

Magnetic draping – from space physics to galaxy clusters and cosmology

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in collaboration with

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Outline

- 1 **Magnetic draping**
 - Physics
 - Solar system
 - Galaxy clusters
- 2 **Spiral galaxies**
 - Polarized radio ridges
 - Draping and synchrotron emission
 - Magnetic coherence scale
- 3 **Implications**
 - Magnetic field orientations
 - Kinetic plasma instabilities
 - Cosmological evolution of galaxy clusters

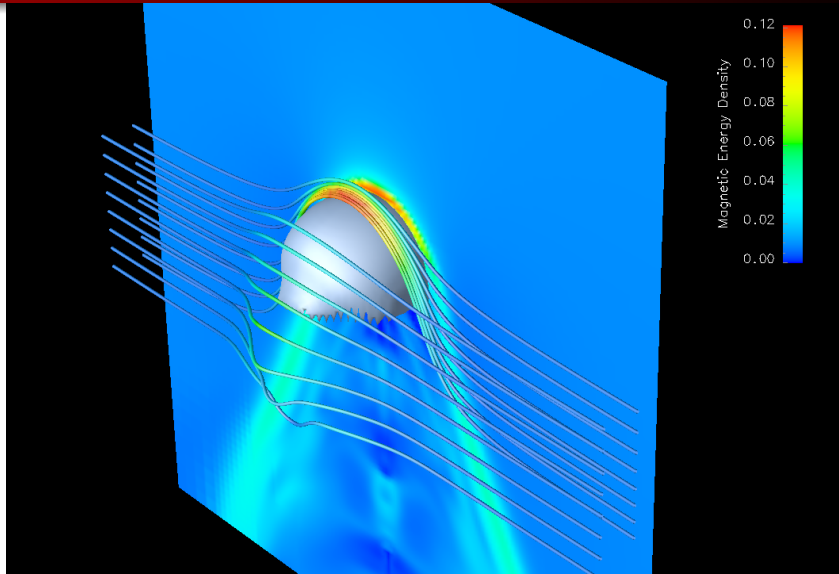


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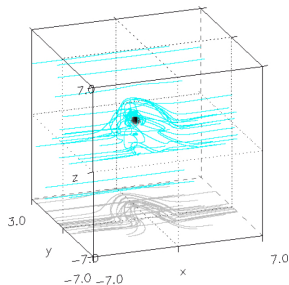


Draping field lines around a moving object



Draping of solar wind field around the Earth

- the Earth's dipolar field shields the surface from penetrating cosmic rays
- the magnetic dipole has reversed sign some hundreds of times over the last 400 million years, which may correspond to breakdowns of the dynamo action

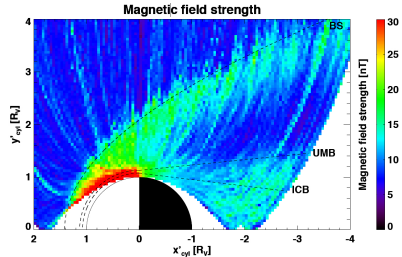


Birk, Lesch & Konz (2004)

- 3D plasma-neutral gas simulations show that the solar wind can induce very fast (~ 10 min) a strong magnetic field in the previously completely unmagnetized Earth's ionosphere
- **Earth magnetic polarity reversals may not be catastrophic to life!**

Draping of the interplanetary field over Venus

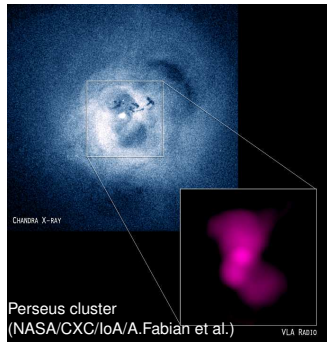
- Venus and Mars do not have a global magnetic field
- *right*: spatial distribution of the magnetic field strength in the plasma environment surrounding Venus (Venus Express)
- draping of solar wind magnetic field around Venus/Mars leads to the **formation of magnetic pile-up region and the magnetotail** → enhanced magnetic field strength in the planets' wake



Guicking et al. (2010)

Puzzles in galaxy clusters

- radio bubbles, seen as X-ray cavities, are observed out to large distances and have very sharp interfaces: **hydrodynamic instabilities should disrupt them**
- high-resolution X-ray data reveal 'cold fronts' with sharp edges in temperature and density: they are **not expected to remain sharp in the presence of diffusion and thermal conduction** for $\gtrsim 10^8$ yrs



→ **Could bubble/core motions sweep up enough magnetic field to suppress instabilities and diffusion/conduction across the interface?**

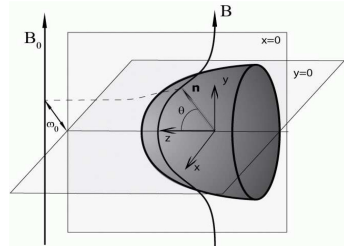


Idea: magnetic draping comes to rescue (Lyutikov 2004)

- analytics, B-profile along the stagnation line:

$$\frac{B}{\rho} = \frac{1}{\sqrt{1 - \frac{R^3}{r^3}}} \frac{B_0}{\rho_0}, \quad l_{\text{drape}} \approx \frac{1}{\mathcal{M}_A^2} R$$

- formula predicts infinite B-amplification at the contact that is in conflict with the kinematic assumption of negligible back-reaction

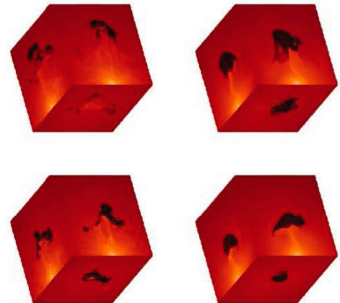


Lyutikov (2004)

→ need MHD simulations to account for the non-linear feedback!

Magnetic draping at radio bubbles

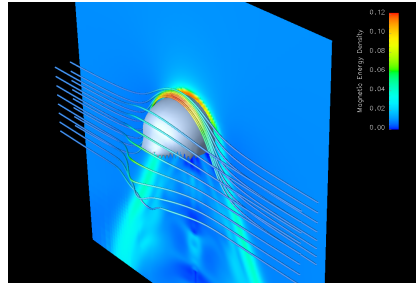
- rising radio bubbles in a hot atmosphere
- shown is the log of the density for the non-draping versus draping case (time increasing upwards)
→ draping suppresses hydrodynamical instabilities in accordance with observations



Ruszkowski et al. (2007)

Magnetic draping at cold fronts

- dense core of a merging cluster performs sloshing motions in the bottom of the larger clusters' potential well
- unavoidable draping of a dynamically strong magnetic layer around the contact surface



Dursi & C.P. (2008)

- **draped magnetic field suppresses diffusion/conduction** across the interface → sharp density/temperature edges!

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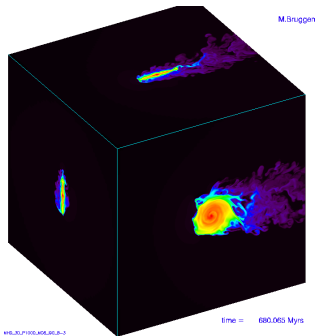
Polarized synchrotron emission in a field spiral: M51



MPIfR Bonn and Hubble Heritage Team

- polarized synchrotron intensity follows the spiral pattern and is strongest in between the spiral arms (NGC 6946)
- the polarization 'B-vectors' are aligned with the spiral structure
- a promising generating mechanism is the *dynamo which transfers mechanical into magnetic energy* (Beck et al. 1996)
- efficient dynamo needs turbulent motions and non-uniform (differential) rotation of the disk

Ram-pressure stripping of cluster spirals

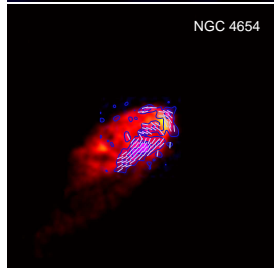
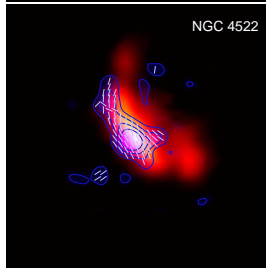
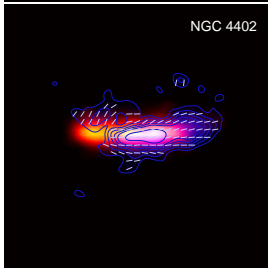
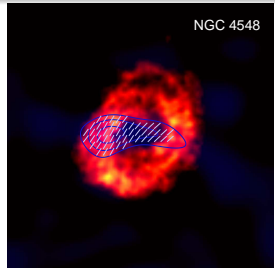
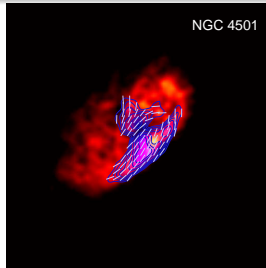
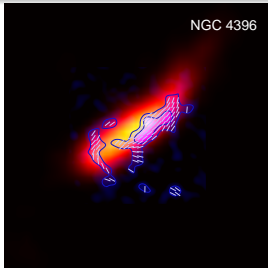


Brueggen (JU Bremen)

- 3D hydrodynamical simulations show that low-density gas in between spiral arms is quickly stripped irrespective of disk radius (Tonnesen & Bryan 2010)
- being flux-frozen into this dilute plasma, the large scale field will also be stripped, leaving behind the small scale field in the star forming regions

→ beam depolarization effects and superposition of causally unconnected star forming patches along the line-of-sight cause the **resulting radio synchrotron emission to be effectively unpolarized**

Polarized synchrotron ridges in Virgo spirals



Vollmer et al. (2007): 6 cm PI (contours) + B-vectors; Chung et al. (2009): HI (red)



Observational evidence and model challenges

- asymmetric distributions of polarized intensity at the leading edge with extraplanar emission, sometimes also at the side
- coherent alignment of polarization vectors over ~ 30 kpc
- stars lead polarized emission, polarized emission leads gas
- HI gas only moderately enhanced (factor $\lesssim 2$), localized ‘HI hot spot’ smaller than the polarized emission region:
$$n_{\text{compr}} \simeq n_{\text{icm}} v_{\text{gal}}^2 / c_{\text{ism}}^2 \simeq 1 \text{ cm}^{-3} \simeq \langle n_{\text{ism}} \rangle$$
- flat radio spectral index (similar to the Milky Way) that steepens towards the edges of the polarized ridge
- no or weak Kelvin-Helmholtz instabilities at interface detectable

→ previous models that use ram-pressure compressed galactic magnetic fields fail to explain most of these points!



Observational evidence and model challenges

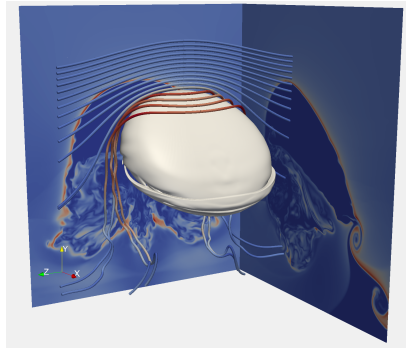
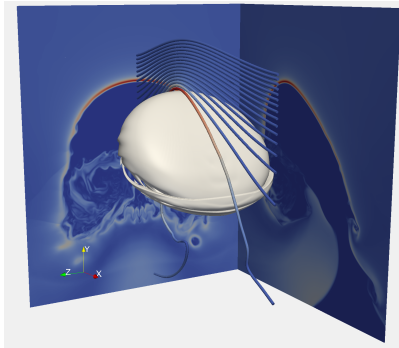
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→ need to consider the full MHD of the interaction spiral galaxy and magnetized ICM !

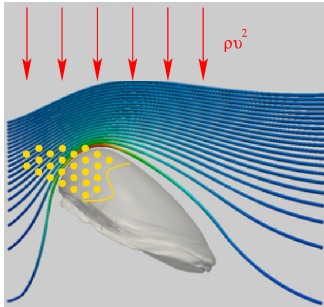


Magnetic draping around a spiral galaxy – MHD



Athena simulations of spiral galaxies interacting with a uniform cluster magnetic field. There is a **sheath of strong field draped around the leading edge** (field strength is color coded).

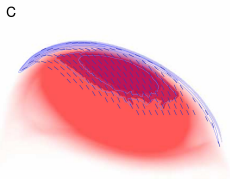
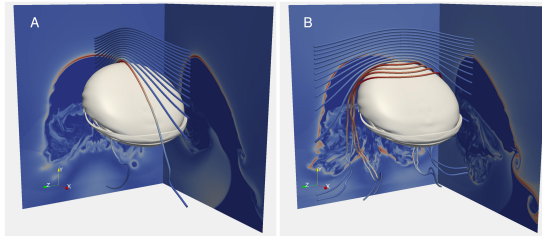
Magnetic draping around a spiral galaxy – physics



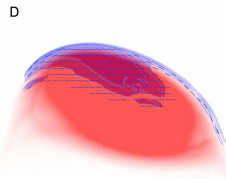
- the galactic ISM is pushed back by the ram pressure wind $\sim \rho v^2$
 - the stars are largely unaffected and lead the gas
 - the draping sheath is formed at the contact of ISM/ICM
 - as stars become SN, their remnants accelerate CRes that populate the field lines in the draping layer
-
- CRes are transported diffusively (along field lines) and advectively as field lines slip over the galaxy
 - CRes emit radio synchrotron radiation in the draped region, tracing out the field lines there → **coherent polarized emission at the galaxies' leading edges**

Magnetic draping and polarized synchrotron emission

Synchrotron B-vectors reflect the upstream orientation of cluster magnetic fields

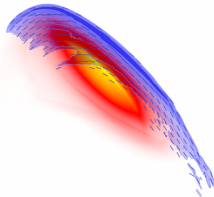


Total PI = 8.227 mJy
Max PI = 218.7 μ Jy/beam

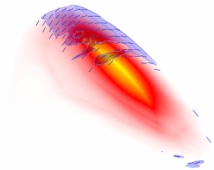
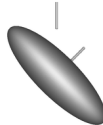


Total PI = 8.440 mJy
Max PI = 334.6 μ Jy/beam

Simulated polarized synchrotron emission



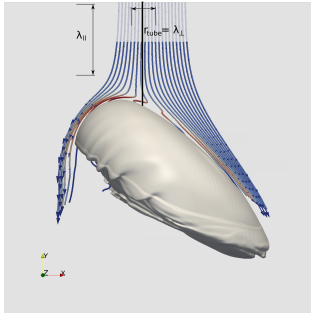
Total PI (mJ) = 23.47
Max PI (μ J/beam) = 3002.



Total PI (mJ) = 4.114
Max PI (μ J/beam) = 133.9

Movie of the simulated polarized synchrotron radiation viewed from various angles and with two field orientations.

Streamlines in the rest frame of the galaxy

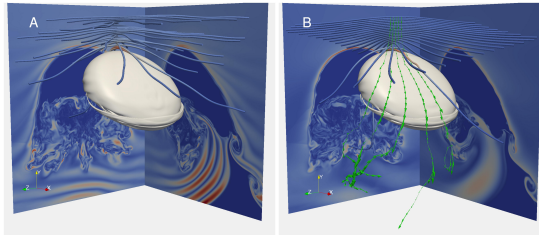


- as the flow approaches the galaxy it decelerates and gets deflected
- only those streamlines initially in a narrow tube of radius $r_{\text{tube}} = \lambda_{\perp}$ from the stagnation line become part of the magnetic draping layer (color coded) \rightarrow constraints on λ_B

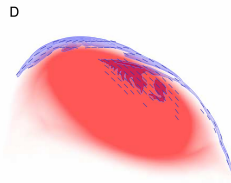
- the streamlines that do not intersect the tube get deflected away from the galaxy, become never part of the drape and eventually get accelerated (Bernoulli effect)
- note the kink feature in some draping-layer field lines due to back reaction as the solution changes from the hydrodynamic potential flow solution to that in the draped layer

Magnetic draping of a non-uniform B-field

(Non-)observation of polarization twist constrains magnetic coherence length



Total PI = 1.586 mJ
Max PI = 67.42 μ J/beam

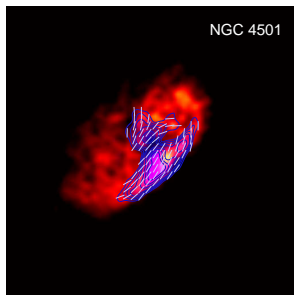


Total PI = 5.927 mJ
Max PI = 304.9 μ J/beam



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Magnetic coherence scale estimate by radio ridges



- observed polarised draping emission
→ field coherence length λ_B is at least galaxy-sized
- if $\lambda_B \sim 2R_{\text{gal}}$, then the change of orientation of field vectors imprint as a change of the polarisation vectors along the vertical direction of the ridge showing a ‘polarisation-twist’
- the reduced speed of the boundary flow means that a small L_{drape} corresponds to a larger length scale of the unperturbed magnetic field ahead of the galaxy NGC 4501

$$L_{\text{coh}} \simeq \eta L_{\text{drape}} v_{\text{gal}} / v_{\text{drape}} = \eta \tau_{\text{syn}} v_{\text{gal}} > 100 \text{ kpc},$$

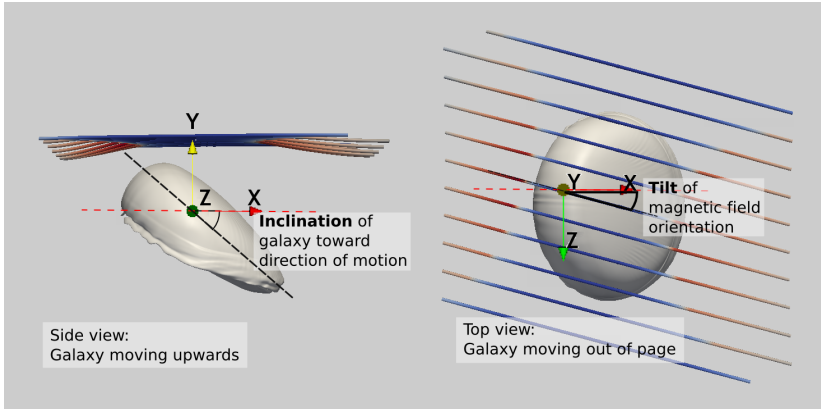
with $\tau_{\text{syn}} \simeq 5 \times 10^7 \text{ yr}$, $v_{\text{gal}} \simeq 1000 \text{ km/s}$, and a geometric factor $\eta \simeq 2$

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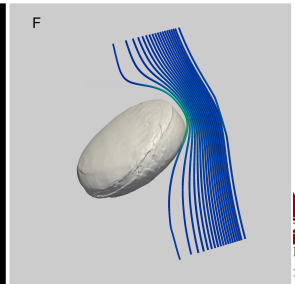
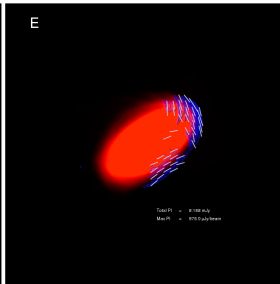
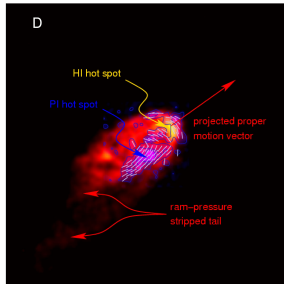
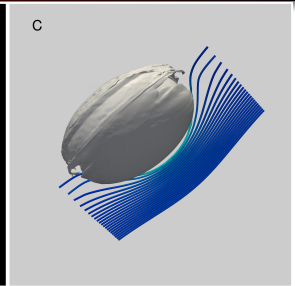
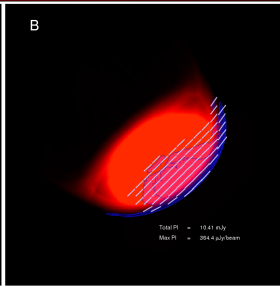
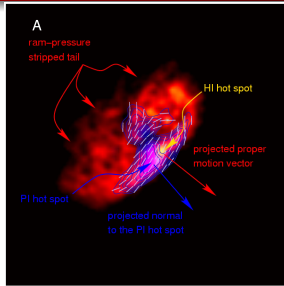
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Varying galaxy inclination and magnetic tilt

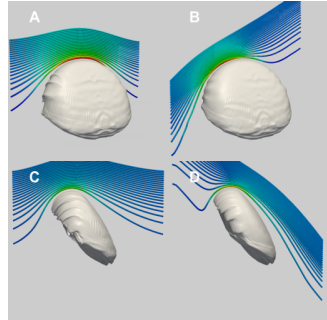


Observations versus simulations

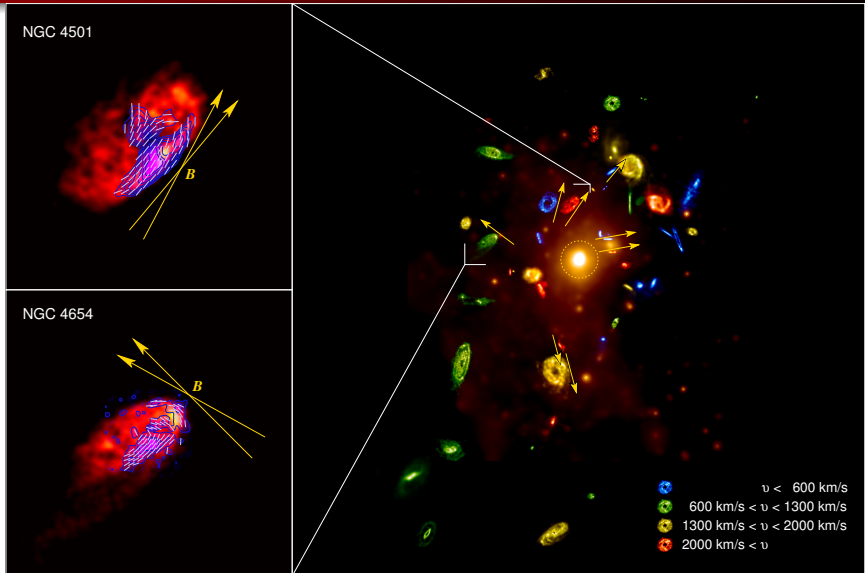


Biases in inferring the field orientation

- uncertainties in estimating the 3D velocity: v_r , ram-pressure stripped gas visible in HI morphology $\rightarrow \hat{\mathbf{v}}_t$
- *direction-of-motion asymmetry*: magnetic field components in the direction of motion bias the location of $B_{\text{max, drape}}$ (figure to the right): draping is absent if $\mathbf{B} \parallel \mathbf{v}_{\text{gal}}$
- *geometric bias*: polarized synchrotron emission only sensitive to traverse magnetic field B_t (\perp to LOS) \rightarrow maximum polarised intensity may bias the location of $B_{\text{max, drape}}$ towards the location in the drape with large B_t



Mapping out the magnetic field in Virgo



Discussion of radial field geometry

- The alignment of the field in the plane of the sky is **significantly more radial than expected from random chance**. Considering the sum of deviations from radial alignment gives a chance coincidence of less than 1.7% ($\sim 2.2 \sigma$).¹
- For the **three nearby galaxy pairs** in the data set, **all have very similar field orientations**.
- The isotropic distribution with respect to the centre (M87) is **difficult to explain with the past activity of the central AGN**.

→ Which effect causes this field geometry?

¹Caveat: this statistical analysis does not include systematic uncertainties such as line-of-sight effects.



Magneto-thermal instability: the idea

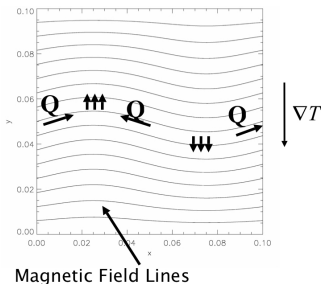


Figure from I. Parrish

Convective stability in a gravitational field:

- Classical Schwarzschild criterion:
 $\frac{dS}{dz} > 0$
- long MFP, Balbus criterion: $\frac{dT}{dz} > 0$
- **new instability causes field lines to reorient radially → efficient thermal conduction radially (close to Spitzer)**

The non-linear behavior of the MTI (Parrish & Stone 2007).

- **Adiabatic boundary conditions for $T(r)$** : the instability can exhaust the source of free energy → isothermal profile
- **Fixed boundary conditions for $T(r)$** : field lines stay preferentially radially aligned (35 deg mean deviation from radial)

Magneto-thermal instability: the idea

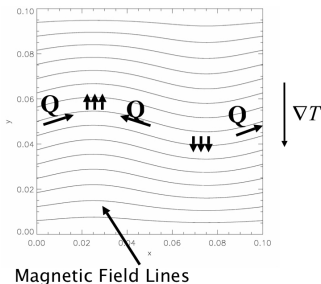


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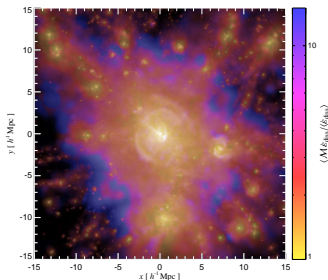
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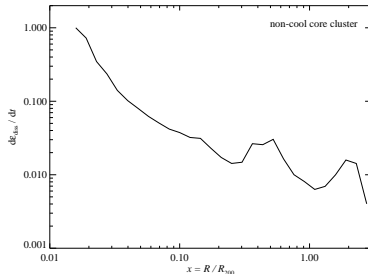


Gravitational shock wave heating

The **observed temperature profile in clusters is decreasing outwards** which is the necessary condition for MTI to operate \rightarrow *gravitational heating can stabilize the temperature profile:*



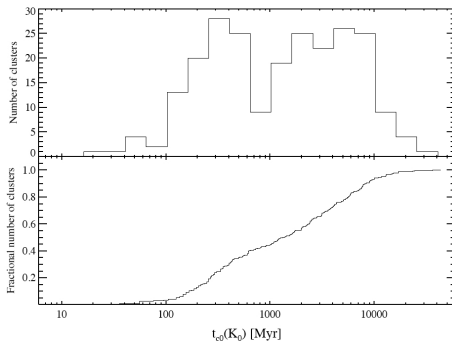
Mach number distribution weighted by ϵ_{diss} .



Energy flux through shock surface
 $\dot{E}_{\text{diss}}/R^2 \sim \rho v^3 \rightarrow$ increase towards the center



Implications for thermal stability of galaxy clusters



Cavagnolo et al. (2009)

- radial fields in non-cool core clusters (NCCs) imply efficient thermal conduction that **stabilizes these systems against entering a cool-core state**: $\tau_{cond} = \lambda^2 / \chi_C \simeq 2.3 \times 10^7 \text{ yr} (\lambda/100 \text{ kpc})^2$, where χ_C is the Spitzer thermal diffusivity (using $kT = 10 \text{ keV}$, $n = 5 \times 10^{-3} \text{ cm}^{-3}$)
- current cosmological cluster simulations fail to reproduce NCCs that have no AGN activity → **MHD + anisotropic conduction**



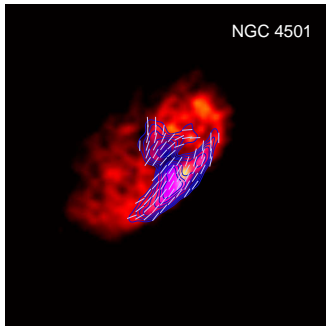
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Speculation: evolutionary sequence of galaxy clusters

- After a merging event of a non-cool core cluster, the **injected turbulence decays on an eddy turnover time**
 $\tau_{\text{eddy}} \simeq L_{\text{eddy}}/v_{\text{turb}} \sim 300 \text{ kpc}/(300 \text{ km/s}) \sim 1 \text{ Gyr}.$
- The **magneto-thermal instability grows on a similar timescale** of less than 1 Gyr and the magnetic field becomes radially oriented.
- The **efficient thermal conduction stabilizes this cluster** until a cooling instability in the center may cause the cluster to enter a cooling core state – similar to Virgo now – and requires possibly feedback by an active galactic nuclei to be stabilized.

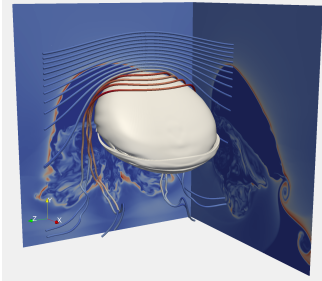


Conclusions on magnetic draping around galaxies



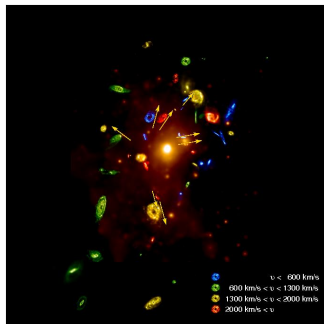
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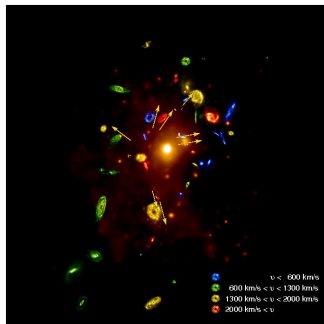
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- application to the Virgo cluster shows that the magnetic field is preferentially aligned radially

Conclusions on magnetic draping around galaxies



- draping of cluster magnetic fields naturally explains polarization ridges at Virgo spirals
 - this represents a new tool for measuring the in situ 3D orientation and coherence scale of cluster magnetic fields
 - application to the Virgo cluster shows that the magnetic field is preferentially aligned radially
-
- this finding is suggestive that the MTI may be operating and implies efficient thermal conduction close to the Spitzer value
 - it also proposes that non-cool core clusters are stabilized by thermal conduction

Literature for the talk

- Pfrommer & Dursi, 2010, Nature Phys., published online, arXiv:0911.2476, *Detecting the orientation of magnetic fields in galaxy clusters*
- Dursi & Pfrommer, 2008, ApJ, 677, 993, *Draping of cluster magnetic fields over bullets and bubbles - morphology and dynamic effects*

