

Origin of the observed entropy in the Universe and SMpp/BSMpp particles

- coarse-grained coherent condensate breaks up into finegrained fluctuations
- particle creation = instability => stretch and break via mode-mode coupling

Dick Bond?

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Bond+Braden+Frolov+Huang+Morrison

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LOS LOSING

Calling

The Kinematics of Inflation, Preheating and Heating: a Playground for Kolmogorov-Sinai and Shannon Entropies Dick Bond @ ipmu18 11 29



what are the degrees of freedom / parameters of the ultra early Universe? TBD

begin-inflate => inflate => end-inflate => preheat => non-equilibrium heat+entropy

=> Standard Model particle physics QG plasma radiation dominated

=> dark matter dominated structure via gravitational instability => dark energy now

 $d\zeta(x,t) = (d\mathbf{E}+\mathbf{p}d\mathbf{V})/3(\mathbf{E}+\mathbf{p}\mathbf{V}) = d \ln \mathbf{\rho_c}/3(\mathbf{1}+\mathbf{w}) + \text{Trace } d\alpha'_i$

fit into a UV-complete theory (ultra-high energy to the Planck scale) strings, landscape, .. & IR-complete theory (post-inflation heating -> quark/gluon plasma)??? TBD

role of (1) instabilities after inflation

entropy generation via the breakup of the coherent low-k inflaton condensate into incoherent high-k fluctuations at a "shock-in-time" => nonGaussianity

role of (2) instabilities during inflation phenomenology of in-states propagating through localized unstable potential structures to out-states, like scattering theory => nonGaussianity

(3) |cg <=> fg> condensate/fluctuation framework, for both using coherent states classical-like approach with **hbar**. includes **Bogoliubov** transformations for fluctuations as condensate evolves => particle creation interpretation in both heating and inflating regimes.

single field V heating slow, oscillating *fine-grain fluctuations = S* generation

> *longitudinal instability* KS yes but no LSS modulation a =

A visualized 2D slice in lattice simulation



su-W-ck&NR=1 www.youtube.com/watch?v=FW

the true quadratic ζ -Websky of the ζ -scape

Planck 2018 X inflation: TTTEEE lowL Epol + CMBlens + BK15 BB + BAO

CMB TT power L~ 20-30 dip => ζ-Spectrum k-dip; includes CMB lensing, parameter marginalization



inflaton V(ϕ)-maps =3MP² H² (1- $\epsilon/3$) HJ eqn, d ϕ /MP/d/na=±sqrt(2 ϵ) along the gradient / Morse flow



Inflaton During and After Inflation Bond@ipmu18 11 29



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Quantum Inflation - if quantum energy then quantum gravity (entangled) then gravitons Phonons *density fluctuations = Trace strain = spatial 3-volume fluctuations*

=> combined entropy-like measure ζ =inflaton



Gravitons tensor perturbations transverse traceless strain $P_{GW} = r P_{\zeta}$ grail r < .06 now, to < .001Isocons when multiple particle-species - orthogonal scalar degrees of freedom to inflaton/phonon Dilatons 4-volume fluctuations - Higgs inflation $L_G(R)$ gravity - conformally-flatten potentials moduli, axions connection to particle physics models "fundamental scalars" ... string theory fermions, vector gauge fields, Higgs Standard model of particle physics . vector perturbations

$$\frac{\mathrm{d}\zeta}{\mathrm{d}t} = \sum_{i} \frac{\nabla \dot{\phi}_{i} \cdot \nabla \phi_{i} + \dot{\phi}_{i} \nabla^{2} \phi_{i}}{3a^{2}(\rho + P)}$$

nonlinear multi-field classical coupled system. evolve using lattice simulations. via **symplectic defrost++ code + spectral code** => high accuracy to unveil small effects after and during inflation

CMB/LSS observations give access to limited partial operator-paths in field-space aka kinematics => *limited glimpse of a UV/IR complete theory aka dynamics*. so far just through the ζ -fluctuation spectrum encoding the quantum diffusion of (longitudinal/inflaton) paths. *Transverse field-motion aka isocons may influence* $\mathcal{P}_{\zeta\zeta}(k)$ and ζ -nonG stochastic dynamical systems theory for kinematics of cg variables

coarse-grained stochastic inflation k~Ha resolution/dynamics relation $d\zeta(x,Ha) = [\mathcal{P}_{\zeta\zeta}(k)]^{1/2} [dlnHa]^{1/2}\eta_{GRD}(x,Ha)$ quantum fluctuations & no drift order parameters from the ultra early U? so far $\zeta \approx GRF$, $\mathcal{P}_{\zeta\zeta}(k)$ 2 ζ -params

 ζ = an adiabatic invariant, conserved even with large field-strains to get nonG need to break adiabaticity => nonlinear fluctuations

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stochastic dynamical systems theory for kinematics of cg variables

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Kinematics of Preheating:

BBFH are developing a trajectory bundle evolution framework for nonG in post-inflation preheating -> heating. field-space path deviations aka strains and their shear can characterize smoothed-bundle "ballistic" chaotic evolutions. arrested by nonlinear fluctuation-mode generation. metric entropy (KS) in the ballistic phase -> Shannon entropy in the fluctuation-mode phase. condensate evolves chaotically? => phase-transition-like melting into phonon+ modes, fast scramble of fundamental fields, evolution of statistically-simple energy-density phonons can one use coherent states to address this quantum mechanically? TBD

 $d\zeta_{fi}(x,t) = dln \rho_c / 3(1+w) + Trace (d\alpha)$

cf. ds_{fi} = ds_{Ksfi} -dTr \mathcal{E}_{fi} = δ Sfi (fg-> cg) FokkerPlanck eqn for Prob(δ Sfi| control) $\zeta_{NL}(x)$ from "isocon" degrees of freedom cf. $\zeta_{NL}(x)$ from inflaton

modulated heating, ballistic chaos, caustics, shock-in-time,

modulators isocon $\chi(x)$, axionic-isocon(x) couplings g(x) super-horizon accessible complex field configs Oscillons, Stringy configurations, curvatons, ... Bubble Collisions



around a minimum is the HOT /heating question 2 filament?

how was matter & entropy generated at the end of acceleration = inflation?



quartic inflaton V(ϕ, χ) = 1/4 $\lambda \phi^4$ + 1/2 g² $\phi^2 \chi^2$



log-normal pdf (density aka ζ), in k-bands too; normal pdf (velocity)





nonG from large-scale modulations of the shock-in-times of preheating



 $\Delta S_{\phi+\chi} \sim 1/2 \sum_{k} ln \det \langle (\phi, \Pi_{\phi}, \chi, \Pi_{\chi}) (\phi, \Pi_{\phi}, \chi, \Pi_{\chi})^{\dagger} \rangle$

A Shocking End to Post Inflation Mean Field Dynamics

Shock-in-space t = const $V_{bulk}^2 > C_s^2 \Rightarrow V_{bulk}^2 < C_s^2$

supersonic \Rightarrow subsonic

Characteristic spatial scale Jump Conditions: $\Delta T^{\mu\nu}$ **Randomizing** Shock Front: ΔS **Mediation**: width via viscosity or collisionless dynamics **post-shock evolution,** slow, of temperature, etc. coherent condensate breakup? Shock-in-time x = const (deviations for nonG) $<\rho>>\delta\rho \Rightarrow <\rho> <<\delta\rho$

Homogeneous \Rightarrow Fluctuations

Characteristic temporal scale Jump Conditions: $\Delta T^{\mu 0}$ **Randomizing** mode cascade & Particle Production: ΔS **Mediation**: width via gradients and nonlinearities

post-shock evolution, slow, of fluctuations
coherent condensate breakup E-phonon soup
a difference: chaotic lead in to shock



dS/dt(t,g) => the Shock-in-time: entropy production rate **Shock**(χcg,eoi(x) |g²/λ)) => Chaotic Billiards: NonG from Parametric Resonance in Preheating B+Frolov, Huang, Kofman 09

B+Braden, Frolov, Huang 19



 $V(\phi,\chi) = 1/4 \lambda \phi^4 + 1/2 g^2 \phi^2 \chi^2$

ζ conserved along trajectories until the "shock-in-time" when high **k** fluctuations (fine-grain) develop from coarse-grain, measure is $s_{Shannon} = -\ln \mathcal{P} \sqrt{g}$

but -Dln \mathcal{P} /dt = Trace d \mathcal{E} /dt does change ~ KS entropy (rate)

stretching of phase strings. begin with anisotropic Gaussian at EoI and watch it stretch, \mathcal{E} grows, rotates, locally OK as distorted ellipsoid, but strain depends upon the central value => phase tubes



caustics in <q^> ballistic **orbits**

 $\langle \delta q^{A} Xt2 | \delta q^{B} Xti \rangle \sim \exp(\mathcal{E}(Xt2 | Xt1))^{A_{C}} \langle \delta q^{C} Xt1 | \delta q^{B} Xti \rangle$ early U parameters: final φ , Π_{φ} , χ , Π_{χ} , In a, In ρ , initial $\chi_{cg,eoi}$, couplings g, λ , ... parameter strain tensor in field space $\mathcal{E}^{A_{C}}(Xt2 | Xt1)$ deformation $e^{\mathcal{E}}$

 $d\mathcal{E}^{A_{C}}/dt$ strain rate ~local Lyapunov coefficients Floquet instability charts instability to have nearby parameters diverge => chaotic billiards Kolmogorov-Sinai entropy: ~ Sum of positive evalues of $d\mathcal{E}/dt$

small $\mathbb{E}^{A_{A}}$ eigenvalues=> coherent trajectory bundles (for a time) = caustics (inverse -> ∞) 1/ [$\partial \alpha$ / $\partial \chi_{cg,eoi}$]; => peaks in ζ ($\chi_{cg,eoi}$) stopping time **tstop** ($\chi_{cg,eoi}$) when $\mathbb{E}^{A_{A}}$ evalues get large <=> local gradients \uparrow

cf. LargeScaleStructure: final Eulerian position <= initial Lagrangian position 1LPT aka Zeldovich: $\partial x/\partial r = \exp(\mathcal{E}) \rightarrow 0$ density $\rho \sim \exp(-Tr(\mathcal{E})) \rightarrow \infty$



ballistic billiards k=0 mode phase space string evolution

2D constrained distribution functions

stopping criterion when coarse-grained entropy of field variables rises $\langle = \rangle$ strain \mathcal{E} high, *ie* when integral of the Kolomgorov-Sinai entropy reaches a threshold - very $\chi_{cg,eoi}$ dependent





V= 1/4 $\lambda \phi^4$ + 1/2 g² $\phi^2 \chi^2$

phase space strings

2D constrained distribution functions



phase string growth in time "parameter strain" integral of Kolmogorov-Sinai entropy rate





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Bond , Huang, Stein; Braden, Morrison, ...

Primordial Non-Gaussianity in CO; large scale => CHIME much larger volume is better



role of (2) instabilities during inflation phenomenology of in-states propagating through localized unstable potential structures to out-states, like scattering theory => nonGaussianity

role of "classical evolution/instabilities" in "particle creation" (1) in heating 🖌 (2) during inflation

Bond+Braden+Frolov+Morrison - work very much in experimental/exploratory phase

map ballistic extreme-parameter-strain view of (pre-) heating into inflating regime i.e., role of classical instabilities during inflation, need to start with quantum fluctuations but then classical "particle creation" what about particles? tie to entropy old & new way: $n_{Ak} \sim \rho_{Ak} / \hbar \omega_{Ak}(t) \sim \exp(\Delta s_{Ak})$

quantum fluctuations are there, but *only as seeds for instabilities* in which trajectories diverge: *little probability Gaussian blobs stretch into highly deformed elongated surfaces. "phase strings" in 2D, 3D*

history (for me). sbb87-89 $\langle \delta \mathscr{P}_{\phi^A \phi^A}(k) | \delta V, \delta m_{eff}^2 \rangle$, $\langle \delta \mathscr{P}_{\zeta\zeta}(k) | \delta V(\phi, \chi) \text{ controls} \rangle$ multifield hybrid, mountains/valleys of extra power. role in non-Gaussianity. role of Higgs etal.

tool: full linearized k evolution, from inside to outside

build on vilenkin/starobinsky stochastic inflation: sb90, 91

used the attractor, aka reduced Hamilton principal function, Sreduced ~ -2Mp² H gives π_A

B91 more general Langevin network leads to P(ϕ A, π _A,H| α) or better P(ϕ A, π _A,H|Ha)

=> nice expression in terms of quantum diffusion velocity/current

today:

dilemma 1: missing terms in crossing horizon (coarse grain / fine grain flow) dilemma 2: outside horizon is a condensate. how to express this dilemma 3: phase coherence as in classical realizations are useful

- incorporate via coherent state condensates and fluctuations upon them?? instabilities with quantum fluctuations the tiny seeds





 $\log \left[\mathcal{P}_{\phi}(k)/H^{2} \right]$















experiment χ -light



experiment χ -heavy





experiment χ -light



experiment χ -light



occupation numbers & particle creation ~ "Gaussian entropy" in the single A-field

$$\ln(n_{\phi^A k} +) \sim \Delta s_{\phi^A k} \sim \frac{1}{2} \operatorname{Trace} \ln[C_{\phi^A \phi^A} C_{\Pi_A \Pi_A} - C_{\phi^A \Pi_A}^2]$$

occupation ~ $n_{Ak} \sim e^{\Delta s_{Ak}}$

old way if well defined mode energies $\omega_{Ak}(t) = \ln(n_{Ak} + 1/2) \sim \ln[\rho_A/\hbar\omega_{Ak}]$

full "Gaussian entropy" in the 2 fields, C are k-mode correlations = power spectra - generalized Sackur-Tetrode

$$\Delta s_{A+B,k} = \frac{1}{2} \operatorname{Trace} \ln [C_{\phi^A \phi^B} C_{\Pi^A \Pi^B} - C_{\phi^A \Pi^B} C_{\Pi^A \phi^B}]$$





phonon occupation





 $\zeta \; \mathsf{PDF} \; (\Delta V{=}0)$



 $\zeta \, \mathsf{PDF}$



experiment χ -heavy

unstable χ cf. trapped



experiment χ -heavy unstable χ cf. trapped



experiment χ-heavy

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unstable χ



trapped χ



unstable χ



trapped χ



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trapped χ



unstable χ



(3) |cg <=> fg> condensate/fluctuation framework, for both using coherent states classical-like approach with hbar. includes Bogoliubov transformations for fluctuations as condensate evolves => particle creation interpretation in both heating and inflating regimes. **condensate/fluctuation framework**, classical-like approach with hbar + Bogoliubov transformations for fluctuations as condensate evolves => particle creation & fluctuation freeze-out into new condensate

stochastic inflation | q_c ; q_f > instead of $dq_c^A = V_c^A dT + K^A_v \sqrt{dT} \eta^{v}(GRD)$

Langevin network evolution step: $q_c(X,T+dT) = q_c(X,T) + V_c dT + \delta q_f$

cf. condensate evolution step $|q_c(T+dT)\rangle = exp(V_cdT) exp(\delta q_f) |q_c(T)\rangle$ schematic $\delta q_f(x,T) = \sum_{k-band} (Q_k^*(x,T)a_k^T - Q_k(x,T)a_k)$ $V_cdT = V_c(x,T)dT(a_x^T - a_x)$

annihilation/creation operators in position and momentum $a_x, a_x^T a_k, a_k^T$

fluctuating part $|\mathbf{q}_{f}\rangle \sim \exp(\sum \delta \mathbf{q}_{f}) | \mathbf{0} \rangle a$ coherent state description? what is the relation to the usual $\mathbf{q}_{f,op} = \sum_{k} (\mathbf{Q}_{k} \ast (\mathbf{x},T) \mathbf{a}_{k}^{T} + \mathbf{Q}_{k} (\mathbf{x},T) \mathbf{a}_{k})$ operator linear in Bunch Davies vacuum operators $\mathbf{a}_{k}, \mathbf{a}_{k}^{T}$ (sign difference)

it is an overcomplete basis representation, but it conforms to a classical lattice simulation of inflation (no bipartite split). still use the gradient expansion for $|q_c(T)>$ & mixed operators $V_c dT$ and $\delta q_f(x,T)$ - promising approach?

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