

The Dark Energy trajectories in the post-EUCLID era

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Outline

Introduction

Parameterizing Dark Energy: beyond the phenomenological w_0-w_a

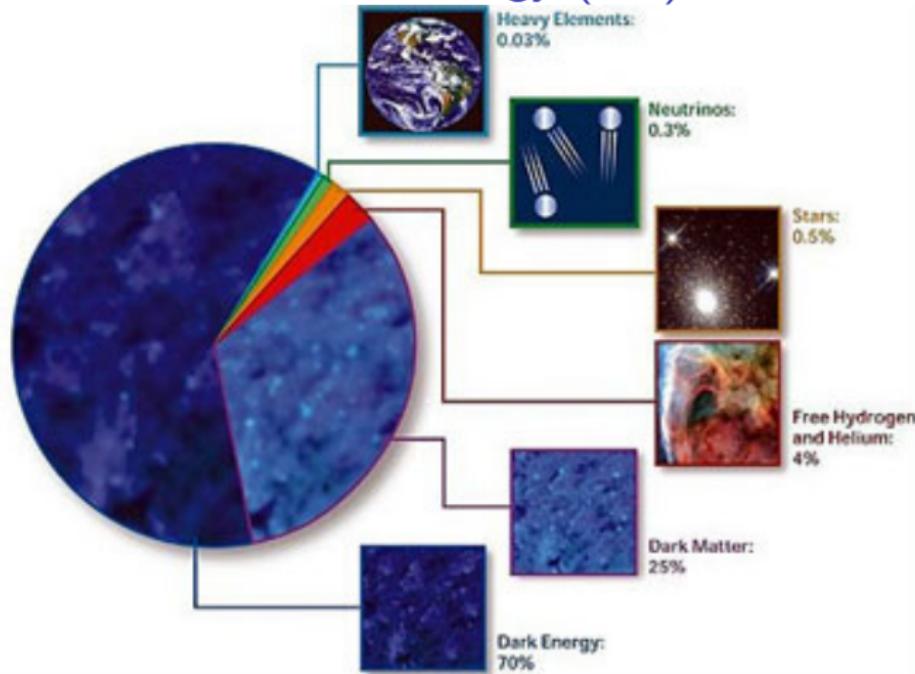
Observational Constraints

current data: SNLS3yr and the latest H_0 measurement

The future observations: Planck CMB + EUCLID LSS (+
21cm + SN)

Conclusion

Cosmic Pie and Dark Energy (DE)



NASA, http://en.wikipedia.org/wiki/File:Cosmological_composition.jpg

Cosmological Constant

$$S = \int \sqrt{-g} d^4x \left[\frac{M_p^2}{2} (R - 2\Lambda) + \mathcal{L}_{\text{matter}}(g^{\mu\nu}, \psi_m) \right]$$
$$\rho_{\text{DE}} = \text{constant}$$

The equation of state (EOS):

$$w_{\text{DE}} \equiv \frac{P_{\text{DE}}}{\rho_{\text{DE}}} = -1$$

Parameterizing the DE EOS

If not $w_{\text{DE}} = -1$, how to parameterize DE?

The oft-used options: constant $w_{\text{DE}} = w_0$ or linear

$$w_{\text{DE}} = w_0 + w_a(1 - a).$$

- ▶ Too many DE models \Rightarrow difficult to do a model-by-model selection.
- ▶ These are good low-redshift approximations for a variety of models.

The take-home message of this talk:

For a wide class of DE models, we have a better choice.

Dark Energy Candidates

A general framework with an extra scalar d.o.f ϕ (scalar-tensor theory)

$$S = \int \sqrt{-g} d^4x \left[\frac{M_p^2}{2} A(\phi; R) + \frac{1}{2} B(\phi) \partial^\mu \phi \partial_\mu \phi - V(\phi) + \mathcal{L}_{\text{matter}}(g^{\mu\nu} e^{2\alpha(\phi)}, \psi_m) \right]$$

Only two free functions are physical:

- ▶ Jordan Frame : $\alpha(\phi) = 0, B(\phi) = \pm 1$.
- ▶ Einstein Frame : $A(\phi) = R, B(\phi) = \pm 1$.

Examples: quintessence; phantom; Brans-Dick theory; $f(R)$ gravity; ...

Dark Energy Candidates

If the coupling between ϕ and matter is negligible:

$$S = \int \sqrt{-g} d^4x \left[\frac{M_p^2}{2} R \pm \frac{1}{2} \partial^\mu \phi \partial_\mu \phi - V(\phi) + \mathcal{L}_{\text{matter}}(g^{\mu\nu}, \psi_m) \right].$$

- +: quintessence.
- : phantom.

- ▶ These are the simplest alternatives to Λ .
- ▶ Many concrete models: Ratra & Peebles 1988; Wetterich 1988; Frieman et al. 1995; Binetruy 1999; Barreiro et al. 2000; Brax & Martin 1999; Copeland et al. 2000; de La Macorra & Stephan-Otto 2001; Carroll et al. 2003; Caldwell 2002; Caldwell et al. 2003; Linder 2006; Copeland et al. 2006; Huterer & Peiris 2007; Linder 2008 ...

Quintessence: two classes of models

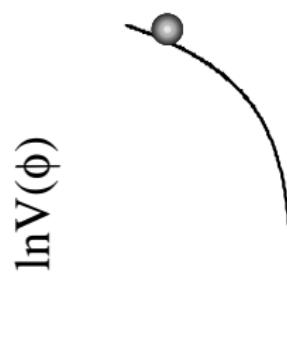
tracking:



ϕ

early-universe: fast-roll
late-universe: slow-roll

thawing:



ϕ

early-universe: frozen
late-universe: slow-roll

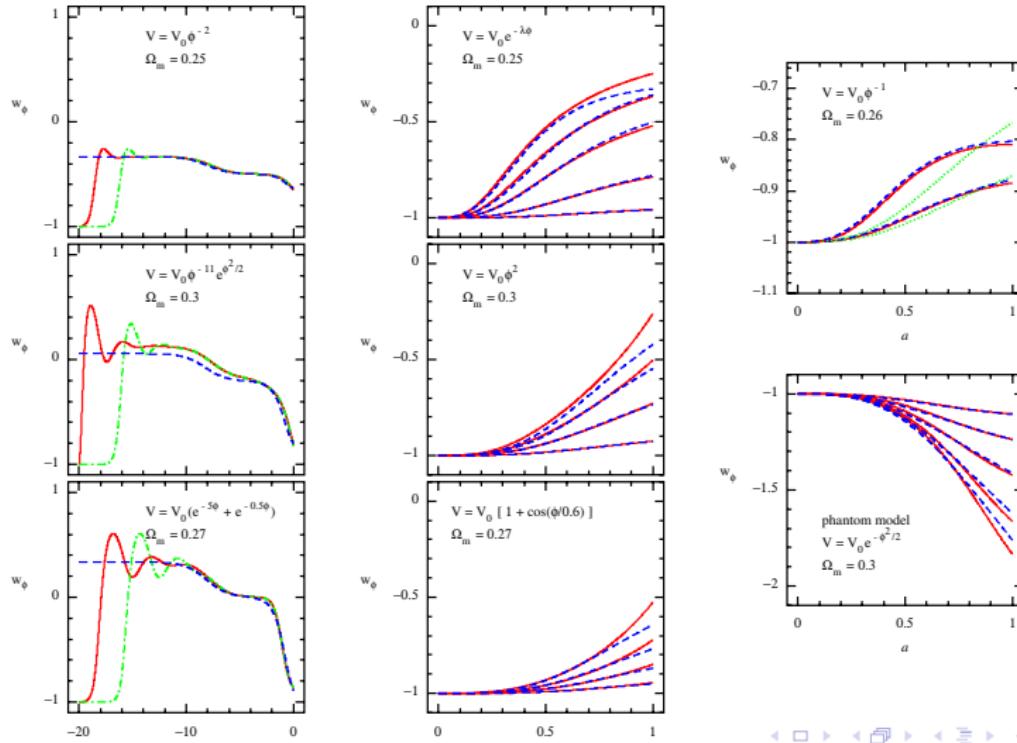
The quintessence/phantom w_{DE} recipe

$$w_{\text{DE}} = w_\phi = F(a; \Omega_m, \epsilon_s, \epsilon_{\phi,\infty}, \zeta_s)$$

- ▶ The slope parameter $\epsilon_s \equiv \pm \frac{M_p^2}{2} \left(\frac{d \ln V}{d\phi} \right)^2$ at low redshift ('+' for quintessence, '-' for phantom).
- ▶ The tracking parameter $\epsilon_{\phi,\infty} \sim |1 + w_\phi|$ at high redshift.
- ▶ The running parameter ζ_s is related to $|d\phi/dt|$ and $\frac{d^2 \ln V}{d\phi^2}$ at low redshift (for thawing models $\zeta_s \propto \frac{d^2 \ln V}{d\phi^2}$).

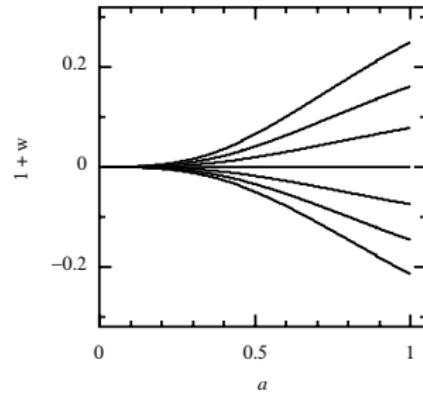
For the explicit expression of F and more details see *Huang, Bond, Kofman, 2011 (ApJ)*.

Comparing our w_{DE} formula with the exact solutions.



The thawing models with slow rolling.

Slow-roll thawing models: only ϵ_s is relevant.



from bottom to top:

$$\epsilon_s = -0.75, -0.5, \dots, 0.75$$

The degeneracy between w_0 and w_a (defined as $\frac{dw}{da}|_{a=1}$).

$$1 + w_0 + w_a \left(0.264 + \frac{0.132}{\Omega_{m0}} \right) = 0 .$$

Cosmological Data Sets

- ▶ Cosmic Microwave Background (**CMB**): WMAP7(2010), ACT(2010), Acbar (2009), QUAD (2009), BICEP (2009), CBI (2008), Boomerang (2006), VSA (2004), MAXIMA (2000)
- ▶ Type Ia Supernova (**SN**): 472 SNs (123 low-z + **242 SNLS3yr** + 93 SDSS1yr + 14 HST)
- ▶ Weak Lensing (**WL**): COSMOS + CFHTLS-wide + RCS + VIRMOS + GaBoDS.
- ▶ Large Scale Structure (**LSS**): SDSS-DR7 LRG (2009).
- ▶ Ly α Forest (**Ly α**): SDSS (P. McDonal 2005, 2006).
- ▶ **HST constraint** $H_0 = 73.8 \pm 2.4 \text{km s}^{-1} \text{Mpc}^{-1}$. (Riess et al 2011)

SNLS3yr

Table 2. Cosmological results assuming a flat universe and constant w for the SNLS3 sample plus BAO and WMAP7

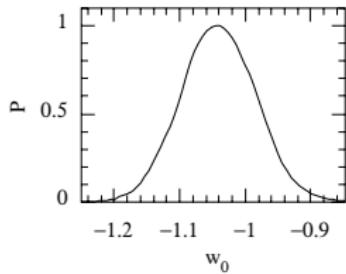
| Fit | α^a | β^a | M_B^1 | M_B^2 | Ω_m | w |
|----------------------|---------------------------|---------------------------|---------|---------|---------------------------|----------------------------|
| Marginalization fits | | | | | | |
| Stat only | $1.450^{+0.112}_{-0.105}$ | $3.164^{+0.096}_{-0.094}$ | -19.164 | -19.227 | $0.276^{+0.016}_{-0.013}$ | $-1.043^{+0.054}_{-0.055}$ |
| Stat + sys | $1.367^{+0.086}_{-0.084}$ | $3.179^{+0.101}_{-0.099}$ | -19.175 | -19.220 | $0.274^{+0.019}_{-0.015}$ | $-1.068^{+0.080}_{-0.082}$ |

2 3 4 5 6 7 8 9 10

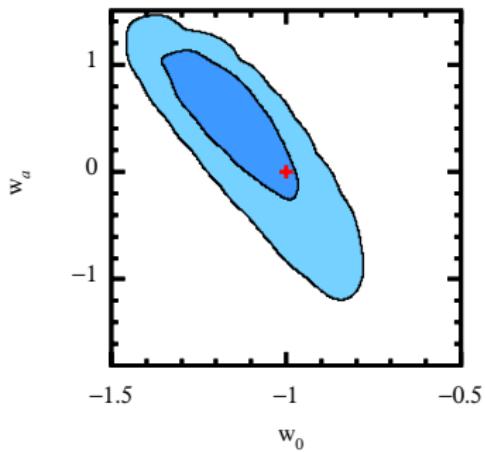
The updated phenomenological w_{DE}

constant w_{DE} :

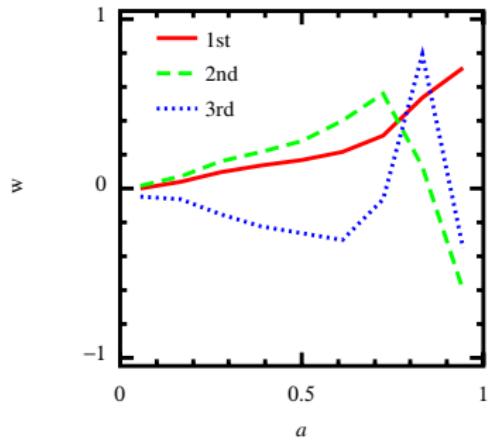
$$w = -1.040^{+0.056}_{-0.056}$$



$w_0 - w_a$:



The eigen modes



9 uniform bins $a \in [0, 1]$.

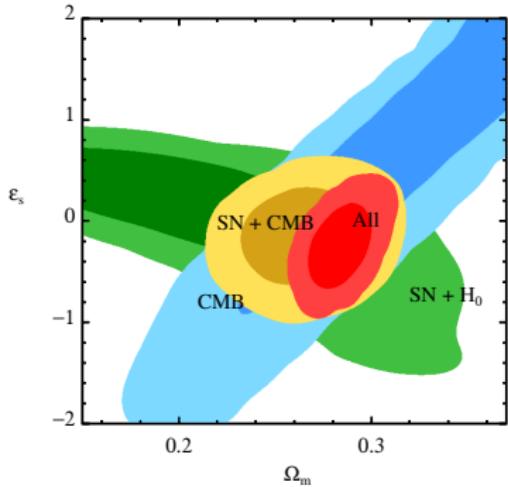
basis: top-hat functions.

$$\sigma_1 = 0.12$$

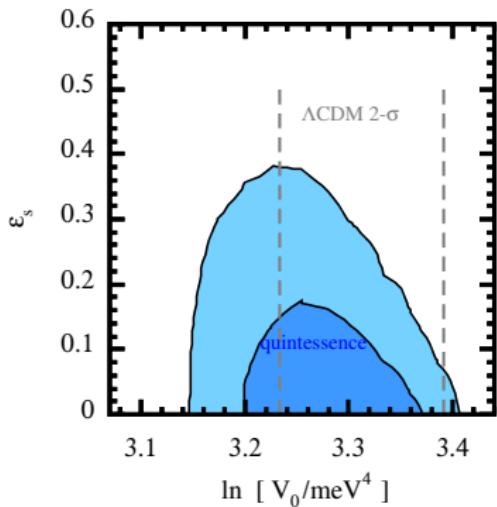
$$\sigma_2 = 0.22$$

$$\sigma_3 = 0.41$$

ϵ_s and $\ln V_0$.

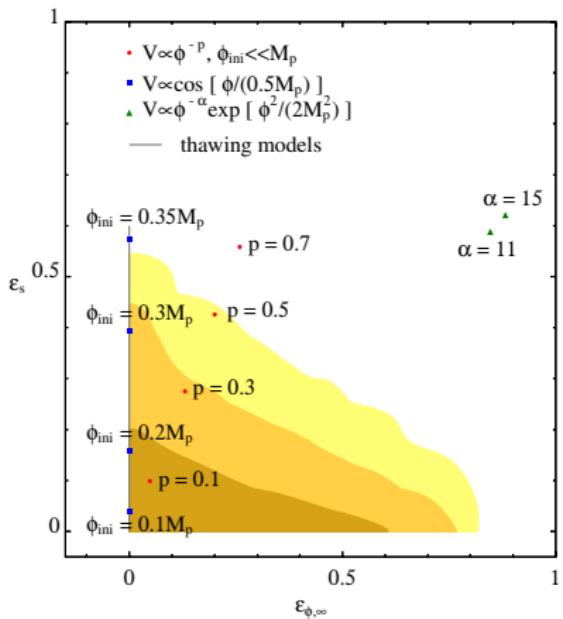


Slow-roll thawing case; assuming $\epsilon_{\phi,\infty} = 0$ and $\zeta_s = 0$.

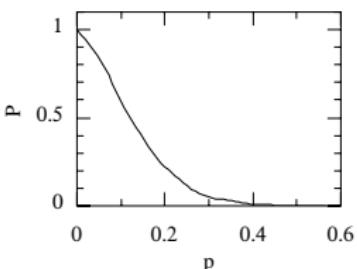


General quintessence; Uniform priors
 on $\Omega_b h^2$, $\Omega_c h^2$ and θ .

Quintessence models on the ϵ_s - $\epsilon_{\phi,\infty}$ plane.

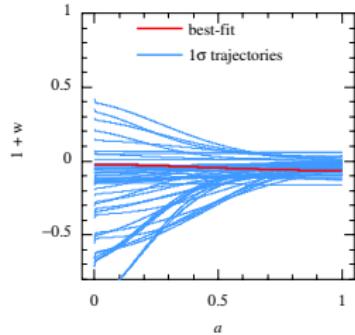


Direct constraint
on the $V \propto \phi^{-p}$
tracking model
is $p < 0.43$ ($3-\sigma$).

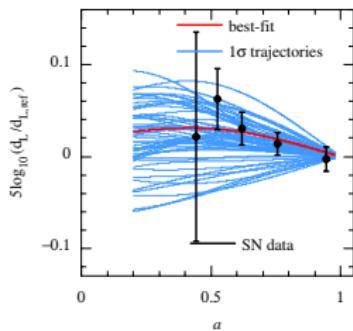


Marginalized over ζ_s and other cosmological parameters.

reconstructed $1+w_{\text{DE}}$ trajectories:



the distance moduli:



Forecasts Mock Data

- ▶ **CMB: Planck2.5yr**, using 3 channels (70GHz, 100GHz, 143GHz), assuming 5% foreground residual (synchrotron + dust), $f_{\text{sky}} = 3/4$, $l_{\text{max}} = 2500$.
- ▶ **LSS: EUCLID** spectroscopic redshift survey; $f_{\text{sky}} = 0.5$, $0.5 < z < 2.1$.
- ▶ **SN: JDEM**, 500 LOWZ ($z < 0.03$) + 2500 HIGHZ ($0.03 < z < 1.7$)
- ▶ **21-cm** survey **CHIME** 200m×200m cylinder radio telescope, 4000 receivers integrated 4 yrs; $f_{\text{sky}} = 0.36$.

The EUCLID LSS survey



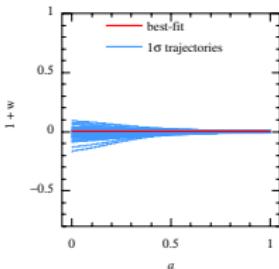
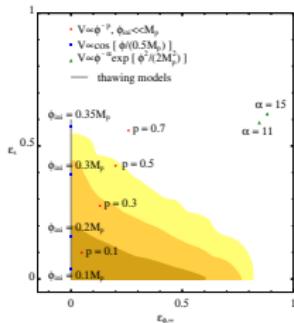
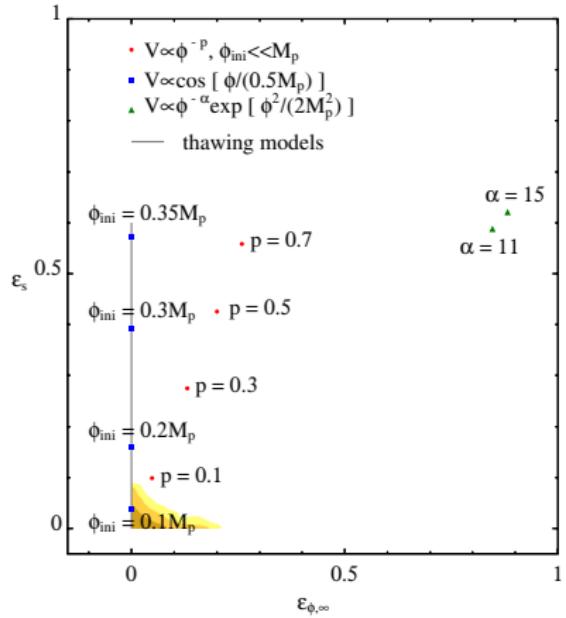
$$P_g(k, z, \mu) = \left(b + \frac{d \ln D}{d \ln a} \mu^2 \right)^2 D^2(z) P_m(k, z=0) e^{-k^2 \mu^2 \sigma^2} .$$

8 redshift bins \times 30 k bins \times 20 μ bins

marginalize over 16 nuisance parameters: $b_1, b_2, \dots, b_8; \sigma_1, \sigma_2, \dots, \sigma_8$.

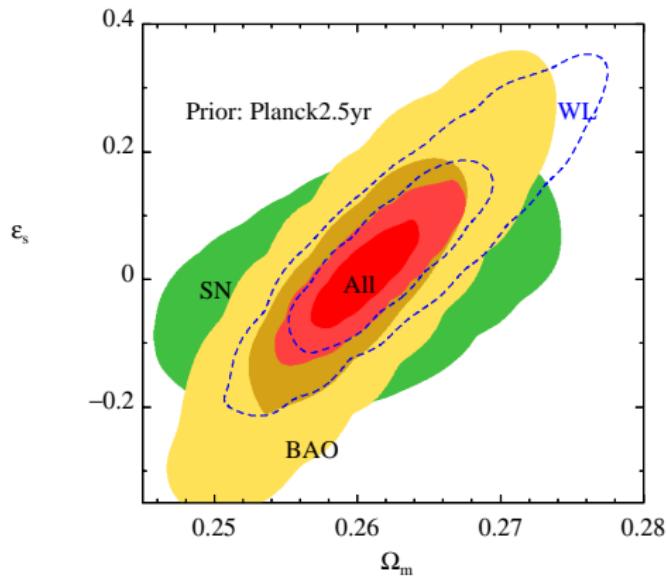
cut-off at quasi nonlinear scales ($k \sim 0.2 \text{ Mpc}^{-1}$).

DE EOS in the post-EUCLID era



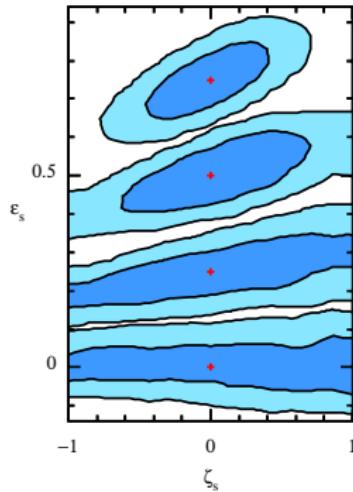
Marginalized over ζ_s and other cosmological parameters.

Comparison between different probes



- ▶ BAO can be achieved with less expensive experiments such as CHIME.
- ▶ LSS and WL are correlated; Simultaneous treatment requires extra work.

How about the “running” parameter ζ_s ?



A slowly rolling field does not “feel”
the curvature of the potential.

Conclusion

- ▶ Both quintessence and phantom models are consistent with current observations. The best-fit model is in the proximity of Λ .
- ▶ The constraints on the slope parameter ϵ_s and tracking parameter $\epsilon_{\phi,\infty}$ can be improved by a factor of about 5 by the future observations.
- ▶ The running parameter (in thawing case, the second derivative of $\ln V$ at low redshift) is not measured today, and will not be measurable in the near-future observations, unless the true model significantly deviates from Λ .
- ▶ The $\epsilon_s-\epsilon_{\phi,\infty}$ space is complementary to the standard w_0-w_a space.