Cosmic rays in clusters of galaxies – Tuning in to the non-thermal Universe

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in collaboration with

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Mar 12, 2008 / Astro Lunch Talk, University of Waterloo





Outline

- Cosmic rays in galaxy clusters
 - Introduction and motivation
 - Cluster simulations and cosmic ray physics
 - Cosmic ray pressure feedback
- Particle acceleration processes
 - Diffusive shock acceleration
 - Stochastic acceleration
 - Particle reactions
- Non-thermal cluster emission
 - Radiative processes
 - Unified model of radio halos and relics
 - High-energy gamma-ray emission





Outline

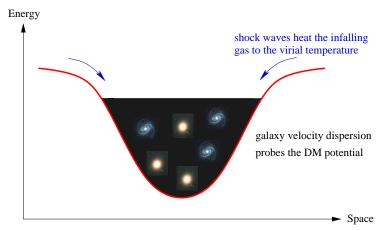
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A theorist's perspective of a galaxy cluster

Galaxy clusters are dynamically evolving dark matter potential wells:







Introduction and motivation

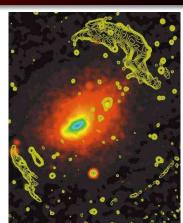
Cluster simulations and cosmic ray physic Cosmic ray pressure feedback

... and how the observer's Universe looks like



1E 0657-56 ("Bullet cluster")

(X-ray: NASA/CXC/CfA/M.Markevitch et al.; Optical: NASA/STScl; Magellan/U.Arizona/D.Clowe et al.; Lensing: NASA/STScl; ESO WFI; Magellan/U.Arizona/D.Clowe et al.)



Abell 3667

(radio: Johnston-Hollitt. X-ray: ROSAT/PSPC.)





Why should we care about cosmic rays in clusters?

It allows us to explore complementary windows to cluster cosmology

- Is high-precision cosmology possible using clusters?
 - Non-equilibrium processes such as cosmic ray pressure and turbulence possibly modify thermal X-ray emission and Sunyaev-Zel'dovich effect.
 - Non-thermal cluster emission will enable constructing a 'gold sample' for cosmology using orthogonal information on the dynamical cluster activity.
- What can we learn from non-thermal cluster emission?
 - Understanding mechanism of diffuse radio and non-thermal X-ray emission of clusters.
 - Estimating the cosmic ray pressure contribution.
 - Fundamental physics: diffusive shock acceleration, large scale magnetic fields, and turbulence.





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Literature for the talk

- Pfrommer, 2008, MNRAS, in print, ArXiv:0707.1693, Simulating cosmic rays in clusters of galaxies – III. Non-thermal scaling relations and comparison to observations
- Pfrommer, Enßlin, Springel, 2008, MNRAS, in print, ArXiv:0707.1707, Simulating cosmic rays in clusters of galaxies II. A unified scheme for radio halos and relics with predictions of the γ-ray emission
- Pfrommer, Enßlin, Springel, Jubelgas, and Dolag, 2007, MNRAS, 378, 385, Simulating cosmic rays in clusters of galaxies – I. Effects on the Sunyaev-Zel'dovich effect and the X-ray emission
- Pfrommer, Springel, Enßlin, Jubelgas 2006, MNRAS, 367, 113,
 Detecting shock waves in cosmological smoothed particle hydrodynamics simulations
- Enßlin, Pfrommer, Springel, and Jubelgas, 2007, A&A, 473, 41,
 Cosmic ray physics in calculations of cosmological structure formation
- Jubelgas, Springel, Enßlin, and Pfrommer, A&A, in print, astro-ph/0603485,
 Cosmic ray feedback in hydrodynamical simulations of galaxy formation





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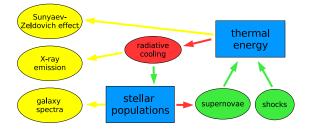
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Radiative simulations – flowchart

Cluster observables:



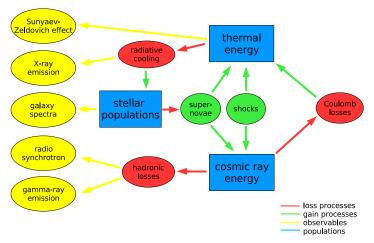






Radiative simulations with cosmic ray (CR) physics

Cluster observables:

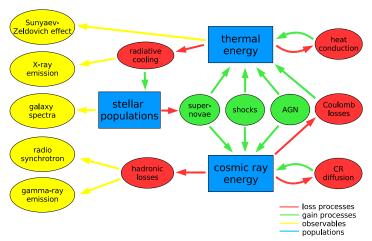






Radiative simulations with extended CR physics

Cluster observables:



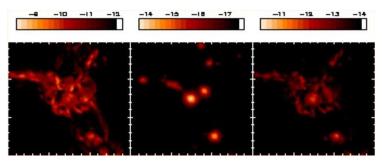


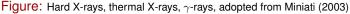


Previous numerical work on cosmic rays in clusters

COSMOCR: A numerical code for cosmic ray studies in computational cosmology (Miniati, 2001):

- advantages: good resolution in momentum space
- drawbacks: CR pressure not accounted for in EoM, insufficient spatial resolution (grid code), non-radiative gas physics







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Our philosophy and description

An accurate description of CRs should follow the evolution of the spectral energy distribution of CRs as a function of time and space, and keep track of their dynamical, non-linear coupling with the hydrodynamics.

We seek a compromise between

- capturing as many physical properties as possible
- requiring as little computational resources as necessary

Assumptions:

- protons dominate the CR population
- a momentum power-law is a typical spectrum
- CR energy & particle number conservation





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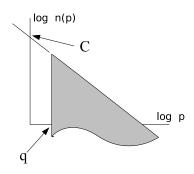
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CR spectral description



$$p = P_{\rm p}/m_{\rm p}\,c$$

$$f(p) = \frac{dN}{dp\,dV} = C\,p^{-\alpha}\theta(p-q)$$

$$egin{aligned} q(
ho) &= \left(rac{
ho}{
ho_0}
ight)^{rac{1}{3}} q_0 \ C(
ho) &= \left(rac{
ho}{
ho_0}
ight)^{rac{lpha+2}{3}} C_0 \end{aligned}$$

$$C(\rho) = \left(\frac{\rho}{\rho_0}\right)^{\frac{3}{3}} C_0$$

$$n_{\rm CR} = \int_0^\infty {\rm d} p \, f(p) = \frac{C \, q^{1-\alpha}}{\alpha-1}$$

$$P_{\mathsf{CR}} = \frac{m_{\mathsf{p}}c^2}{3} \int_0^\infty \mathsf{d}p \, f(p) \, \beta(p) \, p$$

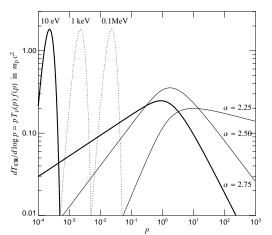
$$=\frac{C m_{\rm p}c^2}{6} \mathcal{B}_{\frac{1}{1+c^2}}\left(\frac{\alpha-2}{2},\frac{3-\alpha}{2}\right)$$





Thermal & CR energy spectra

Kinetic energy per logarithmic momentum interval:

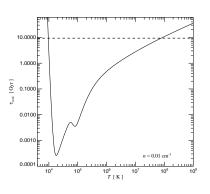




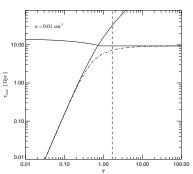


Cooling time scales of CR protons

Cooling of primordial gas:



Cooling of cosmic rays:







Outline

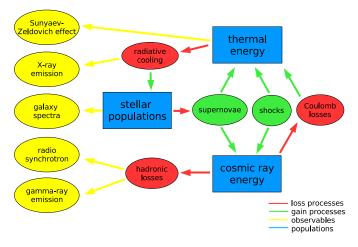
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Radiative simulations with CR physics

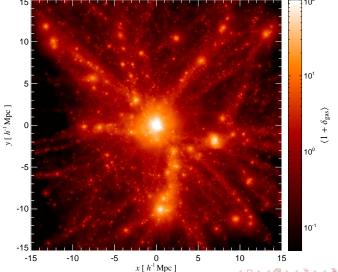
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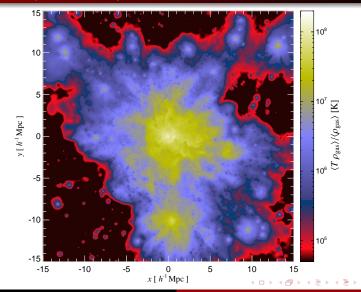


Radiative cool core cluster simulation: gas density



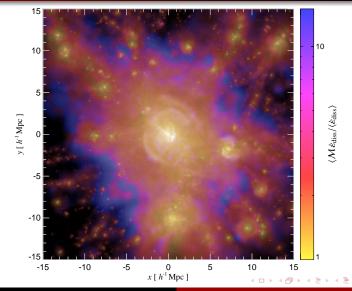


Mass weighted temperature



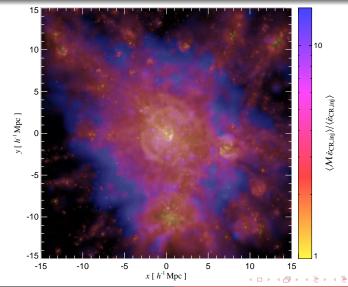


Mach number distribution weighted by $\varepsilon_{\text{diss}}$



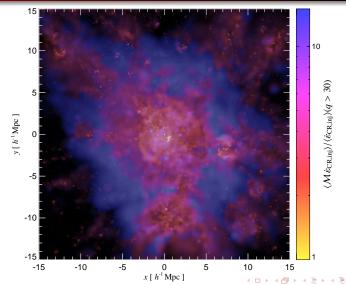


Mach number distribution weighted by $\varepsilon_{\mathrm{CR,inj}}$



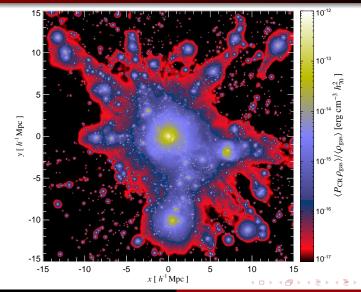


Mach number distribution weighted by $\varepsilon_{CR,ini}(q > 30)$





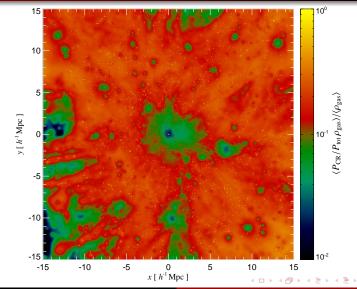
CR pressure P_{CR}





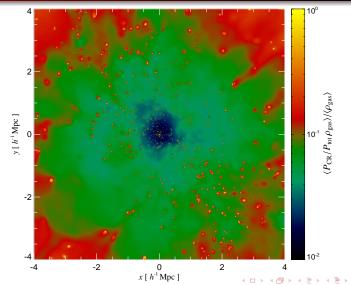
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Relative CR pressure P_{CR}/P_{total}



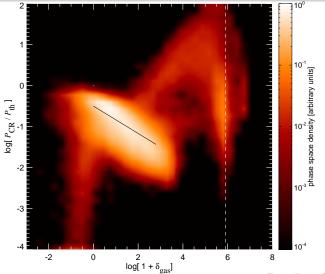


Relative CR pressure P_{CR}/P_{total}





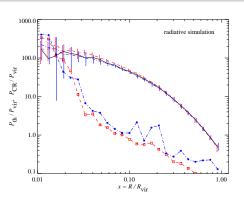
Phase-space diagram of radiative cluster simulation

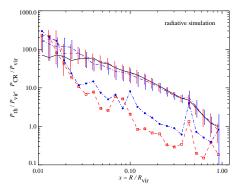






Radiative simulations: pressure profile





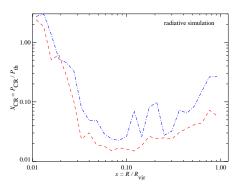
Cool core cluster sample.

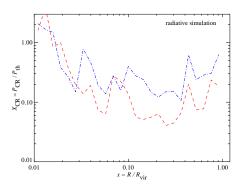
red: only structure formation shock CRs, blue: structure formation & SNe CRs.

Merging cluster sample.



Radiative simulations: relative CR pressure profile





Cool core cluster sample.

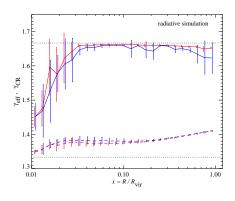
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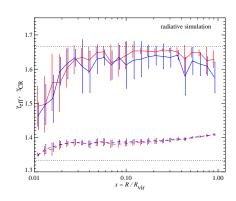
Merging cluster sample.





Radiative simulations: adiabatic index profile





Cool core cluster sample.

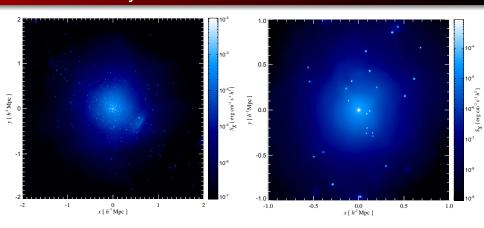
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Merging cluster sample.





Thermal X-ray emission

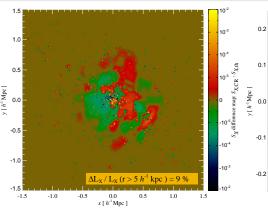


large merging cluster, $\textit{M}_{\rm vir} \simeq 10^{15} \textit{M}_{\odot}/\textit{h}$

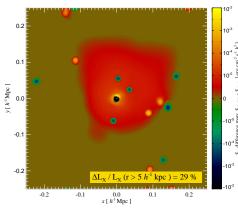
small cool core cluster, $M_{\rm vir} \simeq 10^{14} M_{\odot}/h$



Difference map of S_X : $S_{X,CR} - S_{X,th}$



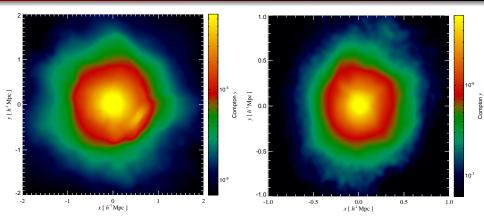
large merging cluster, $M_{\rm vir} \simeq 10^{15} M_{\odot}/h$ \rightarrow contributes to the scatter in the $M-L_{\rm X}$ scaling relation



cool core cluster, $M_{\rm vir} \simeq 10^{14} M_{\odot}/h$ \rightarrow systematic increase of $L_{\rm X}$ for small cool core clusters



Compton y parameter in radiative cluster simulation



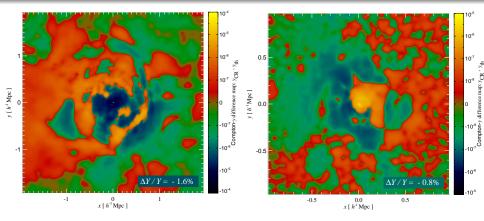
large merging cluster, $M_{\rm vir} \simeq 10^{15} M_{\odot}/h$

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Compton y difference map: $y_{CR} - y_{th}$



large merging cluster, $M_{\rm vir} \simeq 10^{15} M_{\odot}/h$

small cool core cluster, $\textit{M}_{\rm vir} \simeq 10^{14} \textit{M}_{\odot}/h$





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Particle acceleration processes

particles are acclerated via:

- adiabatic compression
- diffusive shock acceleration (Fermi I)
- stochastic acceleration by plasma waves (Fermi II)
- particle reactions (pp $\rightarrow \pi \rightarrow \mu \nu \rightarrow e \nu \nu$)

particles are de-accelerated via

- adiabatic expansion
- radiative cooling (synchrotron, inverse Compton, bremsstrahlung, hadronic interactions)
- non-radiative cooling (Coulomb interactions)





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Diffusive shock acceleration – Fermi 1 mechanism (1)

conditions:

- a collisionless shock wave
- magnetic fields to confine energetic particles
- plasma waves to scatter energetic particles → particle diffusion
- supra-thermal particles

mechanism:

- supra-thermal particles diffuse upstream across shock wave
- each shock crossing energizes particles through momentum transfer from recoil-free scattering off the macroscopic scattering agents
- momentum increases exponential with number of shock crossings
- number of particles decreases exponential with number of crossings







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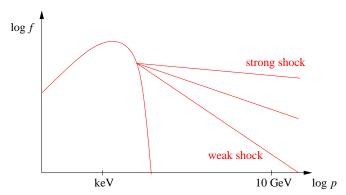




Diffusive shock acceleration – Fermi 1 mechanism (2)

Spectral index depends on the Mach number of the shock,

$$\mathcal{M} = v_{\sf shock}/c_{\sf s}$$
:

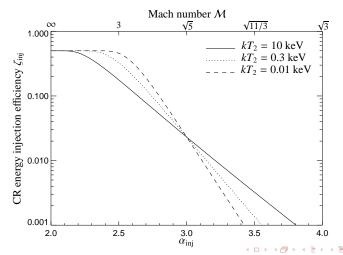






Diffusive shock acceleration – efficiency (3)

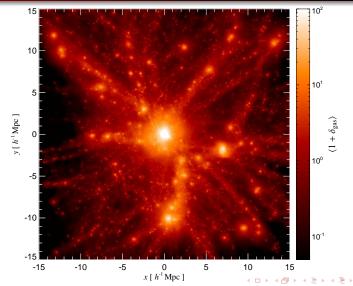
CR proton energy injection efficiency, $\zeta_{\rm inj} = \varepsilon_{\rm CR}/\varepsilon_{\rm diss}$:





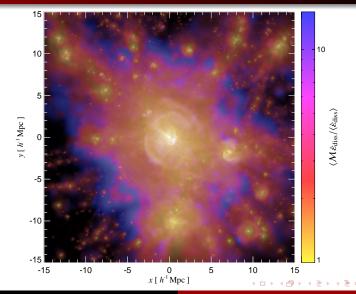


Radiative cool core cluster simulation: gas density



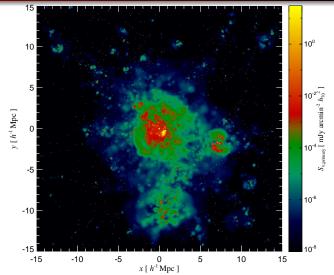


Cosmic web: Mach number





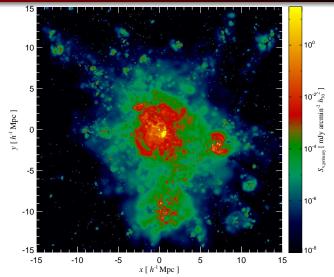
Radio web: primary CRe (1.4 GHz)







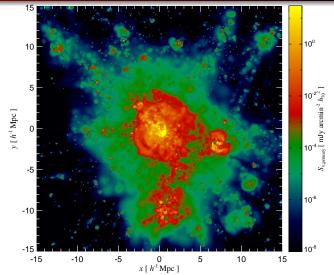
Radio web: primary CRe (150 MHz)







Radio web: primary CRe (15 MHz)

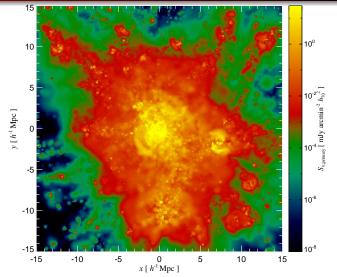






Diffusive shock acceleration Stochastic acceleration

Radio web: primary CRe (15 MHz), slower magnetic decline







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Stochastic acceleration: recipe (1)

conditions:

- super-thermal or better relativistic particles
- magnetic fields to confine them
- high level of plasma waves to scatter them via gyro-resonances

mechanism

- head on wave-particle collision energises particle
- tail on wave-particle collision de-energise particle
- statistically more head-on than tail-on collisions

→ net energy gain due to diffusion in momentum space advantage: plamsa waves are everywhere!





Stochastic acceleration: recipe (1)

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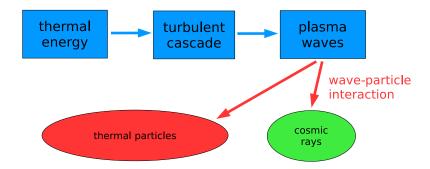
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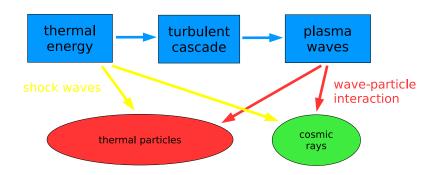
Stochastic acceleration: cartoon (2)







Stochastic acceleration: cartoon (2)







Stochastic acceleration: problems (3)

problems:

- low efficiency (2nd order in ratio of wave to particle velocity)
- waves like to cascade to small scales
- small-scale waves dissipate into the thermal pool
- wave energy budget is usually tight
- at locations with high wave density (e.g. shocks), more efficient acceleration mechanism may be in operation (e.g. DSA)

nevertheless: cluster radio halos may be due to stochastic re-acceleration of 0.2 MeV electrons (e.g. Brunetti et al.)





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Particle reactions

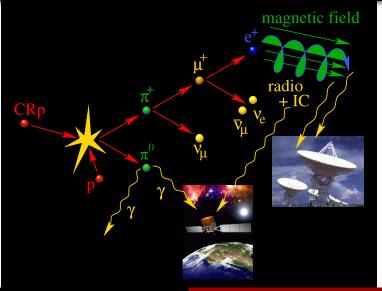
relativistic proton populations can often be expected, since

- acceleration mechanisms work for protons . . .
 - ... as efficient as for electrons (adiabatic compression) or
 - ... more efficient than for electrons (DSA, stochastic acc.)
- galactic CR protons are observed to have 100 times higher energy density than electrons
- CR protons are very inert against radiative losses and therefore long-lived (~ Hubble time in galaxy clusters, longer outside)
- → an energetic CR proton population should exist in clusters



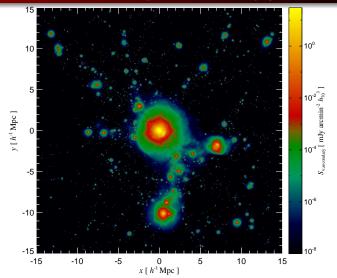


Hadronic cosmic ray proton interaction





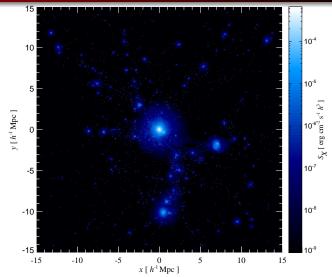
Cluster radio emission by hadronically produced CRe







Thermal X-ray emission







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Non-thermal emission from clusters

Exploring the memory of structure formation

The thermal plasma lost most information on how cosmic structure formation proceeded due to the dissipative processes. The thermal observables, X-ray emission and the Sunyaev-Zel'dovich effect, tell us only very indirectly (if at all) about the cosmic history. In contrast, non-thermal processes retain their cosmic memory since their particle population is not in equilibrium \rightarrow cluster archaeology.

How can we read out this information about non-thermal populations? \rightarrow new era of multi-frequency experiments, e.g.:

- LOFAR, GMRT, MWA: interferometric array of radio telescopes at low frequencies ($\nu \simeq (15-240)$ MHz)
- Simbol-X: future hard X-ray satellite ($E \simeq (0.5 70)$ keV)
- GLAST: high-energy γ -ray space mission ($E \simeq (0.1-300)$ GeV)
- Imaging air Čerenkov telescopes (TeV photon energies)



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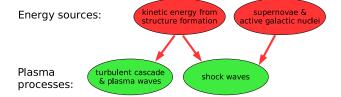
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Cosmic rays and radiative processes

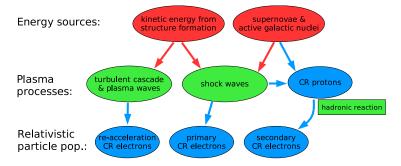






Radiative processes Unified model of radio halos and relic High-energy gamma-ray emission

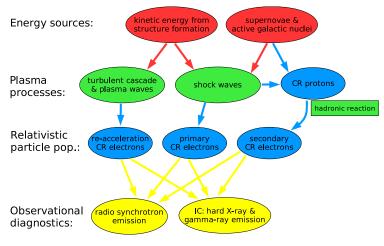
Cosmic rays and radiative processes







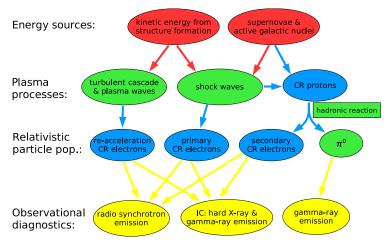
Cosmic rays and radiative processes







Cosmic rays and radiative processes







Outline

- Cosmic rays in galaxy clusters
 - Introduction and motivation
 - Cluster simulations and cosmic ray physics
 - Cosmic ray pressure feedback
- Particle acceleration processes
 - Diffusive shock acceleration
 - Stochastic acceleration
 - Particle reactions
- Non-thermal cluster emission
 - Radiative processes
 - Unified model of radio halos and relics
 - High-energy gamma-ray emission





Previous models for giant radio halos in clusters

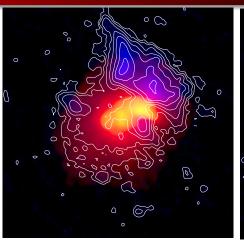
Radio halos show a smooth unpolarized radio emission at Mpc-scales. How are they generated?

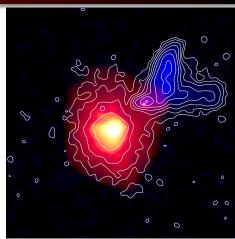
- Primary accelerated CR electrons: synchrotron/IC cooling times too short to account for extended diffuse emission.
- Continuous in-situ acceleration of pre-existing CR electrons either via interactions with magneto-hydrodynamic waves, or through turbulent spectra (Jaffe 1977, Schlickeiser 1987, Brunetti 2001, Brunetti & Lazarian 2007).
- Hadronically produced CR electrons in inelastic collisions of CR protons with the ambient gas (Dennison 1980, Vestrad 1982, Miniati 2001, Pfrommer 2004).

All of these models face theoretical short-comings when comparing to observations.



Which one is the simulation/observation of A2256?

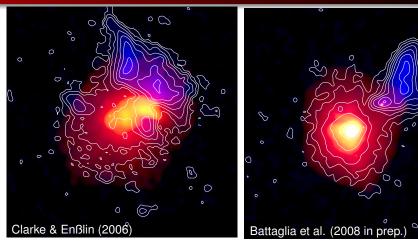




red/yellow: thermal X-ray emission, blue/contours: 1.4 GHz radio emission with giant radio halo and relic



Which one is the simulation/observation of A2256?



red/yellow: thermal X-ray emission, blue/contours: 1.4 GHz radio emission with giant radio halo and relic



Unified model of radio halos and relics

Cluster radio emission varies with dynamical stage of a cluster:

- Cluster relaxes and develops cool core: radio mini-halo develops due to hadronically produced CR electrons, magnetic fields are adiabatically compressed (cooling gas triggers radio mode feedback of AGN that outshines mini-halo → selection effect).
- Cluster experiences major merger: two leading shock waves are produced that become stronger as they break at the shallow peripheral cluster potential → shock-acceleration of primary electrons and development of radio relics.
- Generation of morphologically complex network of virializing shock waves. Lower sound speed in the cluster outskirts lead to strong shocks
 → irregular distribution of primary electrons, MHD turbulence amplifies magnetic fields.
- Giant radio halo develops due to (1) boost of the hadronically generated radio emission in the center (2) irregular radio 'gischt' emission in the cluster outskirts.



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Unified model of radio halos and relics

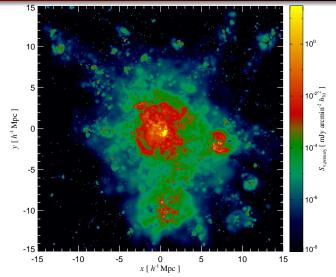
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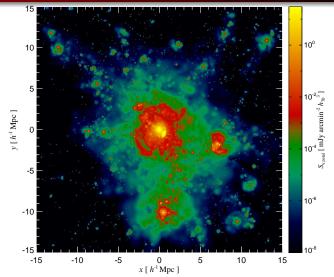
Radio gischt: primary CRe (150 MHz)







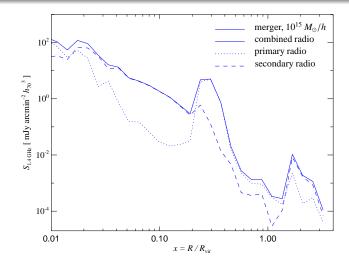
Radio gischt + central hadronic halo = giant radio halo







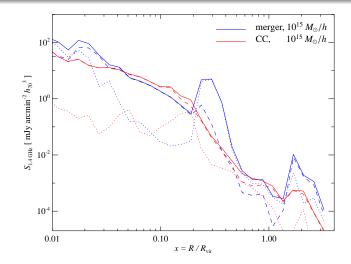
Giant radio halo profile







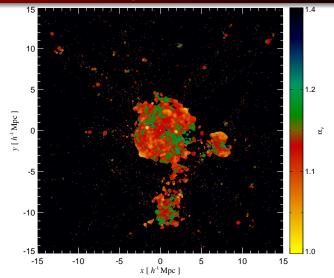
Giant radio halo vs. mini-halo







Radio relics + halos: spectral index







Low-frequency radio emission from clusters

Window into current and past structure formation

Our unified model accounts for . . .

- correlation between merging clusters and giant halos, occurrence of mini-halos in cool core clusters
- observed luminosities of halos/relics for magnetic fields derived from Faraday rotation measurements
- observed morphologies, variations, spectral and polarization properties in radio halos/relics

How we can make use of this information

- Radio relics: produced by primary accelerated CR electrons at formation shocks → probes current dynamical, non-equilibrium activity of forming structures (shocks and magnetic fields)
- Central radio halos: produced by secondary CR electrons in hadronic CR proton interactions → tracing time-integrated non-equilibrium activity, modulated by recent dynamical activities





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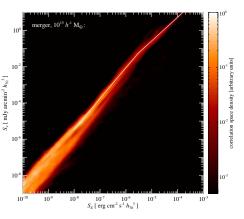
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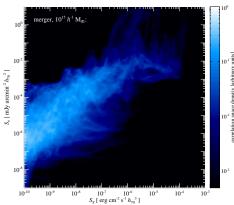




Correlation between X-ray and synchrotron emission



Correlation with secondary 'halo' emission, merging cluster, $M_{\rm vir} \simeq 10^{15} M_{\odot}/h$



Correlation with primary 'relic' emission, merging cluster, $M_{\rm vir} \simeq 10^{15} M_{\odot}/h$





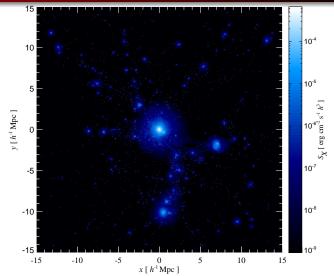
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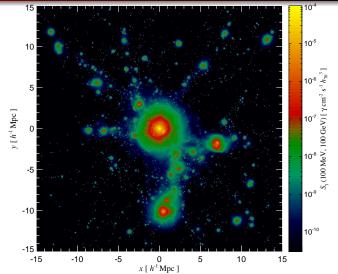
Thermal X-ray emission







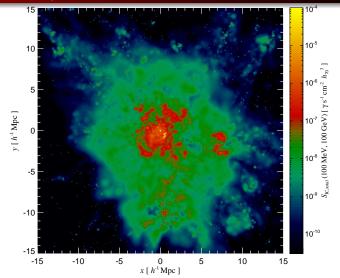
Hadronic γ -ray emission, $E_{\gamma} > 100 \text{ MeV}$







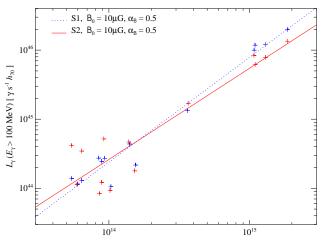
Inverse Compton emission, $E_{IC} > 100 \text{ MeV}$







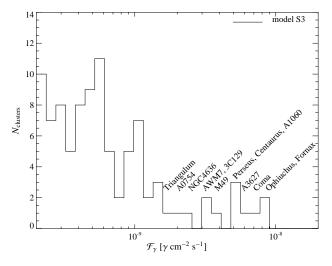
Gamma-ray scaling relations



Scaling relation + complete sample of the brightest X-ray clusters (HIFLUCGS) \rightarrow predictions for GLAST



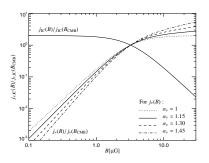
Predicted cluster sample for GLAST







Minimum γ -ray flux in the hadronic model (1)



Synchrotron emissivity of highenergy, steady state electron distribution is independent of the magnetic field for $B\gg B_{\text{CMB}}!$

Synchrotron luminosity:

$$L_{
u} = A_{
u} \int dV \, n_{
m CR} n_{
m gas} rac{arepsilon_{
m B}^{(lpha_{
u}+1)/2}}{arepsilon_{
m CMB} + arepsilon_{
m B}} \
ightarrow A_{
u} \int dV \, n_{
m CR} n_{
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m B} \gg arepsilon_{
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γ -ray luminosity

$$L_{\gamma} = A_{\gamma} \int {
m d} V \, n_{
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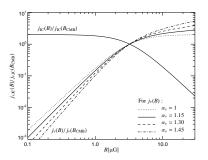
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$$\mathcal{F}_{\gamma,\mathsf{min}} = rac{A_{\gamma}}{A_{
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Synchrotron luminosity:

$$\begin{array}{rcl} \textit{L}_{\nu} & = & \textit{A}_{\nu} \int \textrm{d}\textit{V} \, \textit{n}_{\text{CR}} \textit{n}_{\text{gas}} \frac{\varepsilon_{\textit{B}}^{(\alpha_{\nu}+1)/2}}{\varepsilon_{\text{CMB}} + \varepsilon_{\textit{B}}} \\ & \rightarrow & \textit{A}_{\nu} \int \textrm{d}\textit{V} \, \textit{n}_{\text{CR}} \textit{n}_{\text{gas}} \quad (\varepsilon_{\textit{B}} \gg \varepsilon_{\text{CMB}}) \end{array}$$

 γ -ray luminosity:

$$L_{\gamma} = A_{\gamma} \int \mathrm{d}V \, n_{\mathrm{CR}} n_{\mathrm{gas}}$$

 \rightarrow minimum γ -ray flux:

$$\mathcal{F}_{\gamma, \text{min}} = \frac{A_{\gamma}}{A_{\nu}} \frac{L_{\nu}}{4\pi D^2}$$





Minimum γ -ray flux in the hadronic model (2)

Minimum γ -ray flux ($E_{\gamma} > 100$ MeV) for the Coma cluster:

CR spectral index	2.0	2.3	2.6	2.9
\mathcal{F}_{γ} [10 ⁻¹⁰ γ cm ⁻² s ⁻¹]	0.8	1.6	3.4	7.1

• These limits can be made even tighter when considering energy constraints, $P_B < P_{\rm gas}/20$ and B-fields derived from Faraday rotation studies, $B_0 = 3 \,\mu{\rm G}$:

$$\mathcal{F}_{\gamma, \mathsf{COMA}} \gtrsim 2 \times 10^{-9} \gamma \, \mathit{cm}^{-2} \mathit{s}^{-1} = \mathcal{F}_{\mathsf{GLAST, 2yr}}$$

- Non-detection by GLAST seriously challenges the hadronic model.
- Potential of measuring the CR accleration efficiency for diffusive shock accleration.





Summary – 1. CR pressure feedback

- Oharacteristics of the CRs in clusters:
 - CR proton pressure: time integrated non-equilibrium activities of clusters, modulated by recent mergers.
 - Primary CR electron pressure: resembles current accretion and merging shocks in the virial regions.
- CR pressure modifies the ICM in merging clusters and cooling core regions:
 - Galaxy cluster X-ray emission is enhanced up to 35%, systematic effect in low-mass cooling core clusters.
 - Integrated Sunyaev-Zel'dovich effect remains largely unchanged while the Compton-y profile is more peaked.
 - GLAST should see hadronic γ -ray emission from clusters: measurement of CR protons and origin of radio halos.





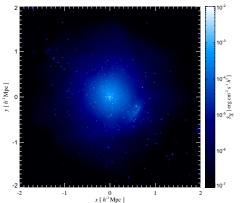
Summary – 2. Non-thermal cluster emission

- Unified model for the generation of giant radio halos, radio mini-halos, and relics:
 - Giant radio halos are dominated in the center by secondary synchrotron emission.
 - Transition to the radio emission from primary electrons in the cluster periphery.
- LOFAR/GMRT are expected to see the radio web emission: origin of cosmic magnetic fields.
- **3** We predict GLAST to detect \sim ten γ -ray clusters: test of the presented scenario
- → exciting experiments allow a complementary view on structure formation as well as fundamental physics!

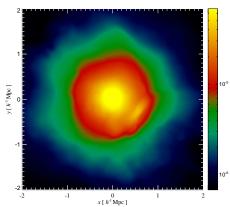




Thermal cluster observables (1)



Thermal bremsstrahlung emission, merging cluster, $M_{\rm vir} \simeq 10^{15} M_{\odot}/h$

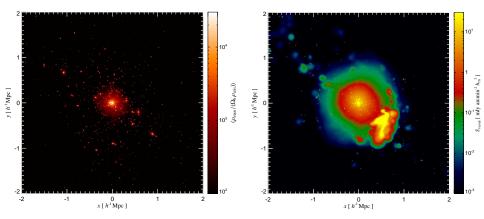


Sunyaev-Zel'dovich effect, $\label{eq:merging} \text{merging cluster, } \textit{M}_{\text{vir}} \simeq 10^{15} \textit{M}_{\odot} / \textit{h}$





Optical and radio synchrotron cluster observables (1)

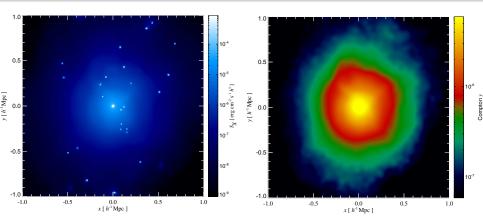


Stellar mass density ("cluster galaxies"), merging cluster, $M_{\rm vir} \simeq 10^{15} M_{\odot}/h$

Radio halo and relic emission, merging cluster, $\textit{M}_{\textrm{vir}} \simeq 10^{15} \textit{M}_{\odot}/\textit{h}$



Thermal cluster observables (2)

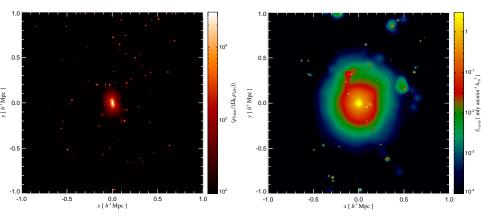


Thermal bremsstrahlung emission, cool core cluster, $M_{\rm vir} \simeq 10^{14} M_{\odot}/h$

Sunyaev-Zel'dovich effect, cool core cluster, $M_{vir} \simeq 10^{14} M_{\odot}/h$



Optical and radio synchrotron cluster observables (2)



Stellar mass density ("cluster galaxies"), cool core cluster, $M_{\rm vir} \simeq 10^{14} M_{\odot}/h$

Radio halo and relic emission, cool core cluster, $\textit{M}_{\textrm{vir}} \simeq 10^{14} \textit{M}_{\odot}/\textit{h}$



