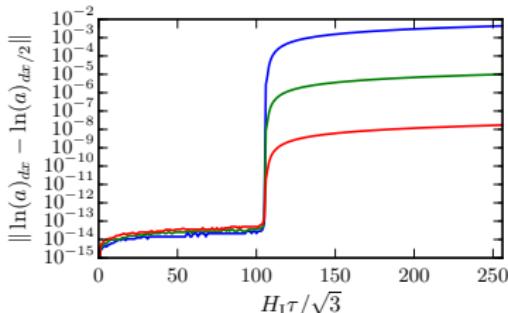
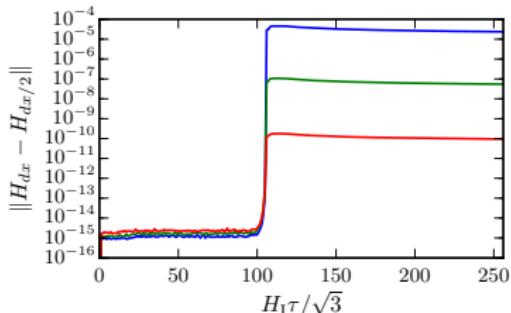
**$\mathcal{O}(1 - 10)\text{s}$  to evolve through 60 e-folds of inflation**

# Convergence Testing : Dynamical Variables

Vary grid spacing  $dx$  at fixed  $dx/dt$

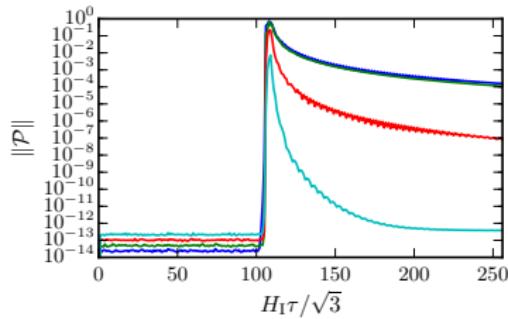
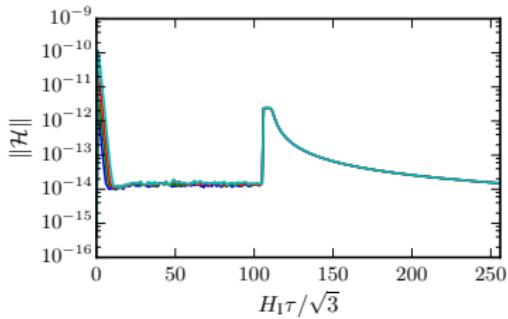


Vary time-step  $dt$  at fixed  $dx$

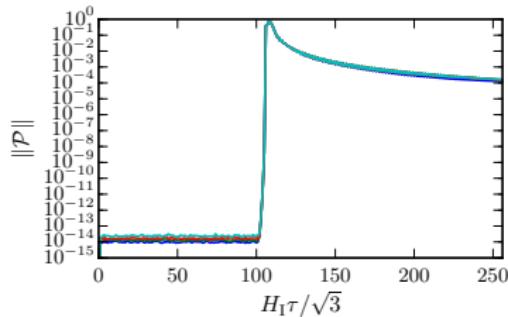
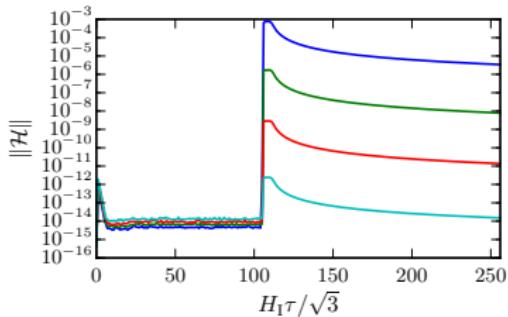


# Convergence Testing : Constraint Preservation

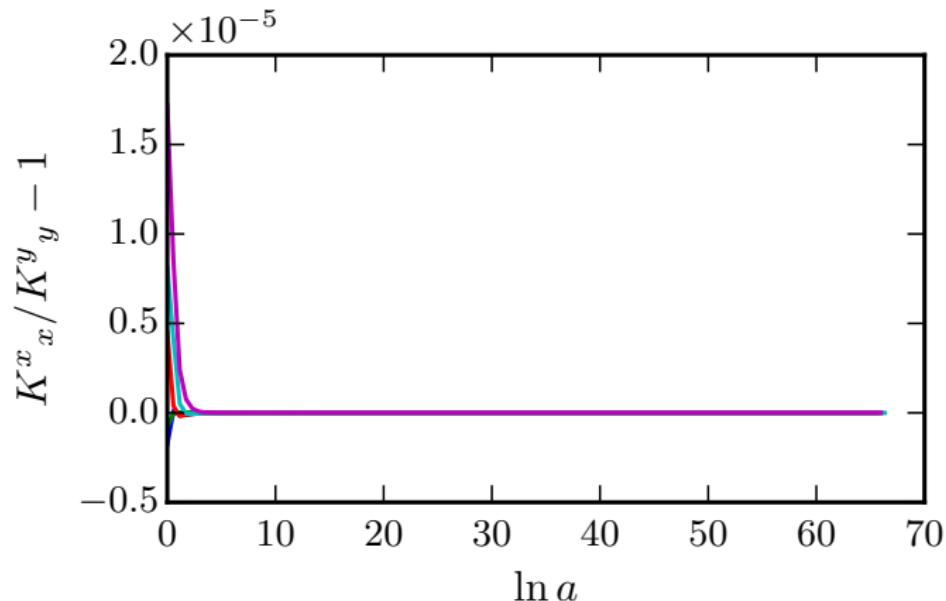
Vary grid spacing  $dx$  at fixed  $dx/dt$



Vary time-step  $dt$  at fixed  $dx$

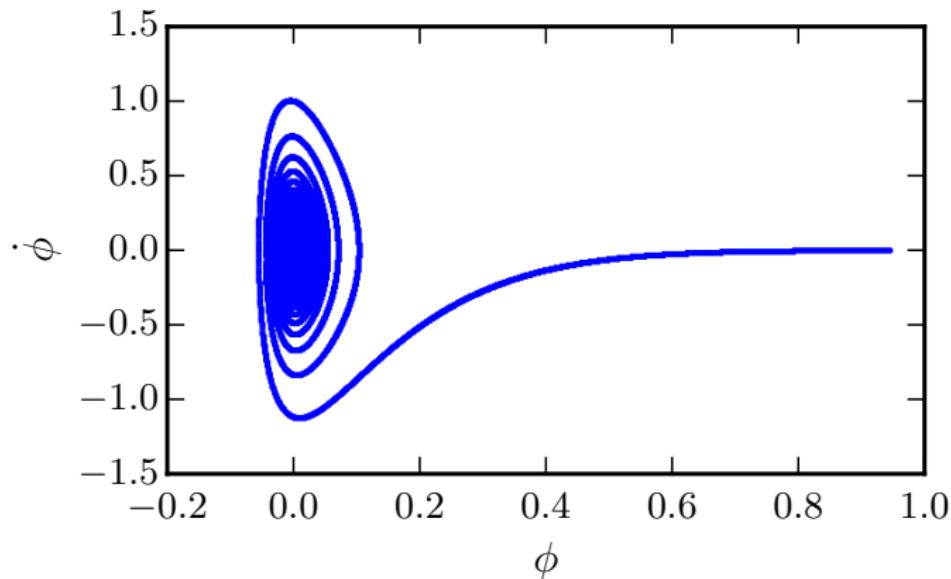


# Self-Gravitating Dynamics: Isotropisation of Expansion

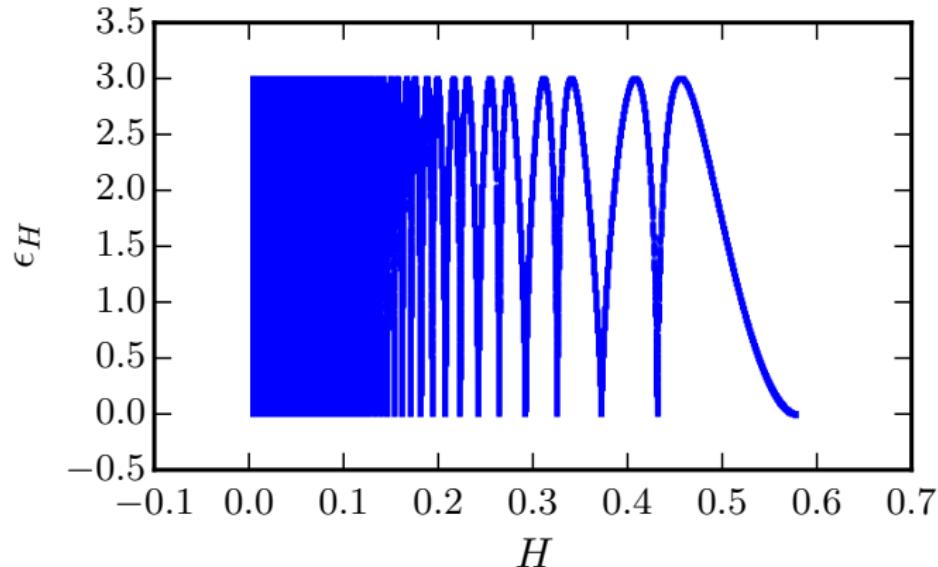


Expansion quickly isotropises itself

# Self-Gravitating Dynamics: Attractor Behaviour

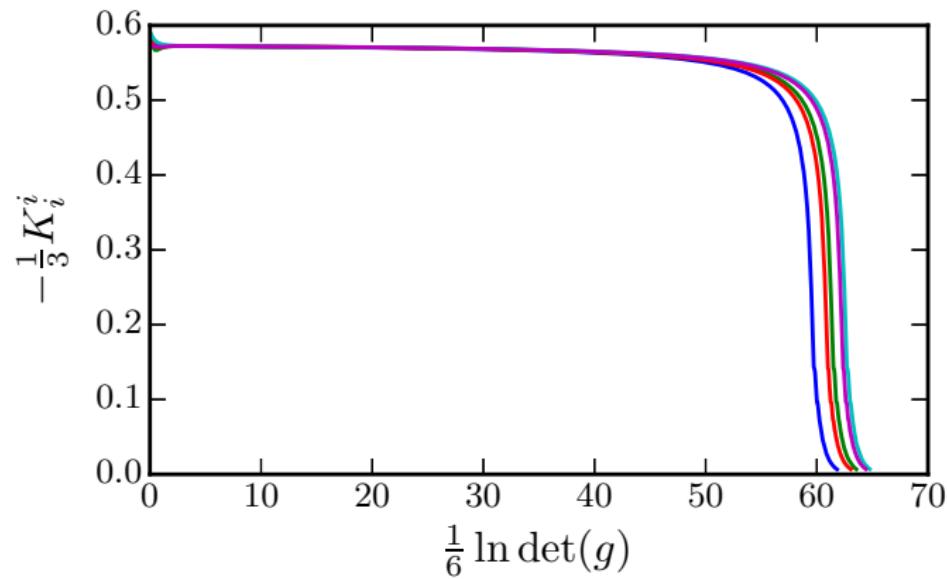


## Self-Gravitating Dynamics: Geometry Attractor



Individual Lattice Sites Evolve Along Attractor  
Fixed  $H$ ,  $\epsilon_H$  are equivalent

# Self-Gravitating Dynamics: Spatially Dependent Expansion History



## Observations: Relation to Local Multipoles

$$\zeta(H) \equiv \delta \ln a|_{a_{\parallel}=1=a_{\perp}}^H$$

Large-scale perturbations →  
evaluate  $\zeta$  on last-scattering surface around a point  $x_0$

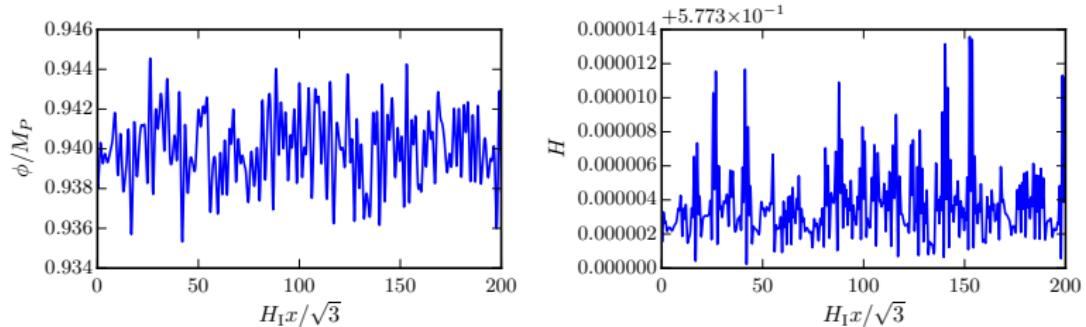
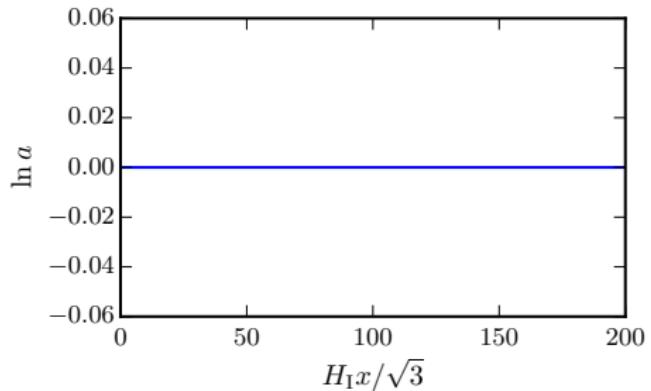
Expand for Large-Scale Fluctuations

$$\zeta(x_0 + r_{ls}) \approx \zeta(x_0) + (H_I r_{ls}) \frac{\partial \zeta}{\partial (H_I x_p)}(x_0) + \frac{(H_I r_{ls})^2}{2} \frac{\partial^2 \zeta}{\partial (H_I x_p)^2}(x_0) + \dots$$

Matches onto Multipoles

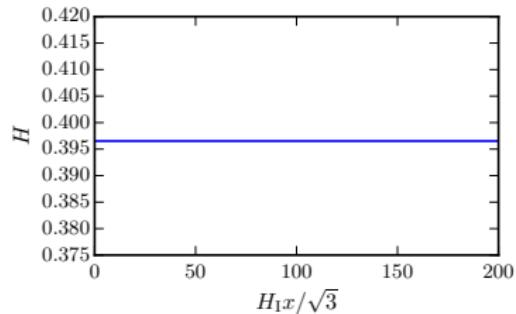
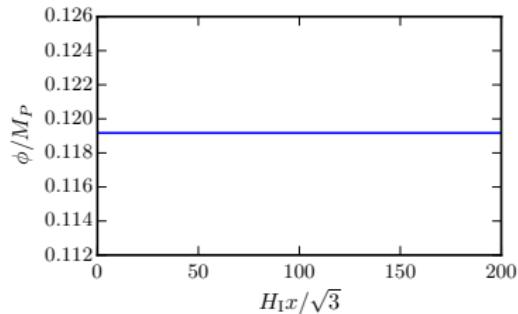
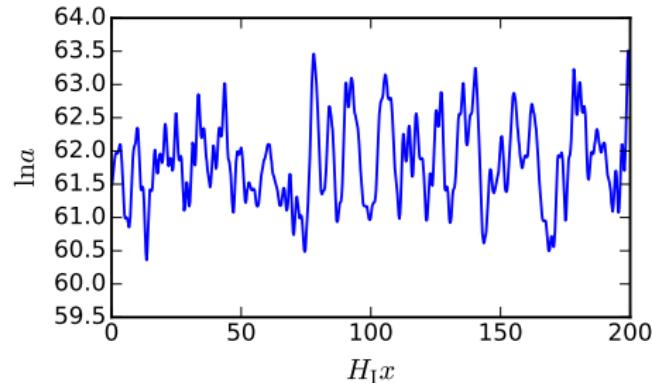
$$a_{20}^{(UL)} \approx F(L_{\text{obs}} H_I)^2 \partial_{x_p}^2 \zeta \simeq F(L_{\text{obs}} H_I)^2 \frac{1}{a_{\parallel}} \partial_x \left[ \frac{\partial_x \zeta}{a_{\parallel}} \right]$$

# Required Evolution



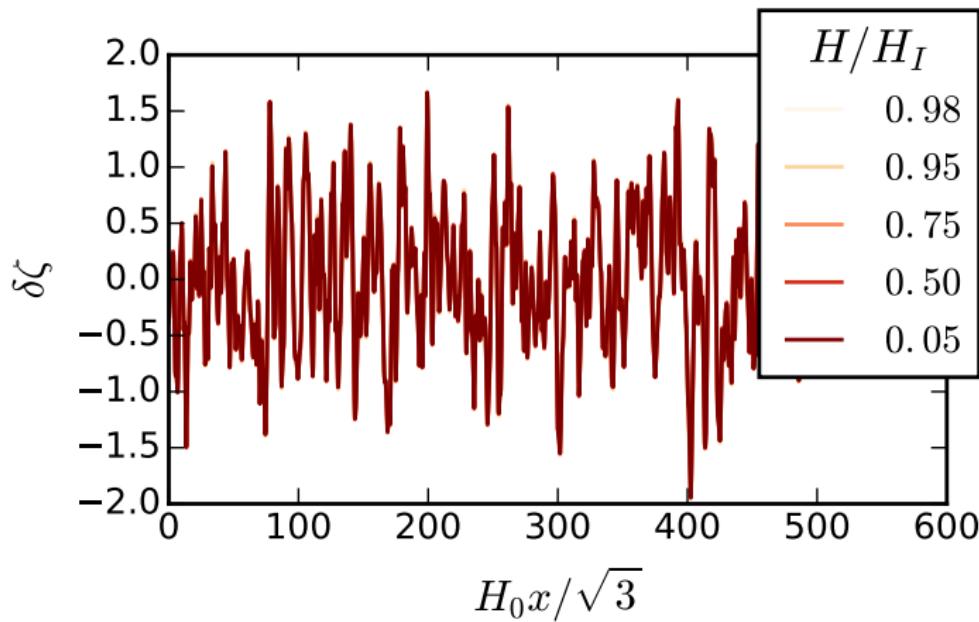
Initial Conditions ( $\tau = 0$ )

# Required Evolution



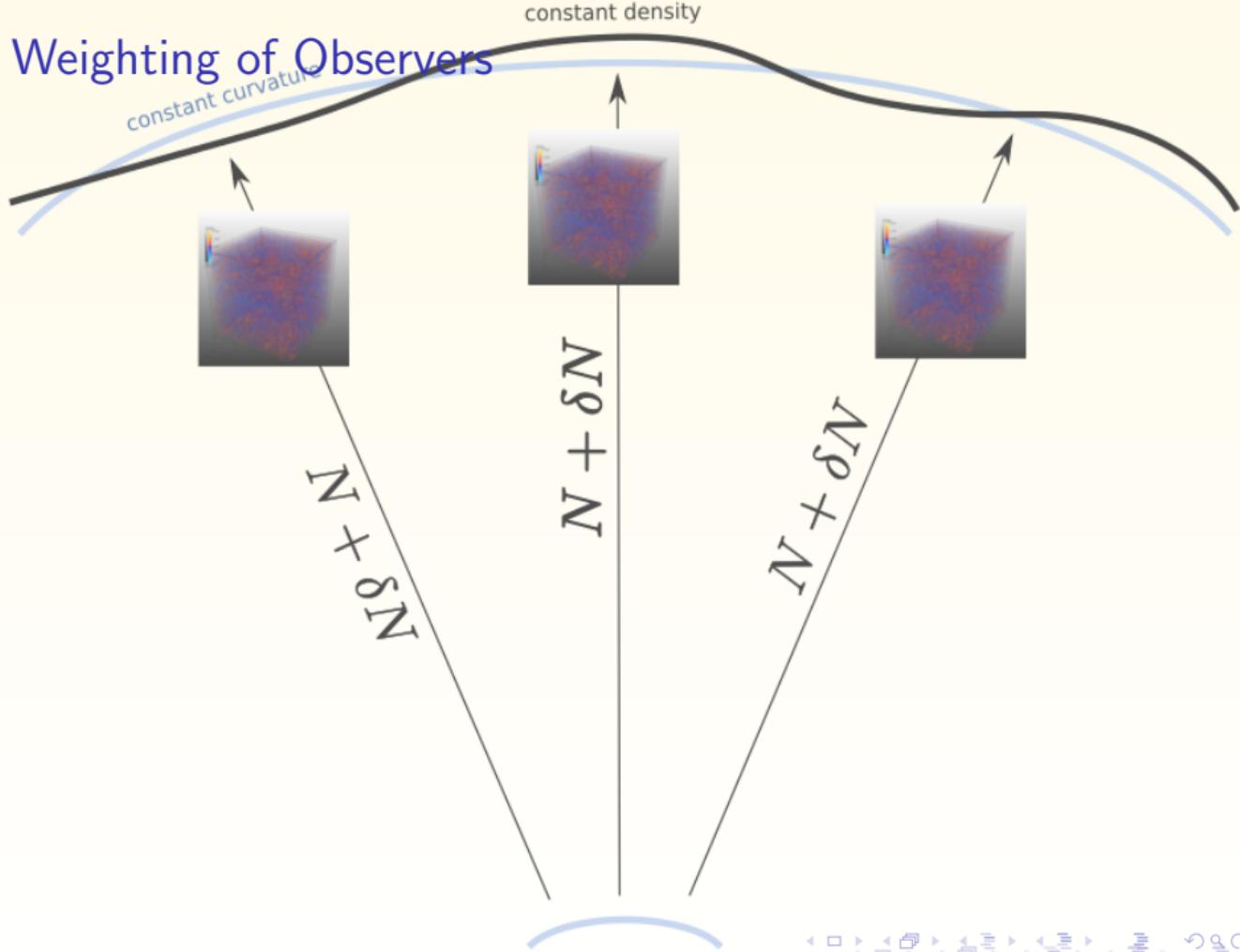
End of Inflation ( $\epsilon_H = -d \ln H / d \ln a = 1$ )

# Freeze-in of Superhorizon Perturbations

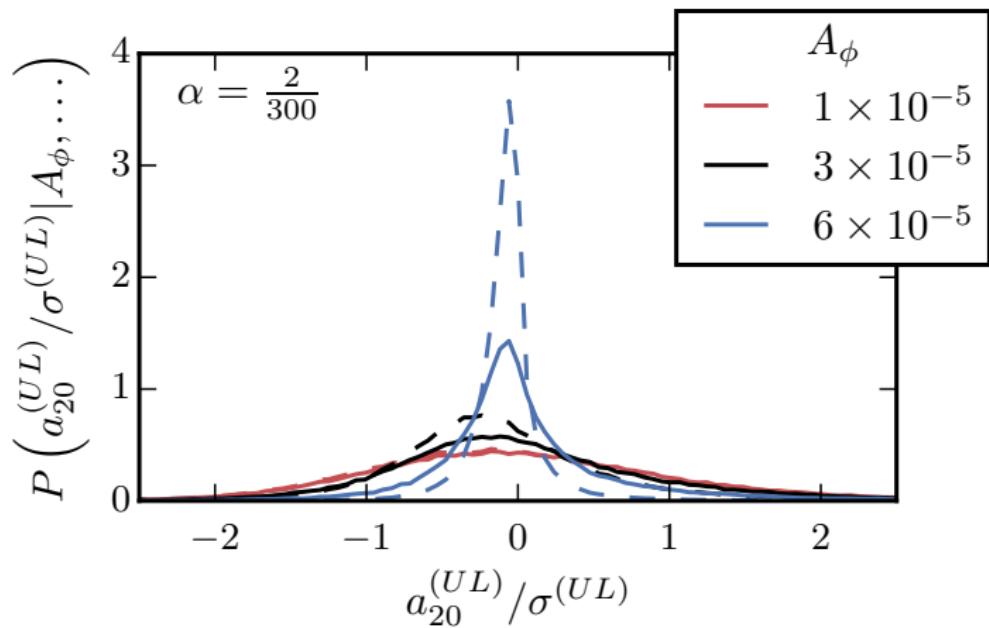


Initial transient modifies result from separate universe approximation

# Weighting of Observers



# Distribution of Superhorizon $a_{20}$



# Statistical Framework

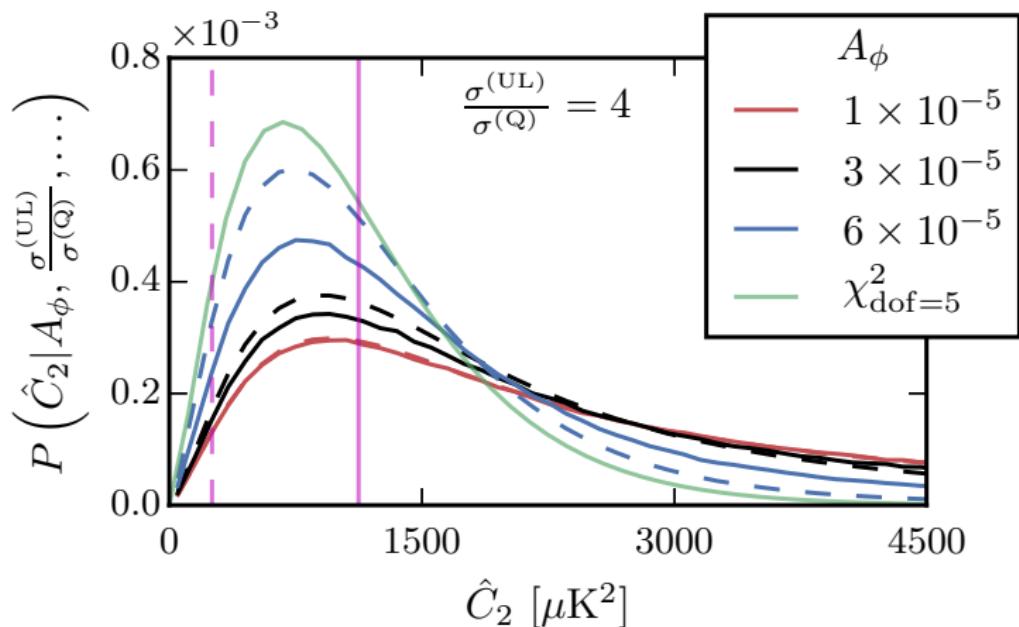
$$\Pr(A_\phi, L_{\text{obs}} | C_2^{\text{obs}}) \propto \mathcal{L}(A_\phi, L_{\text{obs}}, \dots) \Pr(A_\phi, L_{\text{obs}}, \dots)$$
$$\mathcal{L}(A_\phi, L_{\text{obs}}) = \Pr(C_2^{\text{obs}} | A_\phi, L_{\text{obs}}, \dots)$$

Planck high- $\ell$  “theory” and low- $\ell$  “observed”

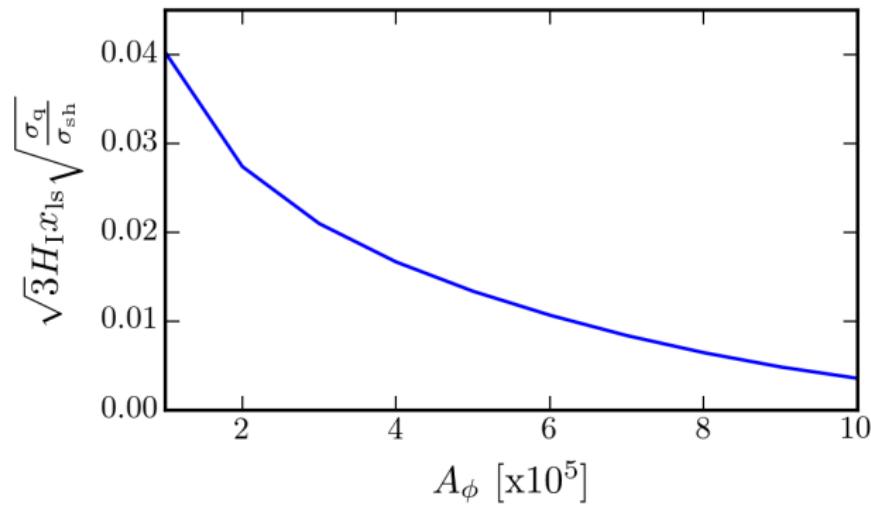
$$C_2^{\text{obs}} = 253.6 \mu\text{K}^2 \quad C_2^{\text{theory}} = 1124.1 \mu\text{K}^2$$

# Distribution of CMB quadrupole

$$\hat{C}_2 = \frac{1}{5} \left[ \left( a_{20}^{(\text{UL})} + a_{20}^{(\text{Q})} \right)^2 + \sum_{m=-2, m \neq 0}^2 \left( a_{2m}^{(\text{Q})} \right)^2 \right]$$

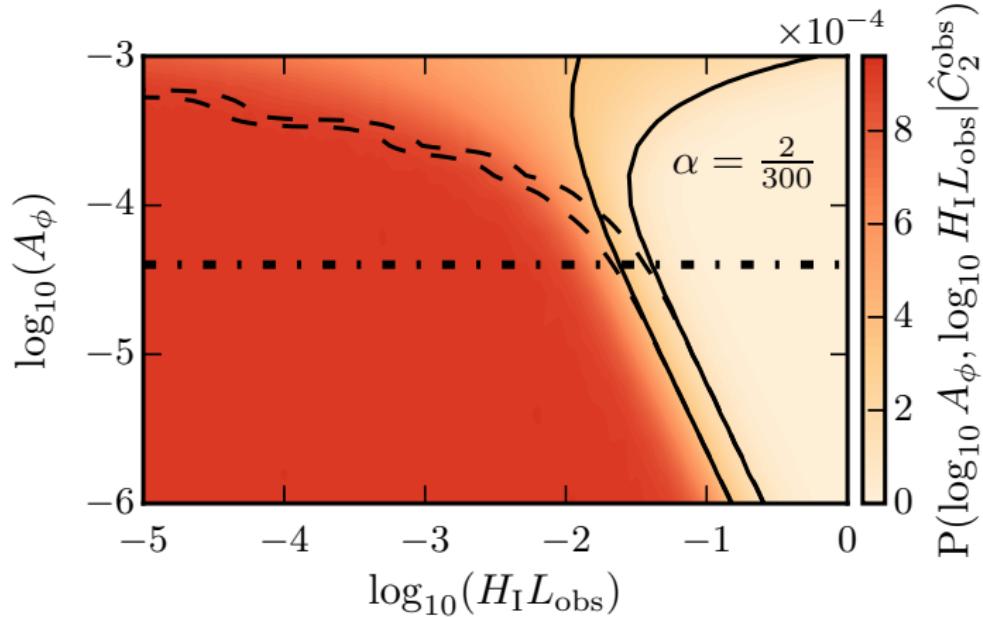


Relative RMS of superhorizon contribution depends of distance to last-scattering surface



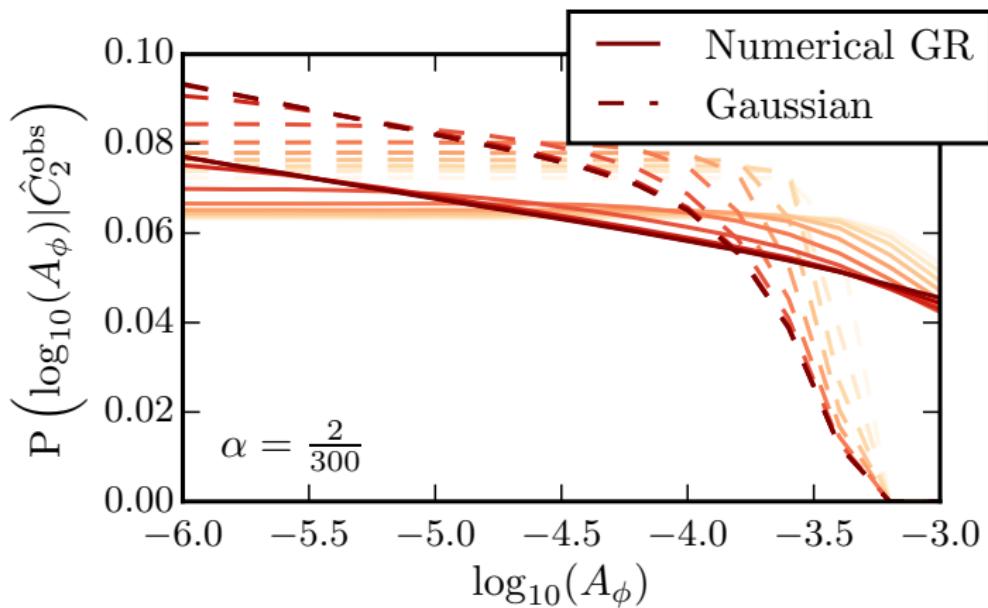
$$a_{20}^{(UL)} \approx F(L_{\text{obs}} H_I)^2 \partial_{x_p}^2 \zeta \simeq F(L_{\text{obs}} H_I)^2 \frac{1}{a_{||}} \partial_x \left[ \frac{\partial_x \zeta}{a_{||}} \right]$$

# Constraints from the CMB Quadrupole



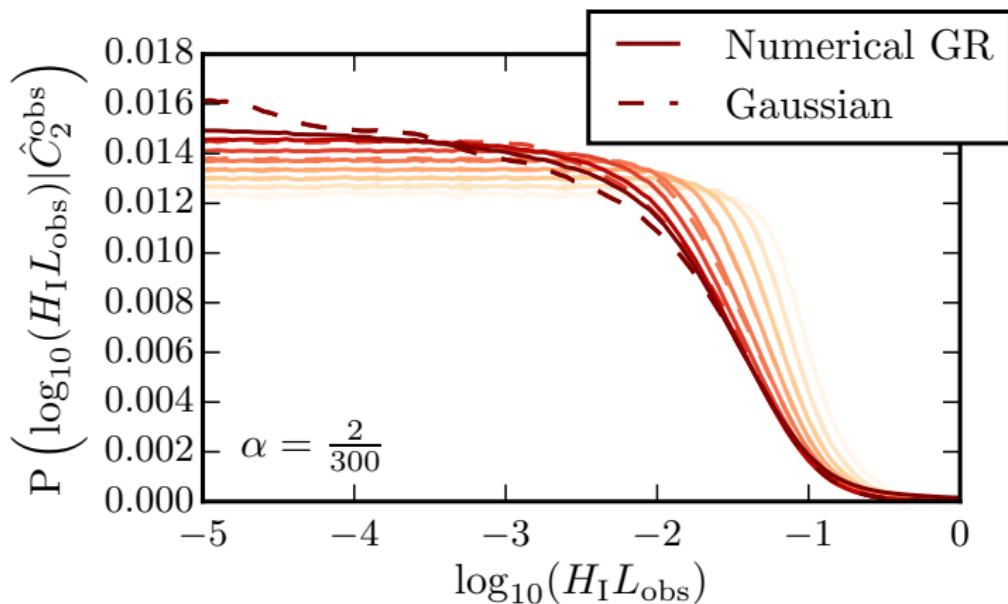
**Strong Modification from Gaussian Ansatz at Large  $A_\phi$**

## Marginalised Constraints



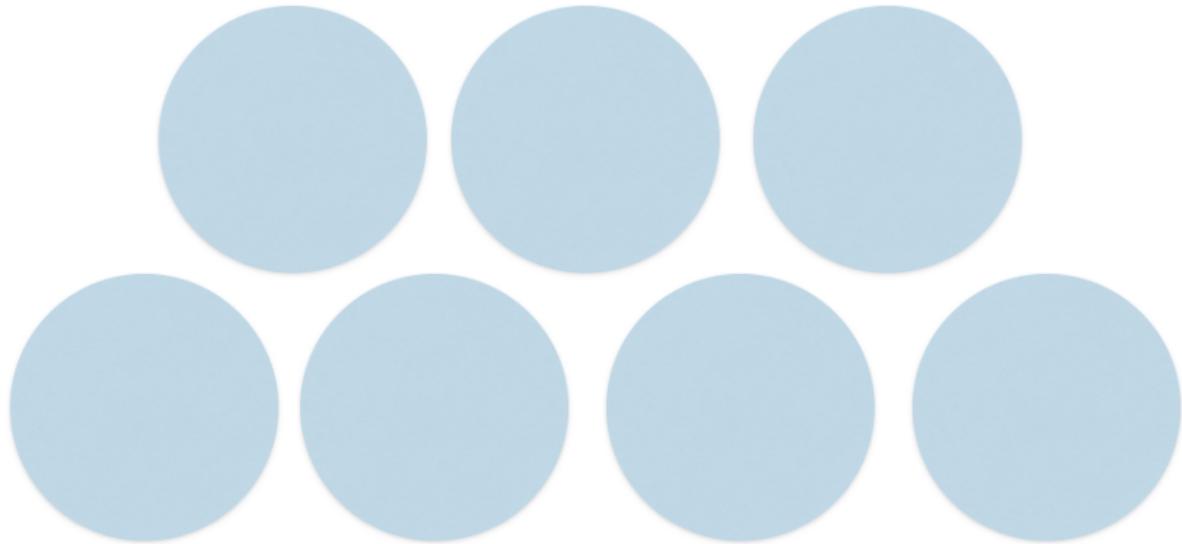
Constraint of  $A_\phi$  as we marginalise over  $L_{\text{obs}}$

## Marginalised Constraints

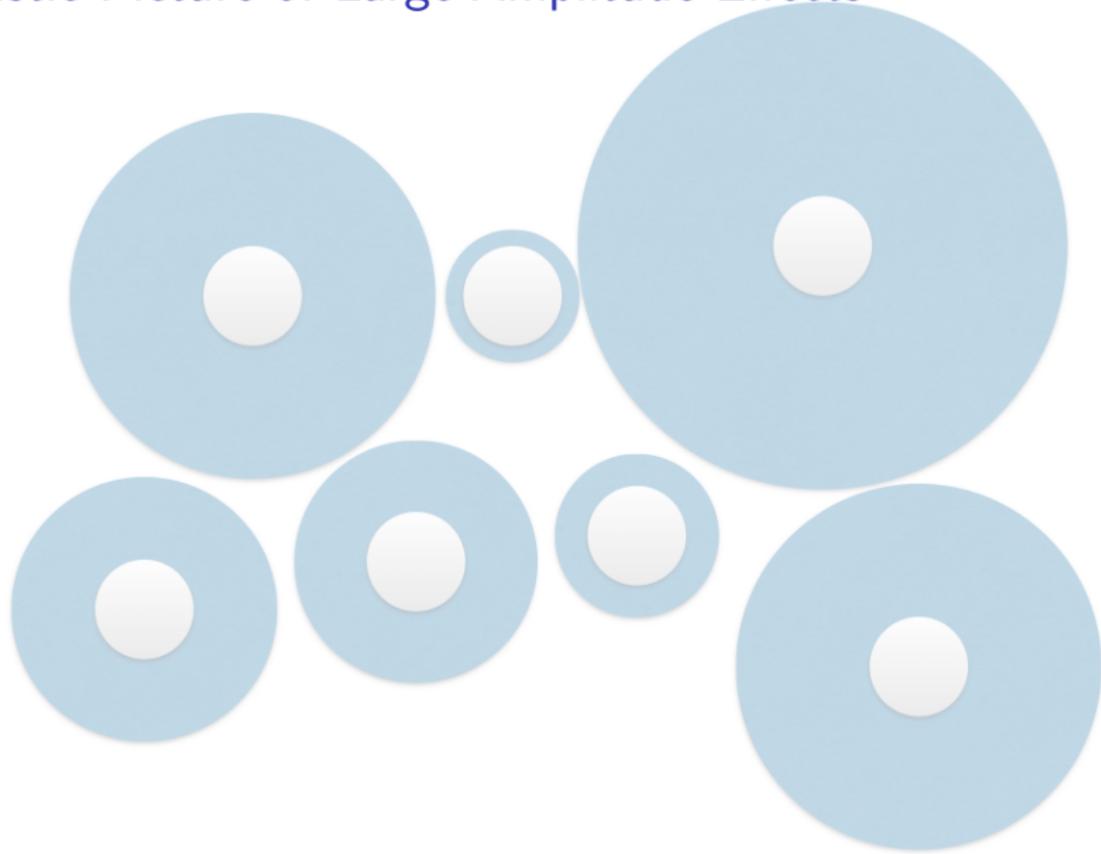


Constraint of  $L_{\text{obs}}$  as we marginalise over  $A_\phi$

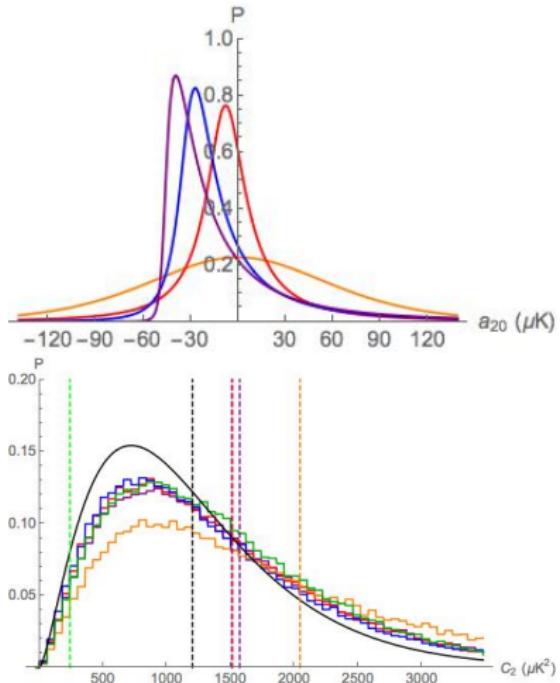
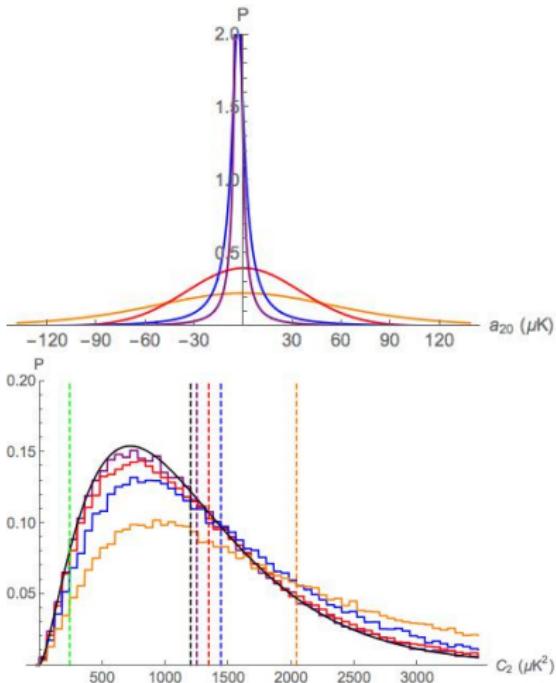
# Heuristic Picture of Large Amplitude Effects



# Heuristic Picture of Large Amplitude Effects

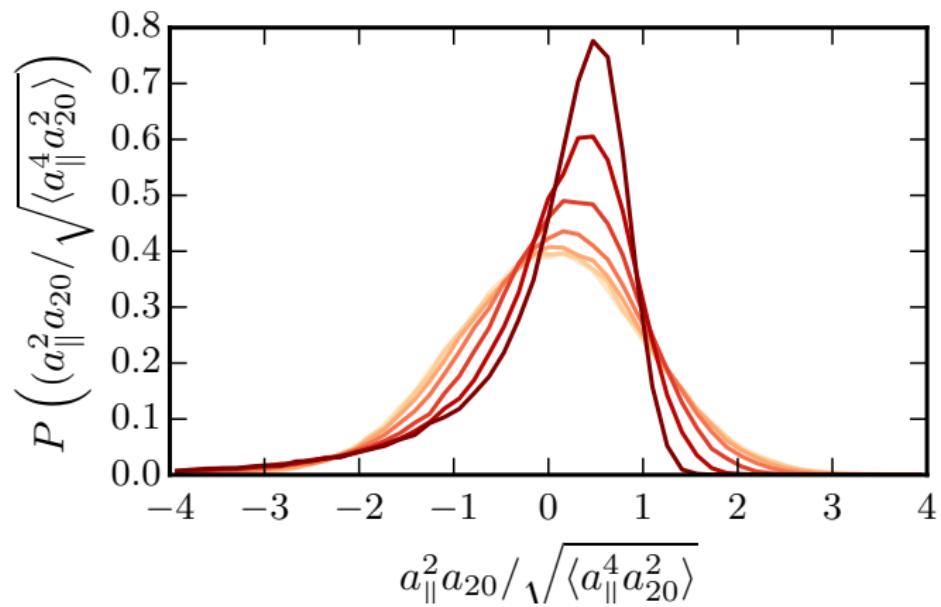


# Modelling With Toy Distributions



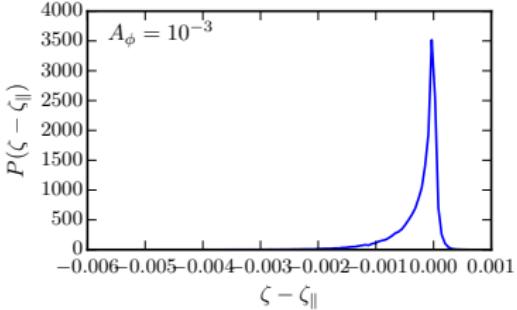
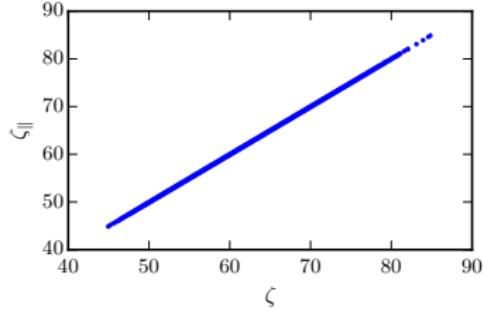
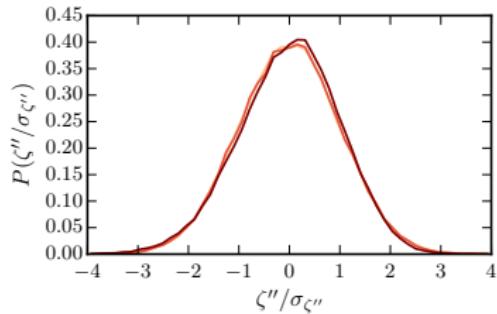
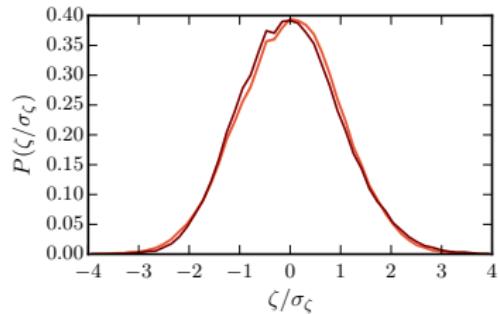
Model behaviour with Johnson distribution

# Removing Inhomogeneous Expansion



# Statistics of $\zeta$

$\zeta$  and comoving derivatives nearly Gaussian



# Analytical Modelling

$$\zeta(H, x_{\text{com}}) \sim \text{GRF}$$

## Large-Scale Approximation for $a_{20}$

$$a_{20}(x_0) \approx A e^{-2\zeta_{||}(x_0)} \left( \zeta''(x_0) - \zeta'_{||}(x_0)\zeta'(x_0) \right)$$

$\zeta, \zeta', \zeta''$  are correlated Gaussian random deviates, with covariance

$$C_\zeta = \begin{bmatrix} \sigma_0^2 & 0 & -\sigma_1^2 \\ 0 & \sigma_1^2 & 0 \\ -\sigma_1^2 & 0 & \sigma_2^2 \end{bmatrix}$$

$$\sigma_i = \int dk k^{2i} \left\langle \left| \tilde{\zeta}_k \right|^2 \right\rangle$$

# Analytical Modelling

$$\zeta(H, x_{\text{com}}) \sim \text{GRF}$$

## Large-Scale Approximation for $a_{20}$

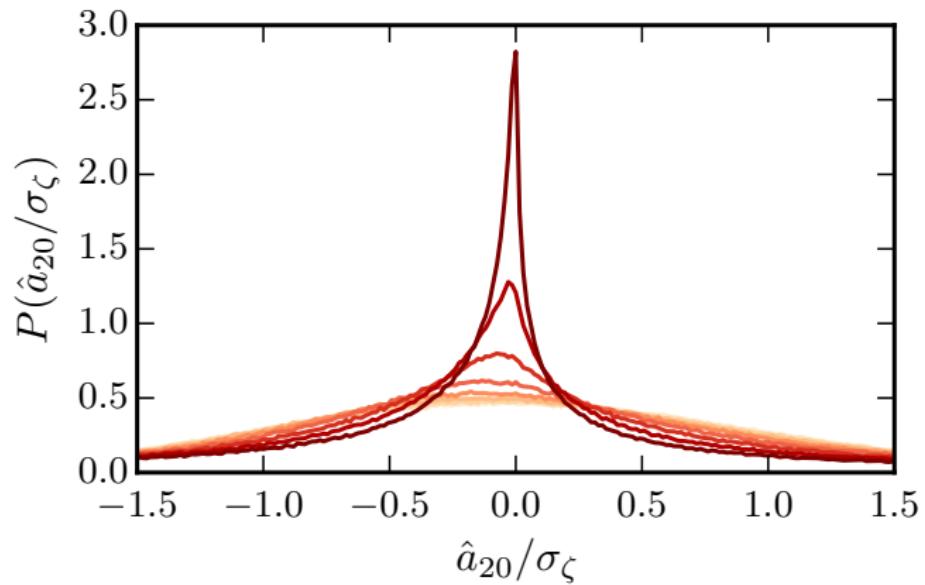
$$a_{20}(x_0) \approx A e^{-2\zeta(x_0)} (\zeta''(x_0) - \zeta'(x_0)^2)$$

$\zeta, \zeta', \zeta''$  are correlated Gaussian random deviates, with covariance

$$C_\zeta = \begin{bmatrix} \sigma_0^2 & 0 & -\sigma_1^2 \\ 0 & \sigma_1^2 & 0 \\ -\sigma_1^2 & 0 & \sigma_2^2 \end{bmatrix}$$

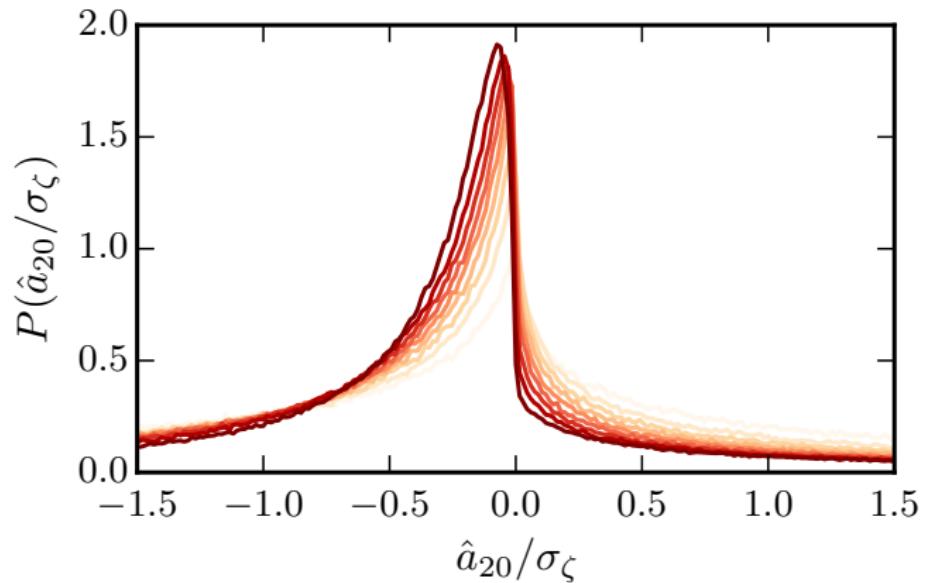
$$\sigma_i = \int dk k^{2i} \left\langle \left| \tilde{\zeta}_k \right|^2 \right\rangle$$

## Anaytic $a_{20}$ Distributions



Vary  $\sigma_\zeta$  at fixed  $\sigma_{\zeta^{(p)}}/\sigma_\zeta$

## Anaytic $a_{20}$ Distributions



Vary  $\sigma_{\zeta'}/\sigma_\zeta$  at fixed  $\sigma_\zeta$  and  $\sigma_{\zeta''}/\sigma_\zeta$

## Future Steps

- ▶ Three-dimensional simulations (in progress)
- ▶ Multi-field models (additional isocurvature modes, non-attractor, etc.)
- ▶ Inclusion of stochastic effects from subhorizon fluctuations (technical challenge)
- ▶ Correlated anomalies

# Conclusions

- ▶ Fully relativistic treatment of highly inhomogeneous initial conditions for inflation
- ▶ Constraints obtained from CMB quadrupole based on full Bayesian analysis
- ▶ Gravitational nonlinearities extremely important in determining final distributions of observables
- ▶ Very large amplitude initial inhomogeneities more poorly constrained than intermediate amplitudes
- ▶ Comoving curvature perturbation  $\zeta$  approximately GRF in appropriate coordinate system