

Radio Halos and Relics in Galaxy Clusters

Christoph Pfrommer¹

in collaboration with Anders Pinzke (DARK, NBI) and Peng Oh (UCSB)

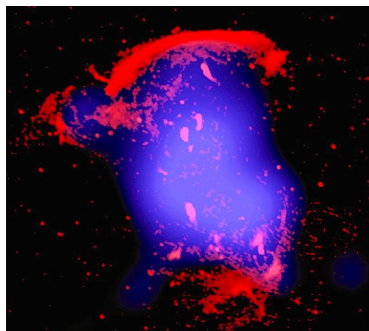
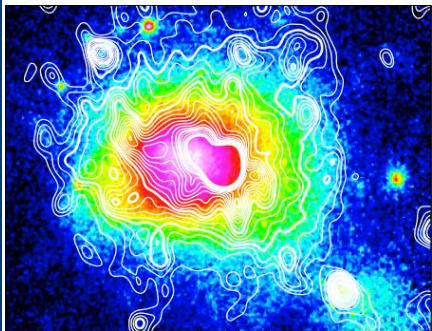
¹Heidelberg Institute for Theoretical Studies, Germany

Jul 2, 2015 / Ringberg Castle: *Cosmic Magnetic Fields*

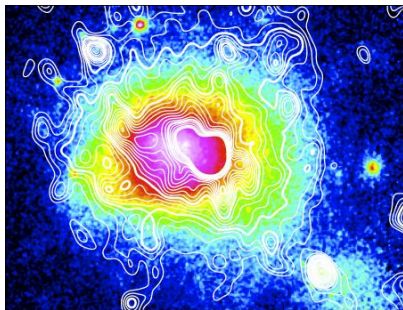




The origin of seed electrons in radio halos and relics



Giant radio halos in a nutshell

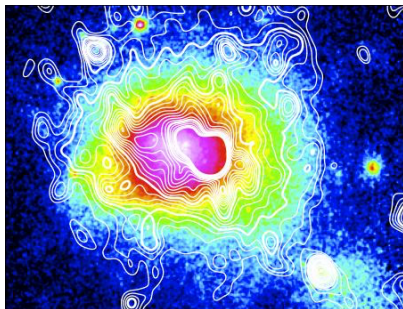


Coma cluster, color: X-ray, contours: radio
X-ray: Snowden/MPE/ROSAT; radio: Brown/Westerbork

- present in > 30 clusters
- Mpc size
- trace X-ray emission
- unpolarized
- steep spectrum $\alpha_\nu \gtrsim 1$
- (sub) $\mu\text{Jy}/\text{arcsec}^2$ surface brightness
- cluster merger connection



Giant radio halos in a nutshell

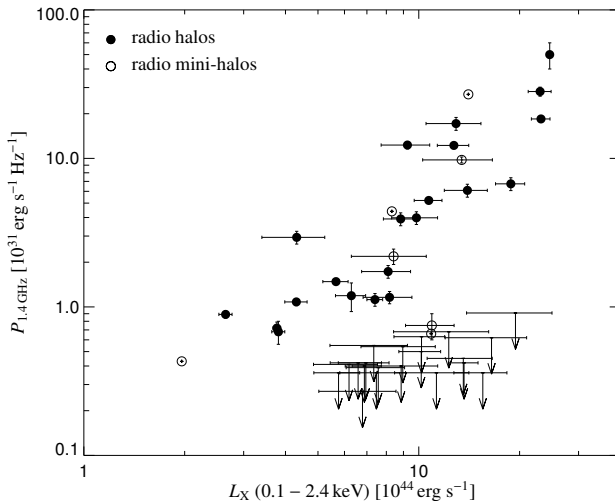


Coma cluster, color: X-ray, contours: radio
X-ray: Snowden/MPE/ROSAT; radio: Brown/Westerbork

- present in > 30 clusters
 - Mpc size
 - trace X-ray emission
 - unpolarized
 - steep spectrum $\alpha_\nu \gtrsim 1$
 - (sub) $\mu\text{Jy}/\text{arcsec}^2$ surface brightness
 - cluster merger connection
-
- evidence of volume-filling magnetic fields in clusters
 - $\tau_{\text{syn}} \lesssim 100 \text{ Myr} \rightarrow$ efficient in-situ electron acceleration



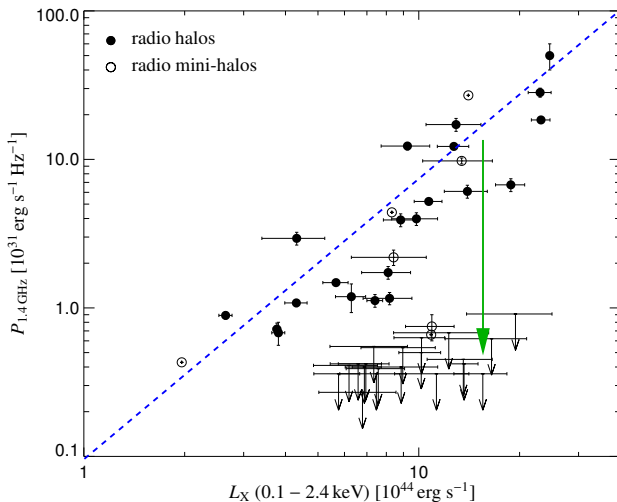
Radio vs. X-ray luminosity – two radio populations



Brunetti+ (2009), Enßlin+ (2011)



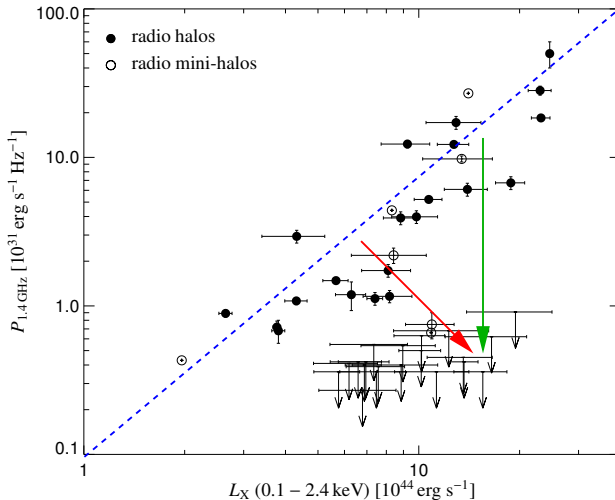
Radio luminosity - X-ray luminosity



Brunetti+ (2009), Enßlin+ (2011)



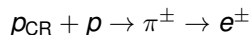
Radio luminosity - X-ray luminosity



Brunetti+ (2009), Enßlin+ (2011)



Radio halo theory – (i) hadronic model

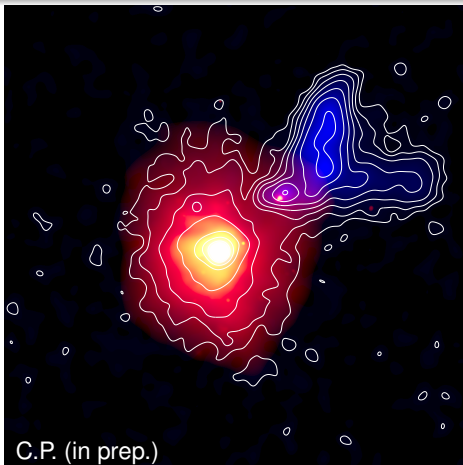
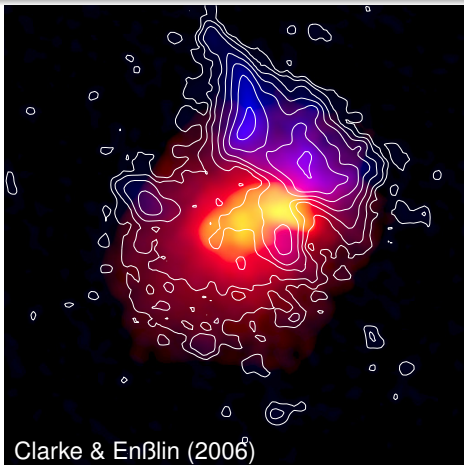


strength:

- all required ingredients available:
shocks to inject CRp, gas protons as targets, magnetic fields
- predicted luminosities and overall morphologies match observations without tuning



Observation – simulation of A2256



red/yellow: thermal X-ray emission,
blue/contours: 1.4 GHz radio emission with giant radio halo and relic



Radio halo theory – (i) hadronic model

$$p_{\text{CR}} + p \rightarrow \pi^{\pm} \rightarrow e^{\pm}$$

strength:

- all required ingredients available:
shocks to inject CRp, gas protons as targets, magnetic fields
- predicted luminosities and overall morphologies match observations without tuning

weakness:

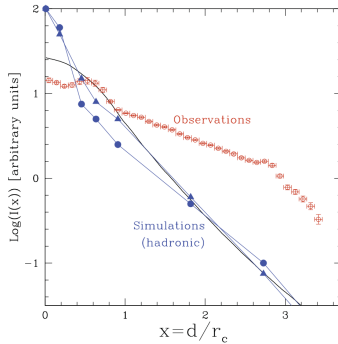
- all clusters should have radio halos
→ putative solution: super-Alfvénic CR streaming (Enßlin+ 2011, Wiener+ 2013)
- does not explain spectral curvature and steep-spectrum sources
→ putative sol.'n: energy-dependent CR diffusion (Enßlin+ 2011, Wiener+ 2013)
- requires increasing CR pressure toward the outskirts of Coma
(Brunetti+ 2013, Zandanel+ 2014)



Coma radio halo: surface brightness profile

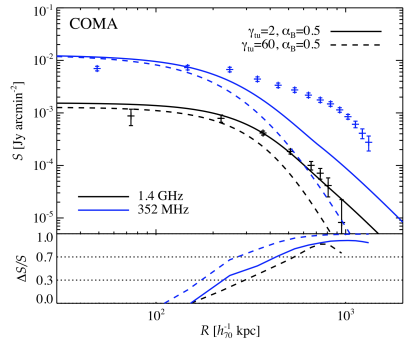
Challenging the hadronic model with extended radio halo profiles?

simulations: pure CR advection



C.P.+ (2008), Pinzke & C.P. (2010), Brunetti+ (2013)

solid: CR streaming $\rightarrow P_{CR} \sim \text{const.}$



Zandanel, C.P., Prada (2014)



Radio halo theory – (ii) re-acceleration model

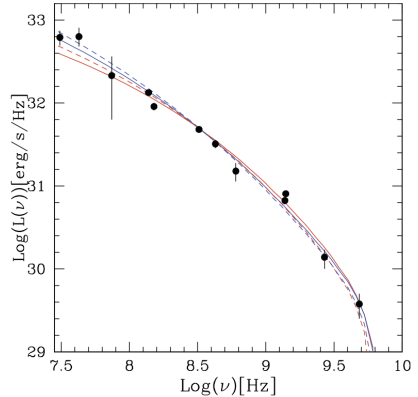
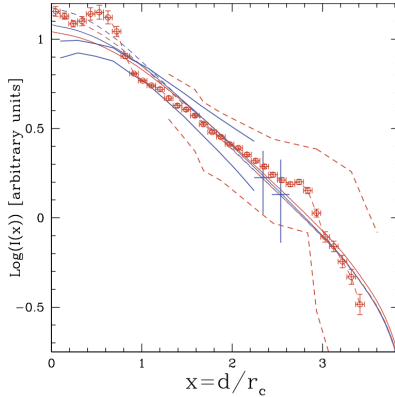
strength:

- all required ingredients available:
radio galaxies & relics to inject CRe, plasma waves to re-accelerate, ...
- reported complex radio spectra emerge naturally
- clusters without halos ← less turbulent



Coma radio halo: re-acceleration model

Good fit to profile and spectrum, **but** many free parameters and assumptions!



Brunetti+ (2013)



Radio halo theory – (ii) re-acceleration model

strength:

- all required ingredients available:
radio galaxies & relics to inject CRe, plasma waves to re-accelerate, ...
- reported complex radio spectra emerge naturally
- clusters without halos ← less turbulent

weakness:

- Fermi II acceleration is inefficient and scales as $(v/c)^2$
- plasma problem: CRe isotropization required by re-acceleration
via transit time damping
- CRe cool rapidly: seed population for re-acceleration?

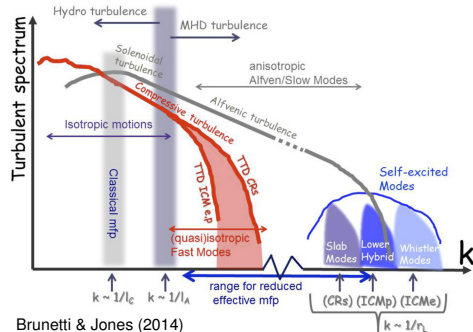


The physics of turbulent re-acceleration

- compressible turbulence can energize particles via gyroresonant interactions

$$\omega - k_{\parallel} v_{\parallel} = n\Omega/\gamma, \quad n = \pm 1, \pm 2, \dots$$

wave vector k_{\parallel} and particle velocity v_{\parallel} are parallel to B and $\Omega = eB/me$



The physics of turbulent re-acceleration

- compressible turbulence can energize particles via gyroresonant interactions

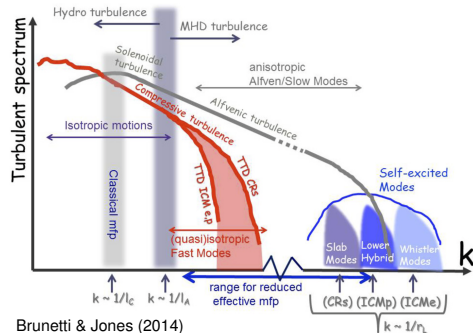
$$\omega - k_{\parallel} v_{\parallel} = n\Omega/\gamma, \quad n = \pm 1, \pm 2, \dots$$

wave vector k_{\parallel} and particle velocity v_{\parallel} are parallel to B and $\Omega = eB/me$

- transit time damping ($n = 0$):

$$v_{\parallel} = \omega/k_{\parallel} = v_{\text{ph},\parallel} \sim c_s$$

→ only *large* pitch-angle CRs can “surf the waves”



The physics of turbulent re-acceleration

- compressible turbulence can energize particles via gyroresonant interactions

$$\omega - k_{\parallel} v_{\parallel} = n\Omega/\gamma, \quad n = \pm 1, \pm 2, \dots$$

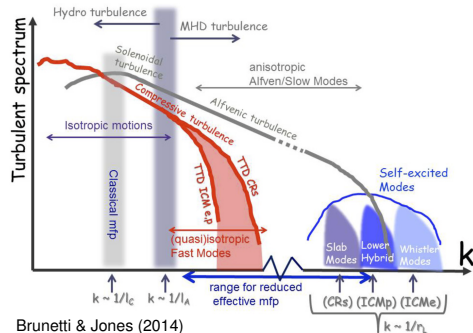
wave vector k_{\parallel} and particle velocity v_{\parallel} are parallel to B and $\Omega = eB/me$

- transit time damping ($n = 0$):

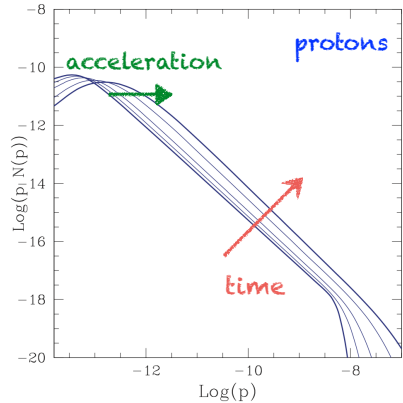
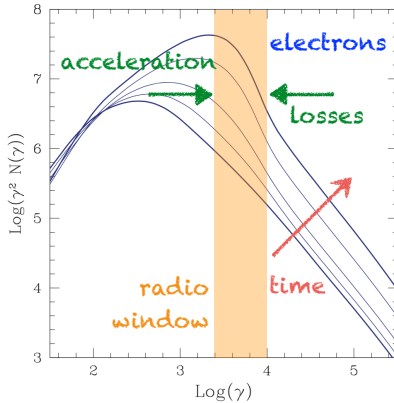
$$v_{\parallel} = \omega/k_{\parallel} = v_{\text{ph},\parallel} \sim c_s$$

→ only *large* pitch-angle CRs can “surf the waves”

- only a fraction of $c_s/c \sim 0.3\%$ goes into CRs, most energy ends up in thermal electrons
- mechanism:** magnetic moment of CRs resonates with the time-varying magnetic field (from the fast modes)



Turbulent re-acceleration: spectral evolution

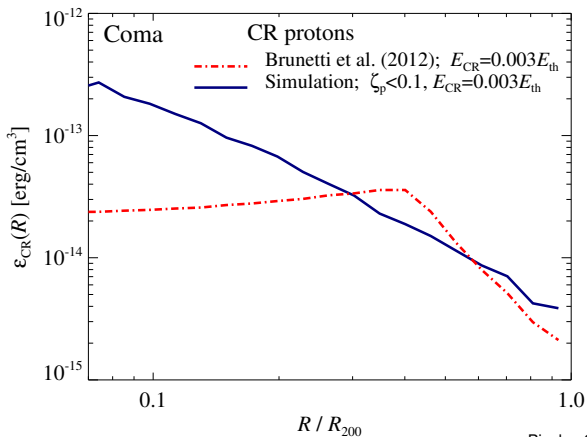


Brunetti & Lazarian (2007, 2011)



But the re-acceleration model has a missing link . . .

. . . it needs seed electrons, which have never been calculated



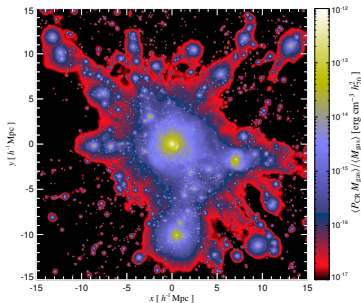
Pinzke, Oh, C.P. (2015)

- cosmological simulations do not reproduce the required population of seed electrons



Method

→ integrate Fokker-Planck equation to follow momentum diffusion in a cosmological simulation with CR proton/electron physics:



P_{CR} in a cosmological zoom simulation of a galaxy cluster (C.P.+ 2008)

$$\begin{aligned} \frac{df_e(\rho, t)}{dt} = & \frac{\partial}{\partial p} \left\{ f_e(\rho, t) \left[\left| \frac{dp}{dt} \right|_c + \frac{p}{3} (\vec{\nabla} \cdot \vec{v}) \right. \right. \\ & \left. \left. + \left| \frac{dp}{dt} \right|_r - \frac{1}{p^2} \frac{\partial}{\partial p} (p^2 D_{pp}) \right] \right\} - (\vec{\nabla} \cdot \vec{v}) f_e(\rho, t) \\ & + \frac{\partial^2}{\partial p^2} [D_{pp} f_e(\rho, t)] + Q_e [p, t; f_p(\rho, t)] \end{aligned}$$

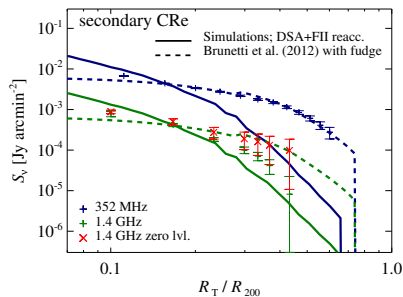
$$\begin{aligned} D_{pp}(\rho, t) = & \frac{\pi}{16} \frac{p^2}{c \rho} \left\langle \frac{\beta |B_k|^2}{16\pi W} \right\rangle l_\theta \int_{k_{\text{cut}}} \mathcal{W}(k) k dk, \\ \mathcal{W}(k) = & \sqrt{2/7} l_0 \rho \langle V_{\text{ph}} \rangle k^{-3/2} \end{aligned}$$



Coma radio halo: multifrequency profiles

even idealized models (Brunetti+ 2013) **have problems:**

→ spectral steepening with radius seen in observations not reproduced with models



Pinzke, Oh, C.P. (2015)

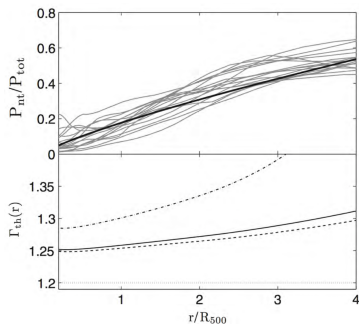
possibilities:

- 1.4 GHz zero-point too high
- observed B -field profile wrong
- new plasma physics

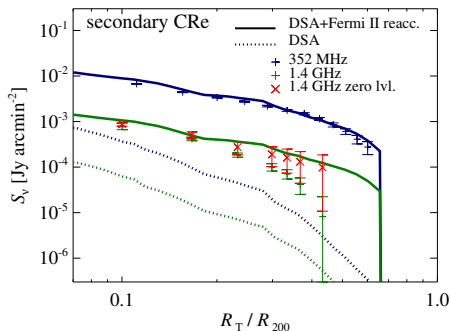
→ can we match the more reliable 352 MHz data? (Brown & Rudnick 2011)



Solution I: changing the turbulent profile



Shaw+ (2010)

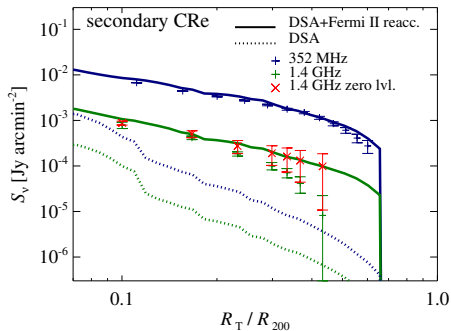
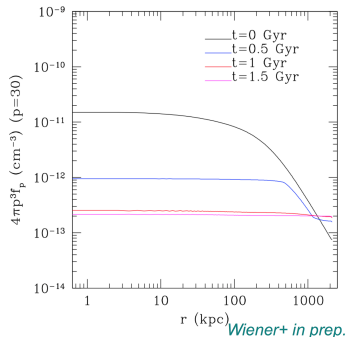


Pinzke, Oh, C.P. (2015)

note: in practice we have to separate compressible turbulence from bulk motions!



Solution II: cosmic-ray streaming



Pinzke, Oh, C.P. (2015)

note: in practice we have to simultaneously simulate cosmic-ray streaming and turbulent re-acceleration!

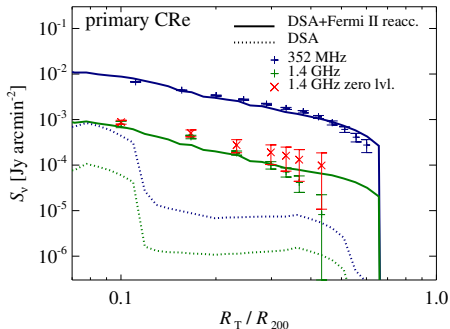


Solution III: primary fossil electrons as seeds

Need high electron acceleration efficiency

recent plasma simulations
with PIC codes ...

- ... find electrons efficiently accelerated in perpendicular shocks (Guo, Sironi, Narayan 2015)
- ... find ions efficiently accelerated in parallel shocks (Caprioli & Spitkovsky 2014)

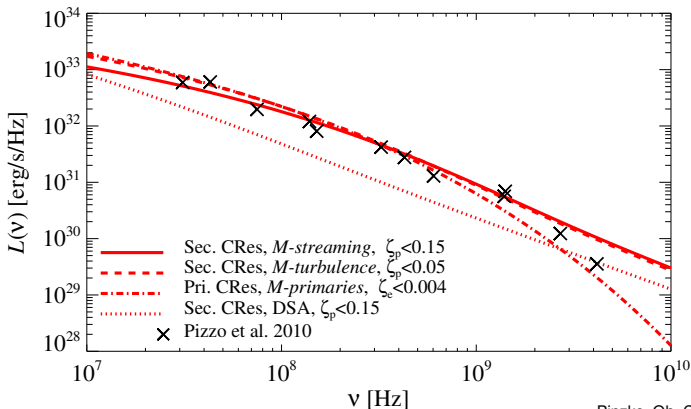


Pinzke, Oh, C.P. (2015)

→ so quasi-perpendicular shock regions might satisfy our requirements!



Coma radio spectrum



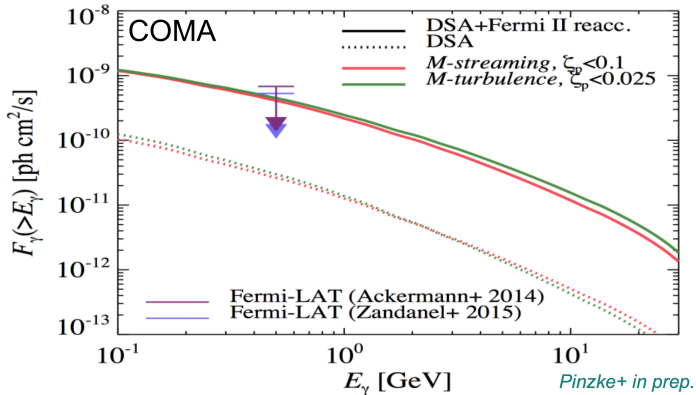
Pinzke, Oh, C.P. (2015)

- all 3 models match the observed radio spectrum
- pure hadronic model fails (only DSA, no turbulent re-acceleration)



How can we disentangle our models?

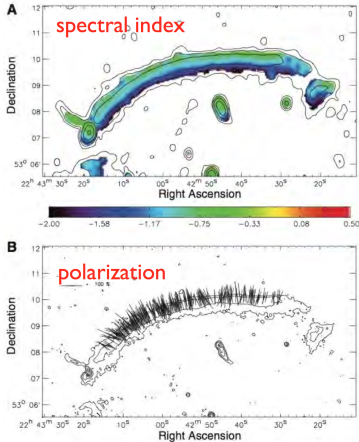
Gamma-ray observations by *Fermi*-LAT are the key



Fermi-LAT can probe M -streaming and M -turbulence in near future!



Radio relics – great tools for studying shock physics

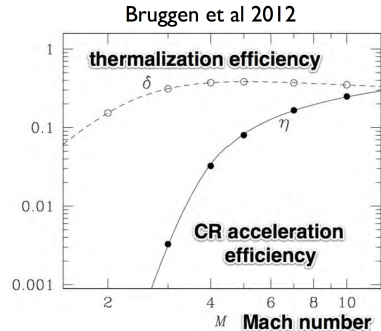
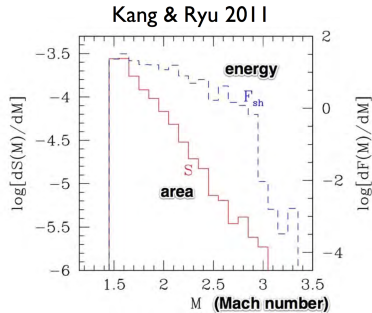


van Weeren+ (2010)

- trace shocks in cluster outskirts
- spectral index: shock Mach number
- spectral ageing: B-field strength
- polarization: B-field orientation



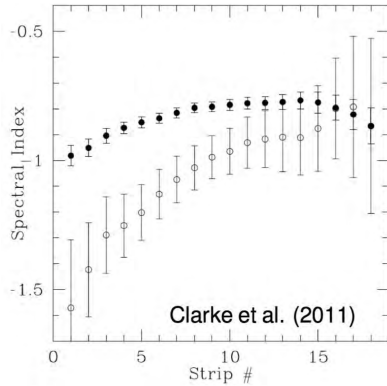
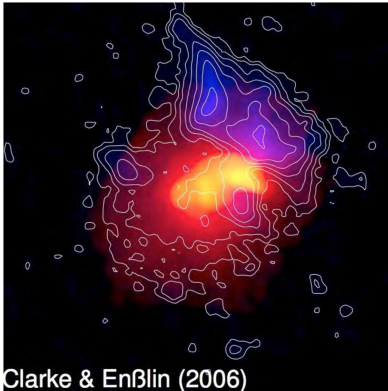
Biggest unknown: shock acceleration efficiency



- merging shocks dominated by low Mach number shocks
- these shocks have low acceleration efficiencies
- how many will LOFAR see?



A poster child: A2256



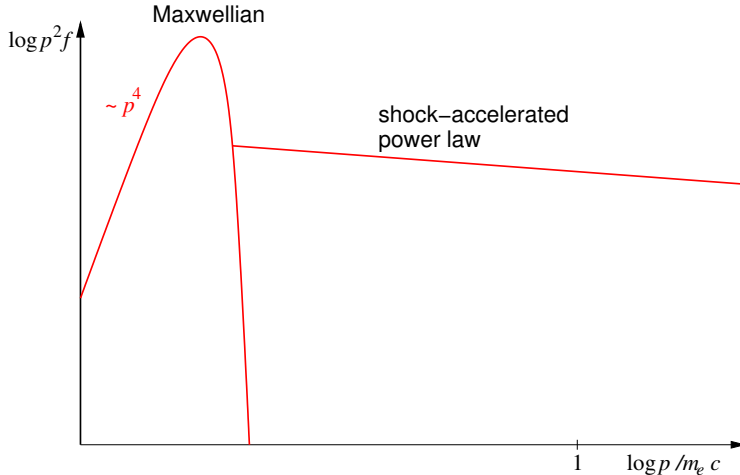
$$\alpha_\nu = 0.85 \quad \rightarrow \quad \mathcal{M} = 2.6:$$

How is this possible?

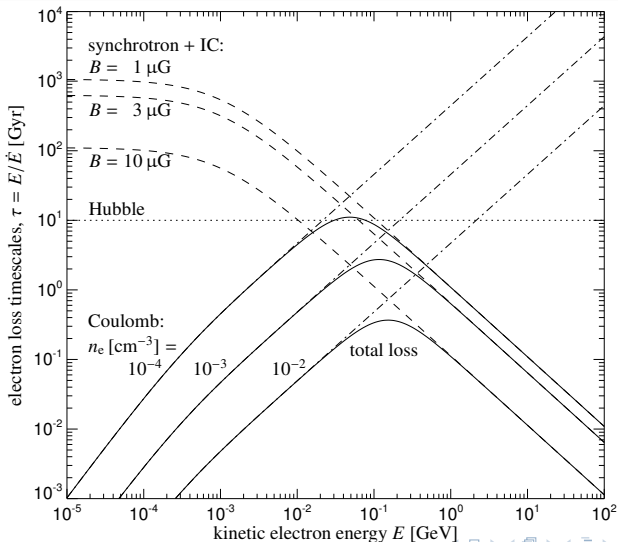


Build-up of the fossil electron distribution

Strong structure formation shocks during the era of cluster formation

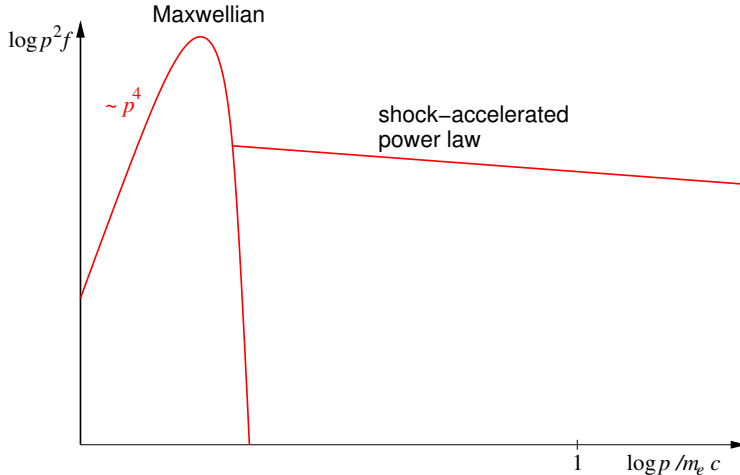


Electron cooling times



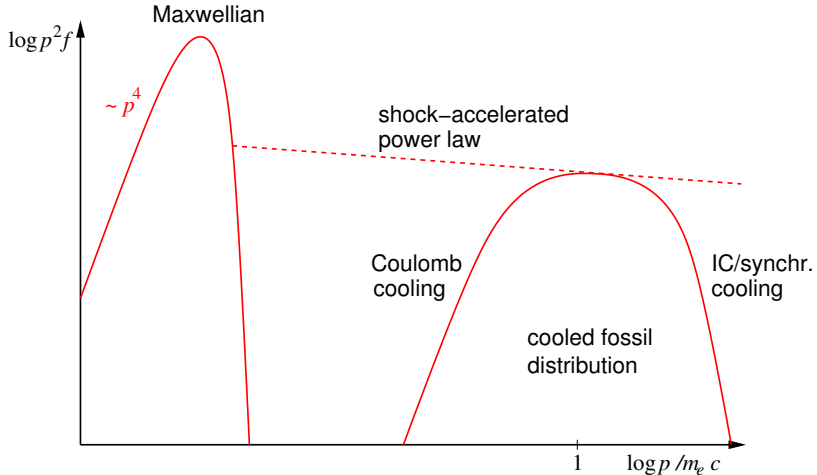
Build-up of the fossil electron distribution

Strong structure formation shocks during the era of cluster formation



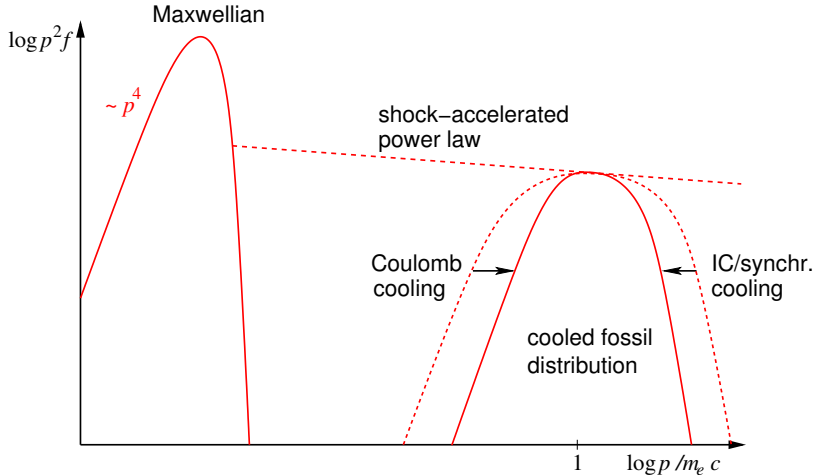
Build-up of the fossil electron distribution

Strong structure formation shocks during the era of cluster formation



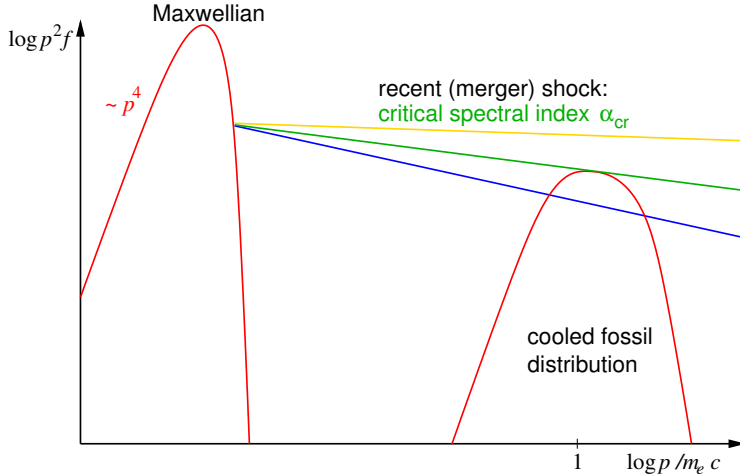
Build-up of the fossil electron distribution

Strong structure formation shocks during the era of cluster formation



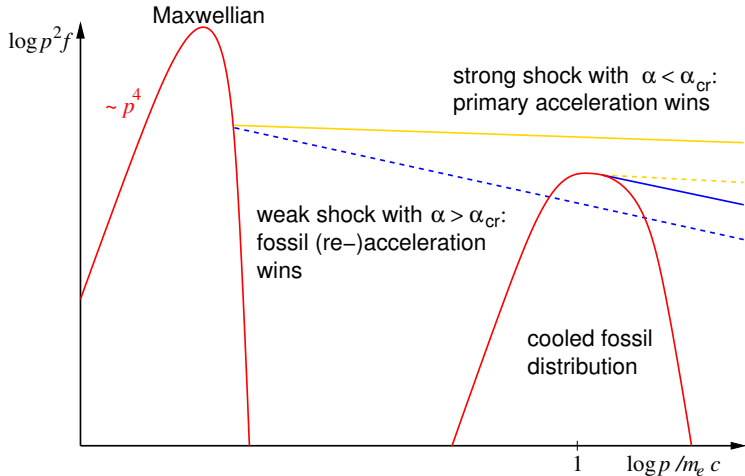
Illuminating radio relics

Re-acceleration of fossil electrons vs. primary acceleration



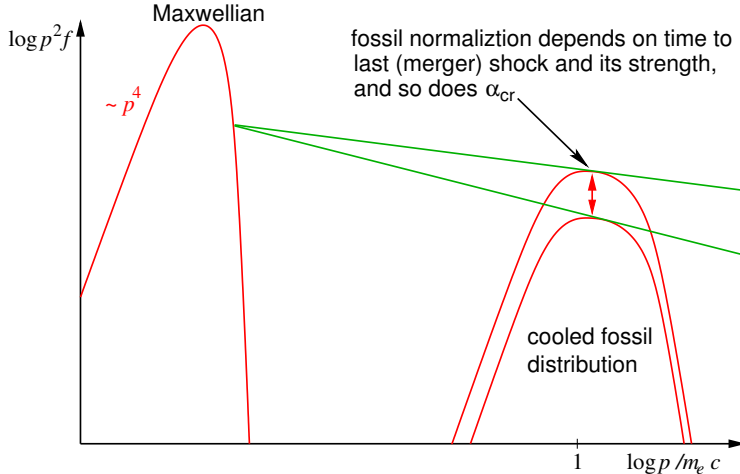
Illuminating radio relics

Re-acceleration of fossil electrons vs. primary acceleration

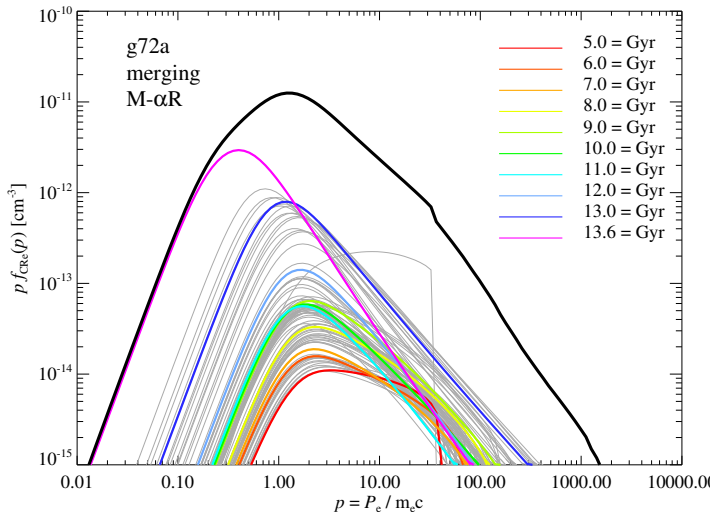


Illuminating radio relics

Re-acceleration of fossil electrons vs. primary acceleration

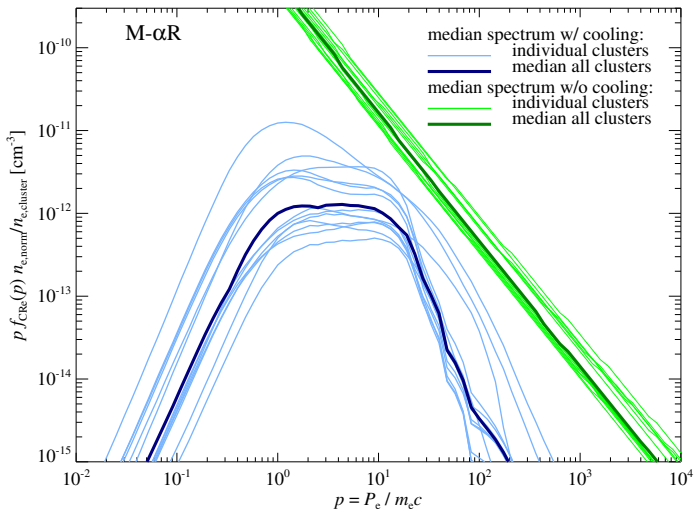


Time evolution of the fossil electron distribution



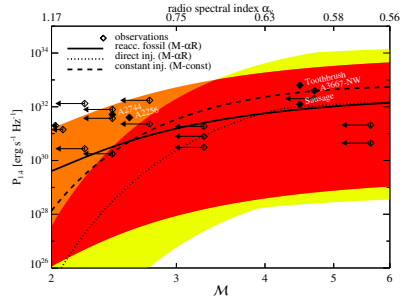
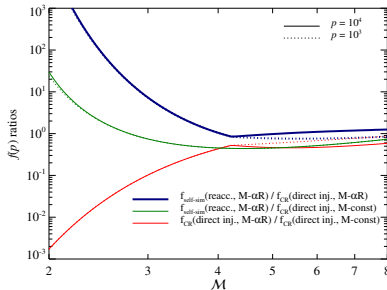
Pinzke, Oh, C.P. (2012)

Fossil CR electron population



Pinzke, Oh, C.P. (2012)

Direct acceleration vs. Fermi-I re-acceleration



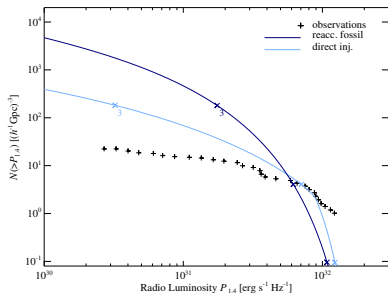
Pinzke, Oh, C.P. (2012)

the bottom line:

- fossil contribution comparable to direct injection at high \mathcal{M}
- fossils dominate at low \mathcal{M}



Radio relics – the future



Pinzke, Oh, C.P. (2012)

→ **the relic luminosity function:**

$$n(> P_{1.4}) = \int dP_{1.4} \frac{dn}{dP_{1.4}}$$

$$\frac{dn}{dP_{1.4}} = \frac{dn}{d\mathcal{M}} \frac{d\mathcal{M}}{dP_{1.4}}$$

depends on the Mach number distribution and the $\mathcal{M} - P_{1.4}$ relation!

bright prospects for LOFAR:

- Fermi-I reacceleration predicts a few 1000 radio relics per Gpc^3
- direct injection predicts a few 100 luminous radio relics



Conclusions on radio halos and relics



- **halos:** producing seed electrons for turbulent reacceleration require modifications to the standard picture:
 - flatter turbulent profile
 - CR streaming
 - high CRe/p injection
- **relics:** fossil electrons could allow radio relics to be seen at low Mach numbers



Literature for the talk

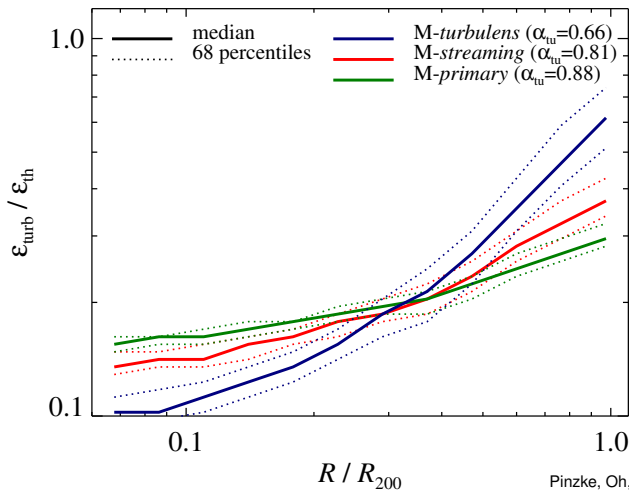
- A. Pinzke, S. P. Oh, C. Pfrommer, *Turbulence and particle acceleration in giant radio halos: the origin of seed electrons*, 2015, submitted, arXiv:1503.07870.
- A. Pinzke, S. P. Oh, C. Pfrommer, *Giant radio relics in galaxy clusters: re-acceleration of fossil relativistic electrons?*, 2013, *MNRAS*, 435, 1061.



Additional slides



Turbulent pressure profiles in our 3 models



Pinzke, Oh, C.P. (2015)

