Constraining the Ultra-Large Scale Structure of the Universe Using Numerial Relativity

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w/ Hiranya Peiris, Matthew Johnson, and Anthony Aguirre based on arXiv:1604.04001 and *in progress*



Ultra-Large Scale Structure



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Ultra-Large Scale Structure



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Local Remnants of Ultra-Large Scale Structure?

Ultra-Large Scale Structure



Local Remnants of Ultra-Large Scale Structure?

- Structure present at start of inflation
- Conversion of structure during or after inflation

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Modelling Initial Conditions

Monte Carlo Sampling: Planar Symmetry

$$ds^{2} = -d\tau^{2} + a_{\parallel}^{2}(x,\tau)dx^{2} + a_{\perp}^{2}(x,\tau)(dy^{2} + dz^{2})$$

Inflaton on $a_{\parallel}(au=0)=1=a_{\perp}(au=0)$

$$\phi(x)=ar{\phi}+\delta \hat{\phi}$$
 $ar{\phi}$ gives ${\cal N}$ e-folds $3H_{
m I}^2\equiv V(ar{\phi})$

Field Fluctuations

$$\delta\hat{\phi}(x_i) = A_{\phi} \sum_{n=1} \hat{G} e^{ik_n x_i} \sqrt{P(k_n)} \qquad \hat{G} = \sqrt{-2 \ln \hat{\beta} e^{2\pi i \hat{\alpha}}}$$

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Modelling Initial Conditions

Monte Carlo Sampling: Planar Symmetry

$$ds^2 = -d au^2 + a_{\parallel}^2(x, au)dx^2 + a_{\perp}^2(x, au)\left(dy^2 + dz^2\right)$$

Inflaton on $a_{\parallel}(au=0)=1=a_{\perp}(au=0)$

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m I}^2\equiv V(ar{\phi})$

Field Fluctuations

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$$P(k) = \Theta(k_{\max} - k) \qquad H_{\rm I}^{-1} k_{\max} = 2\pi\sqrt{3}$$

$$\mathcal{V}(\phi) = \mathcal{V}_0 \left(1 - e^{-\sqrt{rac{2}{3lpha}}rac{\phi}{M_P}}
ight)^2$$

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ight)^2$$



• $\alpha \ll 1 \rightarrow$ small-field model

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$$\mathcal{V}(\phi) = \mathcal{V}_0 \left(1 - e^{-\sqrt{rac{2}{3lpha}}rac{\phi}{M
ho}}
ight)^2$$



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- $\alpha \ll 1 \rightarrow$ small-field model
- $\alpha = 1 \rightarrow \text{Starobinski}$

$$V(\phi) = V_0 \left(1 - e^{-\sqrt{rac{2}{3lpha}}rac{\phi}{M_P}}
ight)^2$$



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- $\alpha \ll 1 \rightarrow$ small-field model
- $\blacktriangleright \ \alpha = \mathbf{1} \to \mathsf{Starobinski}$
- $\blacktriangleright \ \alpha \gg 1 \rightarrow m^2 \phi^2/2$

$$V(\phi) = V_0 \left(1 - e^{-\sqrt{rac{2}{3lpha}}rac{\phi}{M_P}}
ight)^2$$



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Required Evolution



Initial Conditions ($\tau = 0$)

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Required Evolution



End of Inflation ($\epsilon_H = -d \ln H/d \ln a = 1$)

SQR

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Solving Einstein's Equations for a Self-Gravitating Inflaton

Machine precision accuracy

- Gauss-Legendre time-integrator ($\mathcal{O}(dt^{10})$, symplectic)
- Fourier pseudospectral discretisation (exponential convergence)

Fast to allow sampling

- Adaptive time-stepping
- Adaptive grid spacing

 $\mathcal{O}(1s-10s)$ to evolve through 60 e-folds of inflation Machine precision convergence and constraint preservations

Observational Constraints

$$\begin{split} \Pr(A_{\phi}, H_{I}L_{\rm obs} | C_{2}^{\rm obs}, \dots) \propto \mathcal{L}(A_{\phi}, H_{I}L_{\rm obs}) \Pr(A_{\phi}, H_{I}L_{\rm obs} | \dots) \\ \mathcal{L} = \Pr(C_{2}^{\rm obs} | A_{\phi}, H_{I}L_{\rm obs}, \dots) \end{split}$$

- A_{ϕ} : Fluctuation Amplitude $P(k) \propto A_{\phi}^2$
- $H_I L_{\rm obs}$: Uncertain post-inflation expansion history
- ... : $V(\phi)$, spectrum shape, IC hypersurface, $C_2^{\mathrm{high}-\ell}$, etc.

Planck measured C_{ℓ}

$$C_2^{\rm obs} = 253.6 \mu K^2$$
 $C_2^{\rm high-\ell} = 1124.1 \mu K^2$

Numerical GR Input

$$\Pr\left(\hat{C}_2|A_{\phi},H_1L_{obs},\dots\right)$$





Evaluation of CMB Quadrupole

$$\hat{C}_2 = \frac{1}{5} \left[\left(a_{20}^{(\mathrm{UL})} + a_{20}^{(\mathrm{Q})} \right)^2 + \sum_{m=-2, m \neq 0}^{m=2} \left(a_{2m}^{(\mathrm{Q})} \right)^2 \right]$$

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Evaluation of CMB Quadrupole

$$\hat{C}_{2} = \frac{1}{5} \left[\left(a_{20}^{(\text{UL})} + a_{20}^{(\text{Q})} \right)^{2} + \sum_{m=-2, m \neq 0}^{m=2} \left(a_{2m}^{(\text{Q})} \right)^{2} \right]$$

$$a^{
m (Q)}_{2m}$$
 : Gaussian with $\langle (a^{
m (Q)}_{2m})^2
angle = 1124.1 \mu K^2$



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Evaluation of CMB Quadrupole

$$\hat{C}_{2} = \frac{1}{5} \left[\left(a_{20}^{(\text{UL})} + a_{20}^{(\text{Q})} \right)^{2} + \sum_{m=-2, m \neq 0}^{m=2} \left(a_{2m}^{(\text{Q})} \right)^{2} \right]$$

 $a_{20}^{(\text{UL})}$: NR simulations



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Dependence of C_2 on Model Parameters



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Dependence of C_2 on Model Parameters



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Final Posterior



Significant deviations from Gaussian approximation

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Marginalised Constraints on Model Parameters: Amplitude



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Heuristic Explanation: Initial Conditions



Heuristic Explanation: End-of-Inflation



Analytic Approximation



 ζ and comoving derivatives nearly Gaussian

Treat as Gaussian random field

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Analytic Approximation



 ζ and comoving derivatives nearly Gaussian

Treat as Gaussian random field

Large-Scale Approximation for a_{20}

$$a_{20}(x_0) \approx \mathcal{A}e^{-2\zeta(x_0)} \left(\zeta''(x_0) - \zeta'(x_0)^2\right)$$

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Anaytic *a*₂₀ Distributions



Vary σ_{ζ} at fixed $\sigma_{\zeta^{(p)}}/\sigma_{\zeta}$

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Conclusions

- Numerical relativity is a useful framework for making cosmological predictions
 - Sometimes it is a *necessary* tool (deviations from Gaussianity)

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- Robust qualitative conclusions over a variety of inflationary models
- Inflation is effective at hiding large amplitude initial fluctuations
- Gaussianity of ζ in comoving coordinates suggests analytic approach in 3D