

very early U

early to middle to now U

very late U

inflation

string theory/landscape/higher dimensions

dark energy

$V_{\text{eff}}(\psi_{\text{inf}})$? partial shape reconstruction

reconstruct gradient $V_{\text{eff}}(\psi_{\text{inf}})$?

$K_{\text{eff}}(\psi_{\text{inf}})$?

$K_{\text{eff}}(\psi_{\text{inf}})$?

trajectory probability

$-\frac{d \ln \rho_{\text{tot}}}{d \ln a} / 2 \Rightarrow P_s, P_t, V_{\text{eff}}(k),$

$= \mathcal{E}(k) = 1 + q, k \sim H a$

$\psi_{\text{inf}}(k)$

trajectory probability:

~ 1 e-fold \Rightarrow blind is bad

\Rightarrow slow-to-moderate roll

$-\frac{d \ln \rho_{\phi}}{d \ln a} / 2$

$= \mathcal{E}_{\phi}(a) = (1 + w) / 2$

$\epsilon_s = (d \ln V / d \psi)^2 / 4$

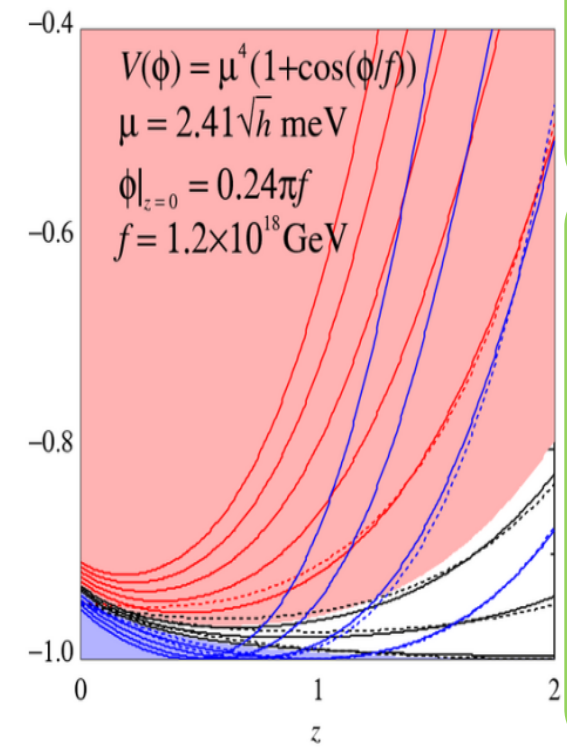
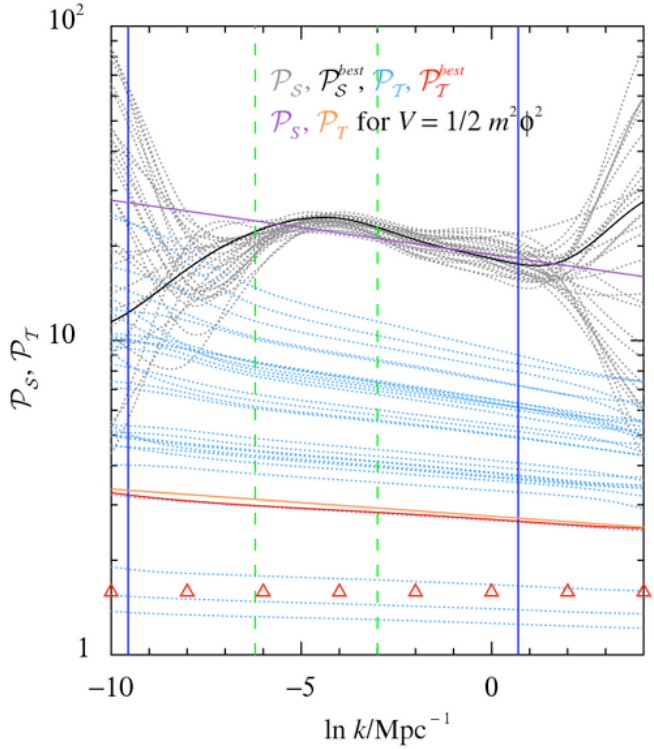
@pivot a_{eq} **yes**

$\zeta_s = \pm 1.001 \frac{d^2 \ln V}{d \psi^2} / 4$

@pivot a_{eq}

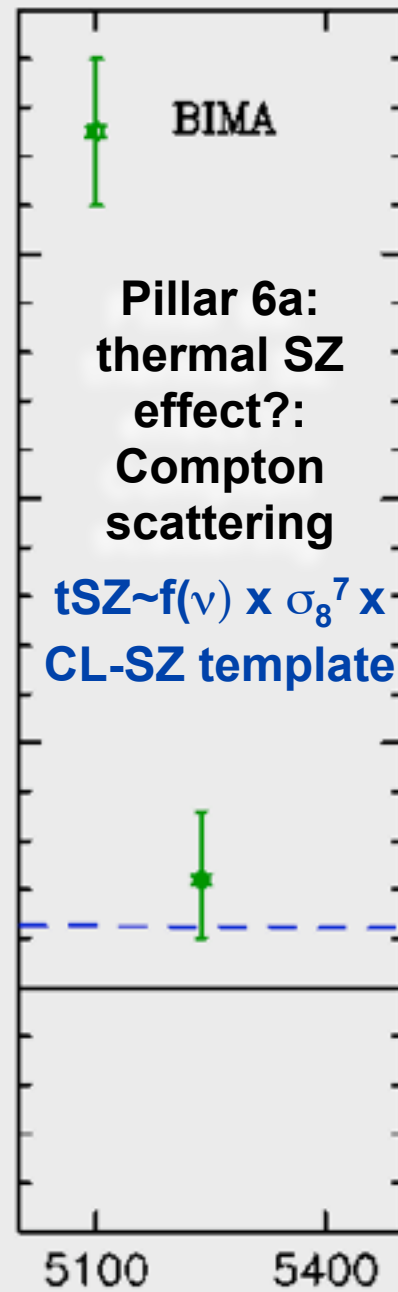
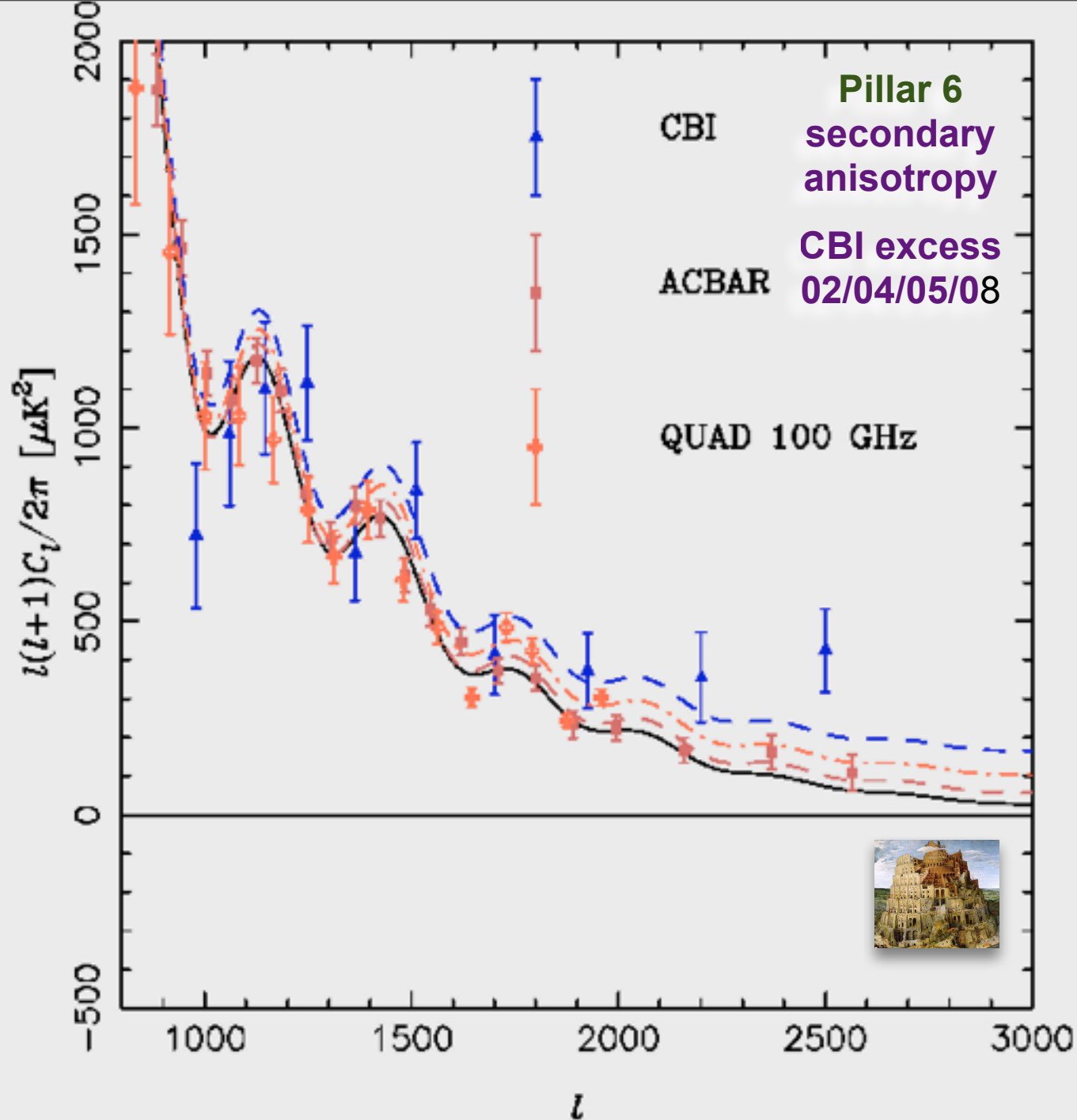
$\zeta_s = \frac{d \ln \epsilon_s}{d \ln a} \times 1/2$

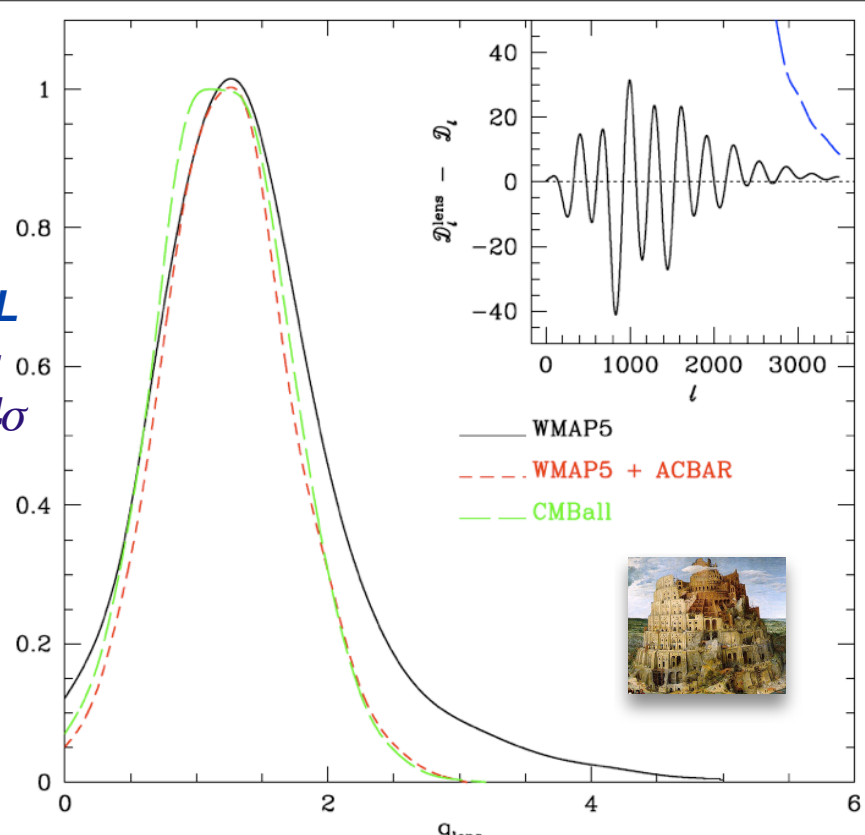
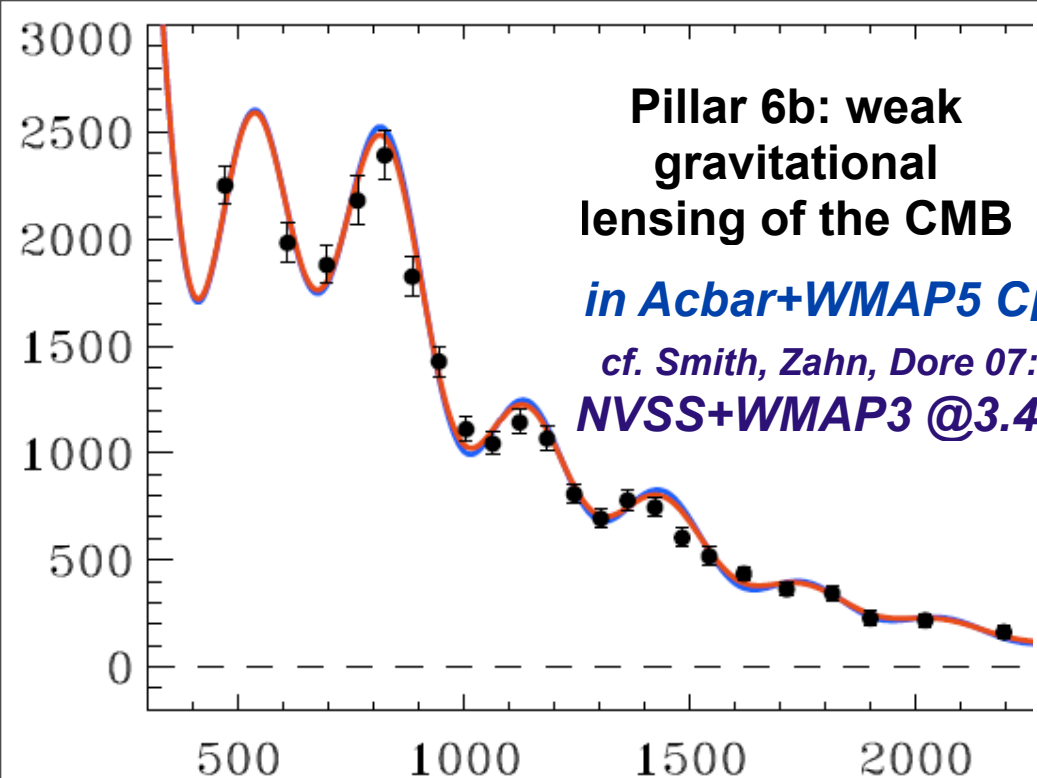
@pivot a_{eq}



Inflation now summary

- the data cannot determine more than 2 w-parameters (1-cs, ind?). general higher order Chebyshev or spline expansion in $1+w$ as for "inflation-then" $\epsilon=(1+q)$ is not that useful. **Parameter eigenmodes** show what is probed
- Any $w(a)$ leads to a viable DE model. The $w(a)=w_0+w_a(1-a)$ phenomenology requires baroque potentials
- Philosophy of HBK08: **backtrack from now ($z=0$) all w-trajectories arising from quintessence ($\epsilon_s > 0$) and the phantom equivalent ($\epsilon_s < 0$); use a 3-parameter model $\epsilon_{\phi} = (1+w(a))^{3/2} = \epsilon_s f(a/a_{\Lambda eq}; a_s/a_{\Lambda eq}; \zeta_s)$ to well-approximate even rather baroque w-trajectories, as well as thawing & freezing trajectories.**
- We ignore constraints on Q-density from photon-decoupling and BBN because further trajectory extrapolation is needed. Can include via a prior on $\Omega_Q(a)$ at z_{dec} and z_{bbn} .**
- For general slow-to-moderate rolling 2 "dynamical parameters" (a_s, ϵ_s) & Ω_Q describe w to a few %. In early-scaling-exit, the information stored in a_s is erased by Hubble drag over the observable range & w can be described by a single parameter ϵ_s . for baroque w-trajectories, add a 3rd param ζ_s ($d \ln \epsilon_s / d \ln a / 2$) - not-determined now & then. freeze-out w at high z, 4th param
- prior-dependence e.g. $\epsilon_{\phi 0} = \begin{matrix} -0.00^{+0.09} \\ -0.13 \end{matrix}$, a_s near 0, $\epsilon_s = \begin{matrix} 0.00^{+0.20} \\ +0.42 \end{matrix}$ since $\epsilon_{\phi} < 0$ of phantom energy, negative kinetic energy is baroque
- Apr08 observations well-centered around a cosmological constant $\epsilon_s = -0.03 \pm 0.28$ $a_s < 0.36$ ($z_s > 2.0$)
cf. $\epsilon_{\phi 0} = -0.00 \pm 0.09$ if constant, $\epsilon_{\phi 0} = -0.015 \pm 0.30$ if a-linear model
- in Planck1yr-CMB+JDEM-SN+DUNE-WL future ϵ_s to ± 0.07 , a_s to < 0.21 ($z_s > 3.7$)**
- cannot reconstruct the quintessence potential**, just the slope ϵ_s & hubble drag info
- late-inflaton field is $<$ Planck mass, but not by a lot
- DE may couple to matter, 5th force constraints are strong, maybe best hope in determining more about DE (chameleon example of dilaton a la Khoury and Weltman 04)





$$C_{\ell}^{\text{lens}} = C_{\ell}^{\text{no-lens}} + q_{\text{lens}} \Delta C_{\ell}^{\text{lens}}$$

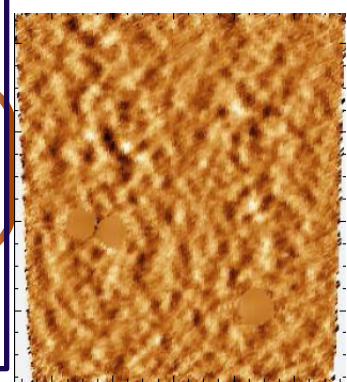
$$\Delta \ln \mathcal{E} = \ln [P(\text{lens} | \text{data}, \text{theory}) / P(\text{no-lens} | \text{data}, \text{theory})]$$

wmap5 $q_{\text{lens}} = 1.34^{+0.27(+1.51)}_{-0.26(-0.85)}$

wmap5+acbar $q_{\text{lens}} = 1.23^{+0.21(+0.83)}_{-0.23(-0.76)}$

CMBall $q_{\text{lens}} = 1.21^{+0.24(+0.82)}_{-0.24(-0.76)}$

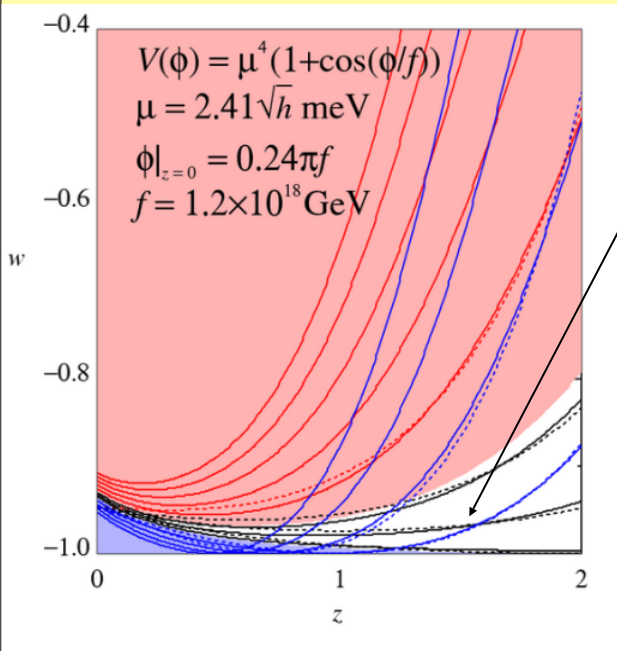
Bayesian evidence	
$\Delta \ln \mathcal{E} = 2.04$	
$\Delta \ln \mathcal{E} = 2.89$	
$\Delta \ln \mathcal{E} = 2.63$	



Late-Inflaton $\epsilon_\phi(a) = \epsilon_s f(a/a_{\Lambda eq}; a_s/a_{\Lambda eq}; \xi_s)$

3-param formula accurately fits **slow-to-moderate roll** & even wild rising baroque late-inflaton trajectories, as well as **thawing & freezing** trajectories. but not oscillating DE

Cosmic Probes Now **CFHTLS SN(Union~300), WL, CMB, BAO, LSS, Ly α**



slow-to-moderate roll OK

wild rise & roll up/down OK

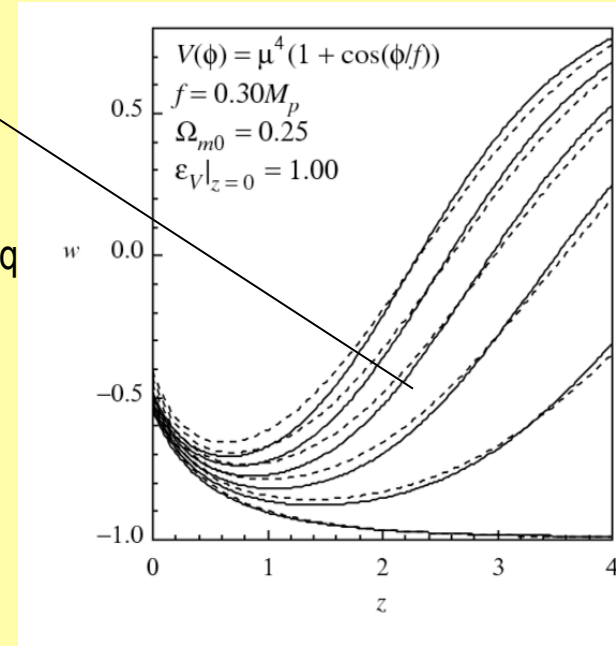
$$\epsilon_v = (\text{dln}V/\text{d}\psi)^2/4 \text{ @pivot } a_{eq}$$

$$\epsilon_s = -0.03 \pm 0.25 \text{ now}$$

$$a_s < 0.36 \text{ (} z_s > 2.3 \text{) now}$$

$$\xi_s = \text{dln}\epsilon_s / \text{dln}a \times 1/2 \text{ @pivot } a_{eq}$$

ill-determined now

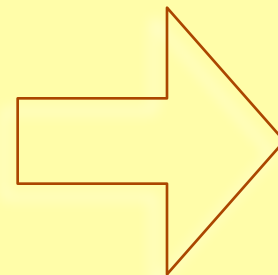


ϵ_s to ± 0.07 then

Planck1+JDEM SN+DUNE WL,

weak $a_s < 0.21$ then, ($z_s > 3.7$)

3rd param ξ_s ill-determined then



cannot reconstruct the quintessence potential, just the

slope ϵ_s & hubble drag info

(late-inflaton field < Planck mass, but sometimes not by a lot)

INFLATION NOW

WHAT IS ALLOWED?

radically variable braking in
acceleration component

$$\varepsilon(k) = (1+q)(a) = -d \ln H / d \ln a$$

Blind trajectory analysis cf. data, then & now

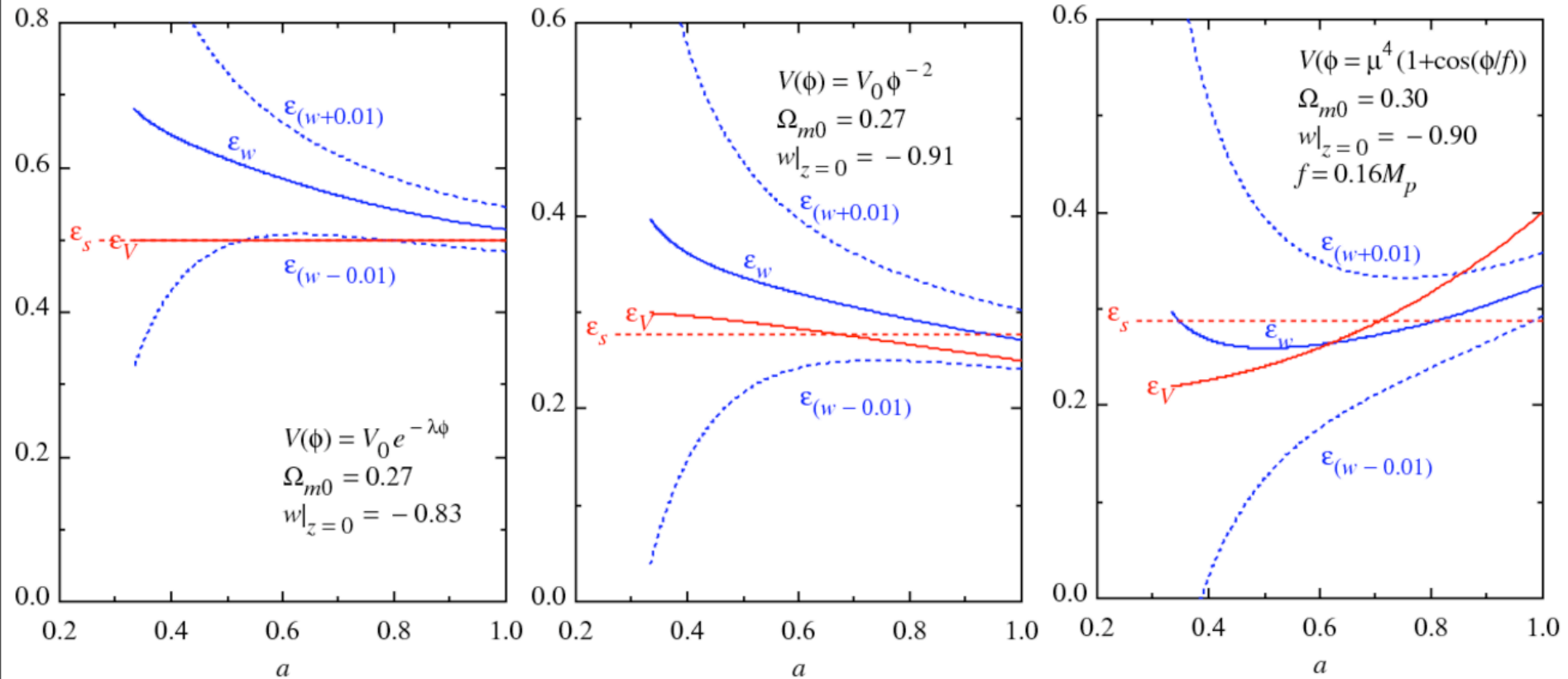
expand $\varepsilon(k)$ in localized mode functions e.g. Chebyshev/B-spline coefficients ε_b

the measures on ε_b matter - choice for “theory prior” = informed priors?

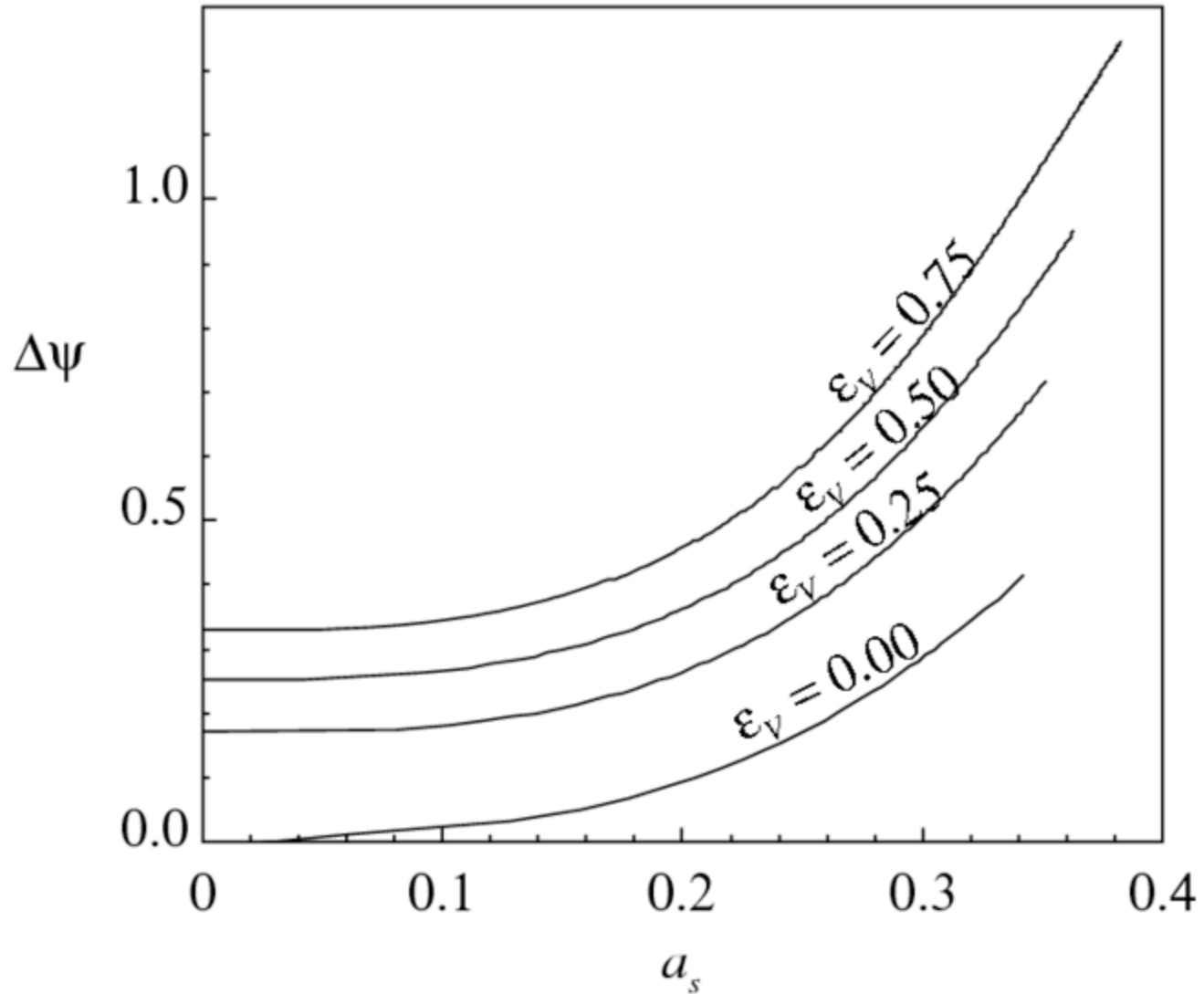
ϵ_V trajectories are slowly varying: why the fits are good

Dynamical $\epsilon_w = \epsilon_\phi \epsilon_s / \epsilon_{\phi\text{-approx}}$ cf. shape $\epsilon_V = (V'/V)^2$ (a) / (16πG)

& ϵ_s is ϵ_V uniformly averaged over $0 < z < 2$ in a



the quintessence field is below the reduced Planck mass



➤ Cosmological
Constant ($w=-1$)

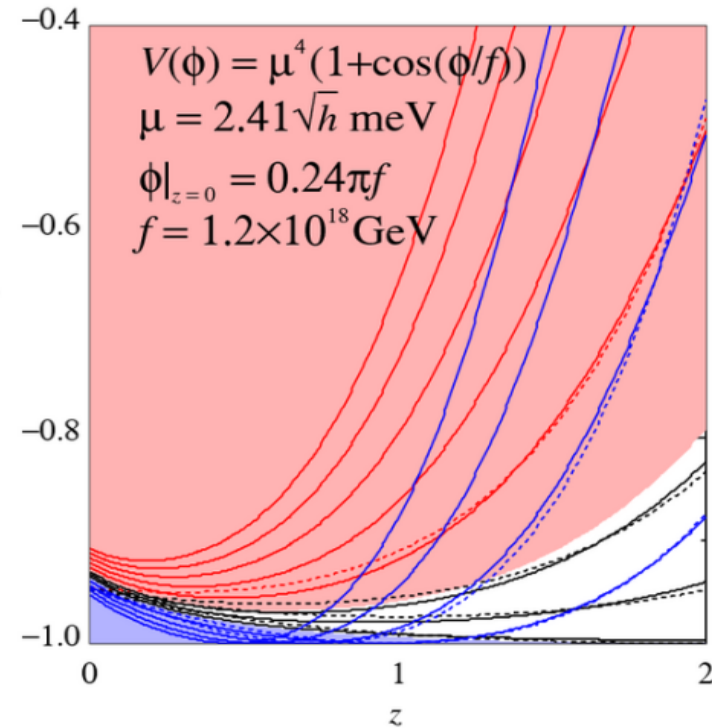
➤ Quintessence
($-1 \leq w \leq 1$)

➤ Phantom field
($w \leq -1$)

➤ Tachyon fields
($-1 \leq w \leq 0$)

➤ K-essence
(no prior on w)

INFLATION NOW PROBES NOW



trajectory probability: ~ 1 e-fold \Rightarrow blind is bad \Rightarrow slow-to-moderate roll $++$

$$-d \ln \rho_\phi / d \ln a \text{ / } 2$$

$$= \epsilon_\phi(a) = (1+w)/2$$

$$= \epsilon_s f(a/a_{\Lambda \text{eq}}; a_s/a_{\Lambda \text{eq}}; \zeta_s)$$

$$\epsilon_s = (d \ln V / d \psi)^2 / 4 \text{ @pivot } a_{\text{eq}}$$

$$\zeta_s = \pm 1.001 d^2 \ln V / d \psi^2 / 4 \text{ @pivot } a_{\text{eq}}$$

$$\zeta_s = d \ln \epsilon_s / d \ln a \times 1/2 \text{ @pivot } a_{\text{eq}}$$

INFLATION NOW

PROBES NOW

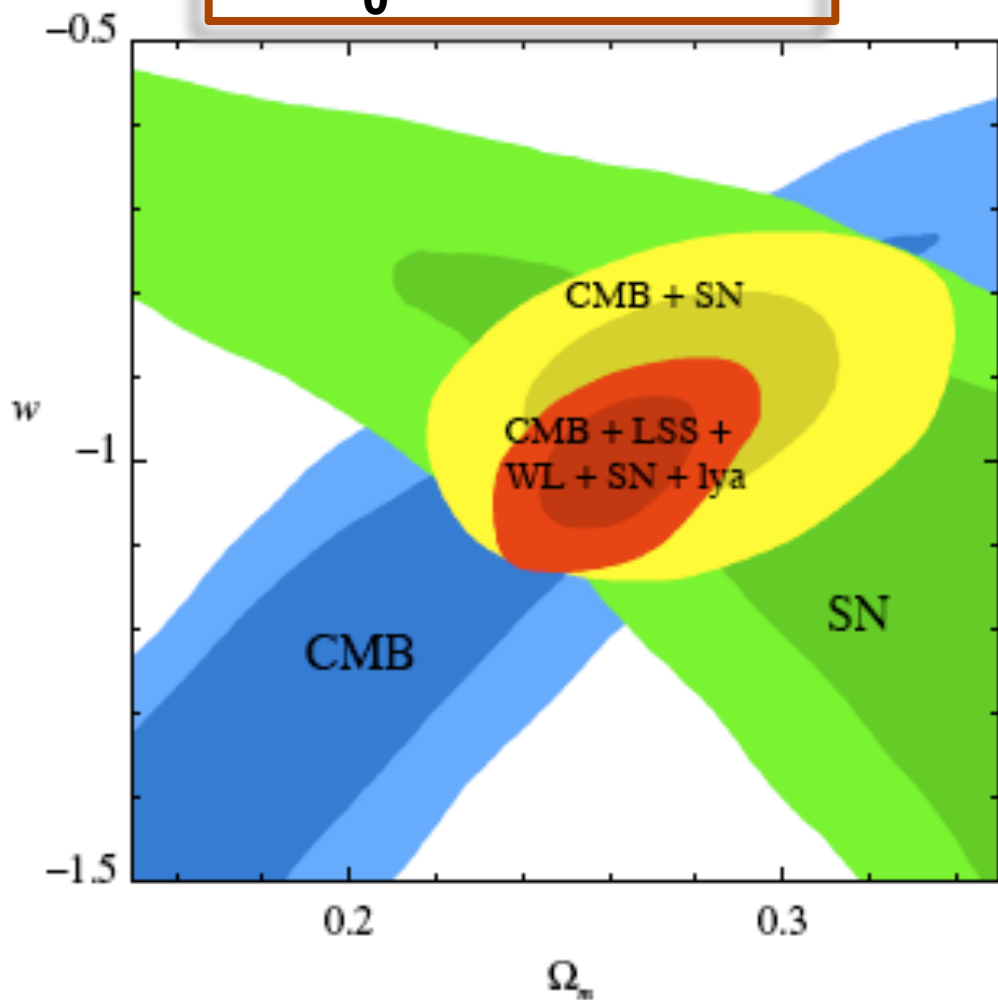
$$1 + \mathbf{w}_0 = -0.0 \pm 0.06$$

$$w(a) \equiv \frac{p(a)}{\rho(a)}$$

$$\mathbf{w}(a) = \mathbf{w}_0 + \mathbf{w}_a(1-a)$$

$$1 + \mathbf{w}_0 = -0.01 \pm 0.19$$

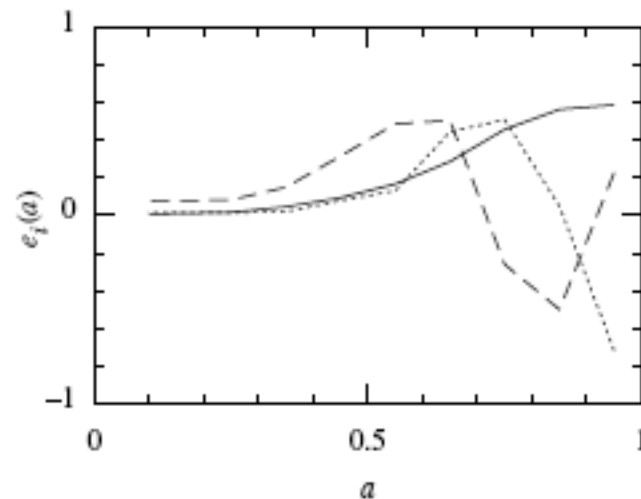
$$\mathbf{w}_a = 0.0 + 0.6 - 0.8$$



piecewise parameterization
 4,9,40 modes in redshift
 9 & 40 into Parameter eigenmodes

data cannot determine >2 EOS parameters
 DETF Albrecht et al06, Crittenden et al06, hbk08

$$\sigma_1 = 0.13 \quad \sigma_2 = 0.33 \quad \sigma_3 = 0.58$$



$\epsilon_{\phi_0} = 0.0 \pm 0.09$ if constant, $\epsilon_{\phi_0} = -0.015 \pm 0.3$ if a-linear model

Inflation Then $\epsilon(k) \sim r/16$

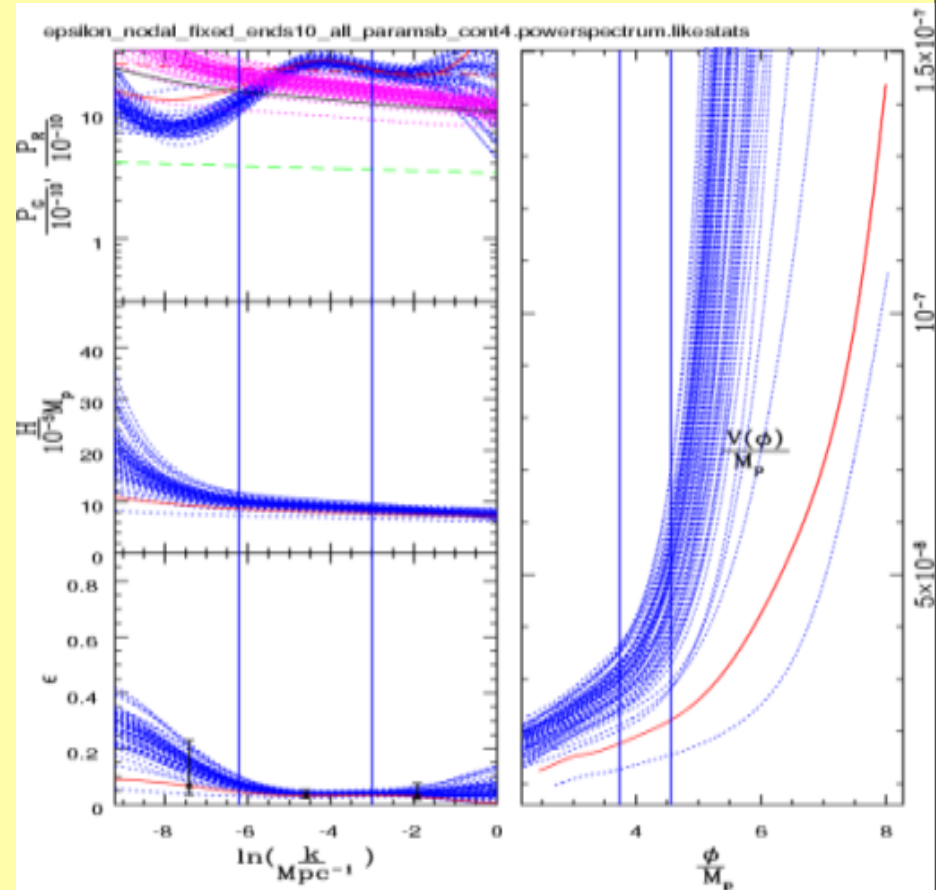
= mode expansion in $\ln H a \sim \ln k$

be blind: all $\epsilon < 1$ trajectories give allowed potential & kinetic energies

~10 good e-folds $k \sim 10^{-4} \text{Mpc}^{-1}$ to $\sim 1 \text{Mpc}^{-1}$

~ 10+ parameters? $H(\phi)$, $V(\phi)$

Bond, Contaldi, Huang, Kofman, Vaudrevange 08



Inflation Now all $\epsilon_\phi < 1$ trajectories give allowed potential & kinetic energies but... do not be blind:... ~1 good e-fold. only ~2params

get $\epsilon_s = (d \ln V / d \psi)^2 / 4$ @ pivot pt

Huang, Bond & Kofman 08