

Cosmological shock waves in SPH simulations

Exploring cosmic ray feedback

Christoph Pfrommer

Canadian Institute for Theoretical Astrophysics, Toronto

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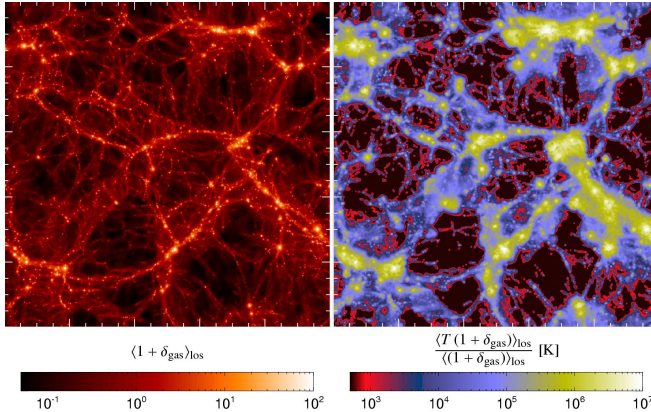
Outline

- 1 **Cosmological shock waves**
 - Cosmic rays in GADGET
 - Mach number finder
 - Cosmological and cluster simulations

- 2 **Cosmic rays in galaxy clusters**
 - Cluster radio halos
 - Energetically preferred CR pressure profiles
 - CR pressure influences Sunyaev-Zel'dovic effect



Cosmic rays in GADGET (Pfrommer, Springel, Enßlin, Jubelgas, 2006, MNRAS)

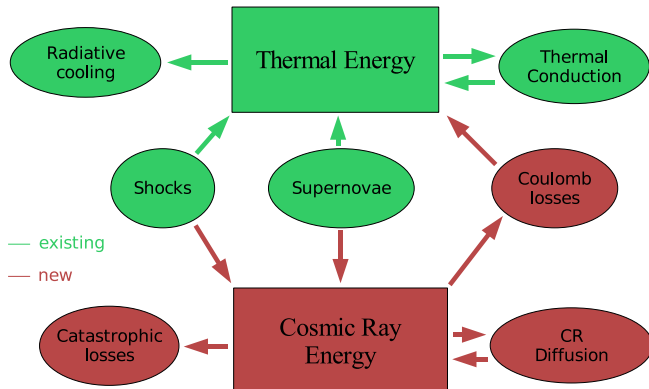


The "cosmic web" today. *Left*: the projected gas density in a cosmological simulation. *Right*: gravitationally heated intracluster medium through cosmological shock waves.



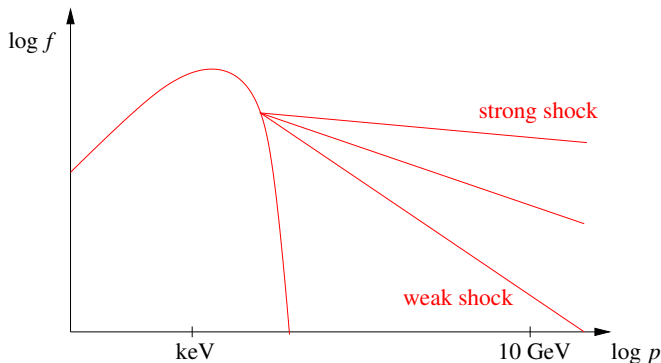
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Cosmic rays in GADGET– flowchart

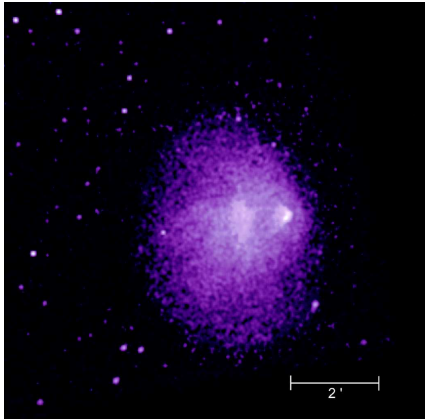


Diffusive shock acceleration – Fermi 1 mechanism

Cosmic rays gain energy $\Delta E/E \propto v_1 - v_2$ through bouncing back and forth the shock front. Accounting for the loss probability $\propto v_2$ of particles leaving the shock downstream leads to power-law CR population.

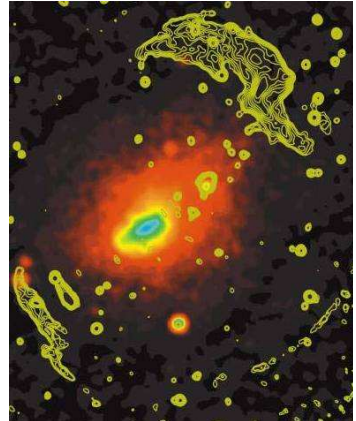


Observations of cluster shock waves



1E 0657-56 ("Bullet cluster")

(NASA/SAO/CXC/M.Markevitch et al.)



Abell 3667

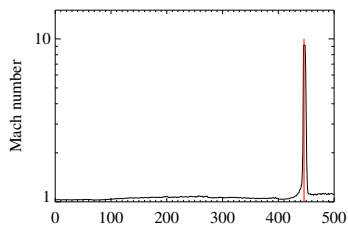
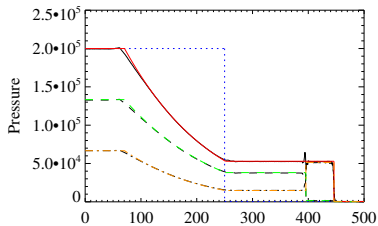
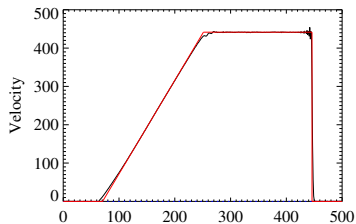
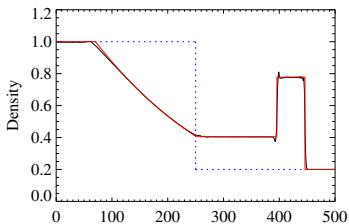
(Radio: Austr.TC Array. X-ray: ROSAT/PSPC.)

Motivation for the Mach number finder

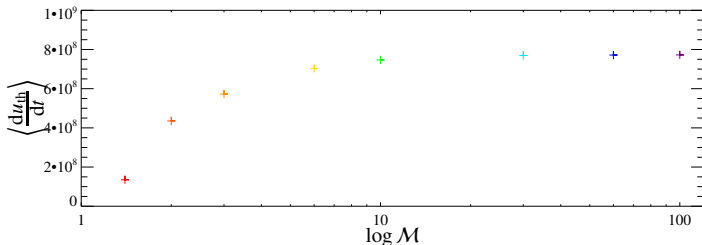
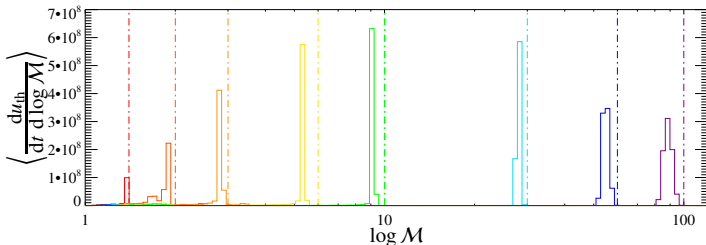
- **cosmological shocks** dissipate gravitational energy into thermal gas energy: where and when is the gas heated, and which shocks are mainly responsible for it?
- **shock waves are tracers** of the large scale structure and contain information about its dynamical history (warm-hot intergalactic medium)
- **shocks accelerate cosmic rays** through diffusive shock acceleration at structure formation shocks: what are the cosmological implications of such a CR component, and does this influence the cosmic thermal history?
- **simulating realistic CR distributions** within galaxy clusters provides detailed predictions for the expected radio synchrotron and γ -ray emission



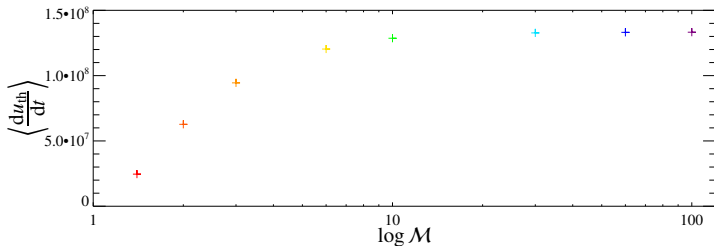
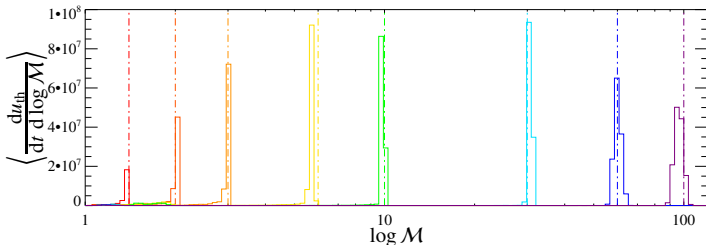
Shock tube (CRs & gas, $\mathcal{M} = 10$): thermodynamics



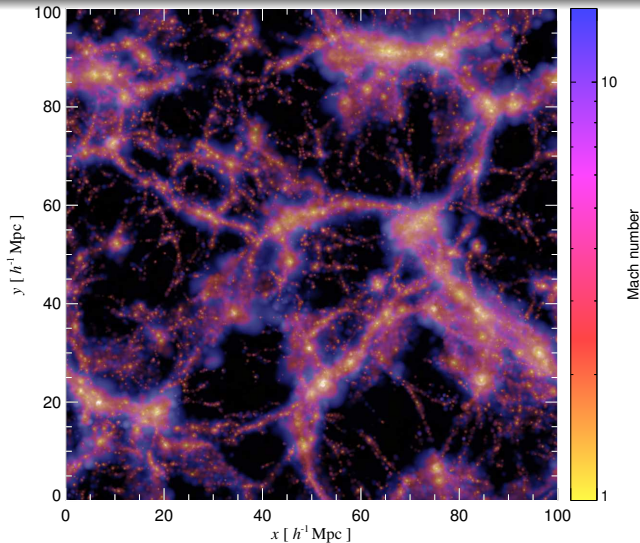
Shock tube (CRs & gas): Mach number statistics



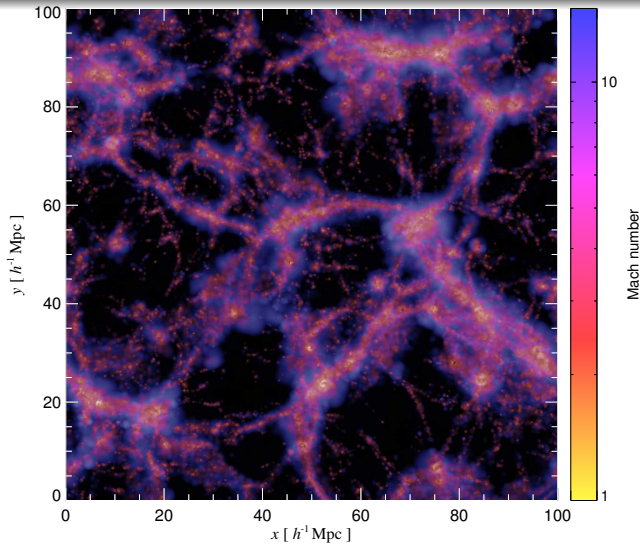
Shock tube (th. gas): Mach number statistics



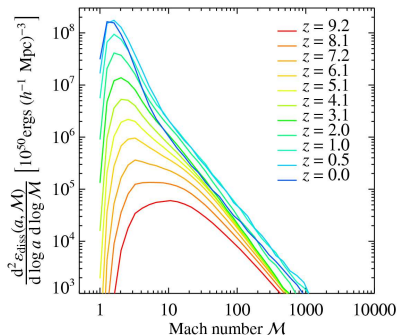
Cosmological Mach numbers: weighted by $\varepsilon_{\text{diss}}$



Cosmological Mach numbers: weighted by ε_{CR}

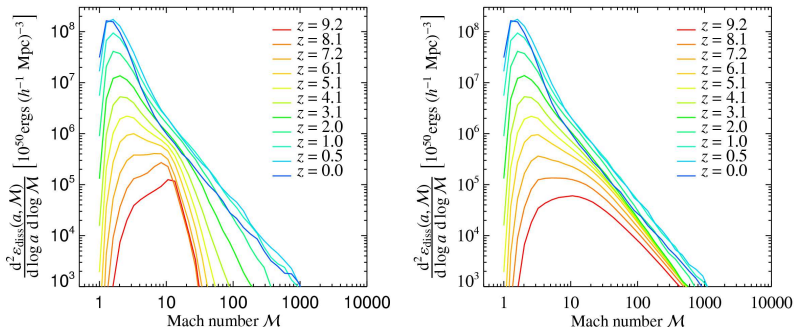


Cosmological Mach number statistics



- more energy is dissipated in weak shocks internal to collapsed structures than in external strong shocks
- more energy is dissipated at later times
- mean Mach number decreases with time

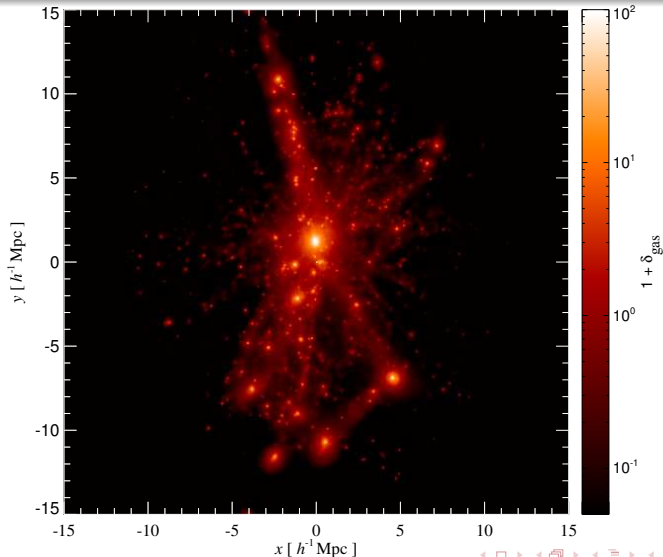
Cosmological statistics: influence of reionization



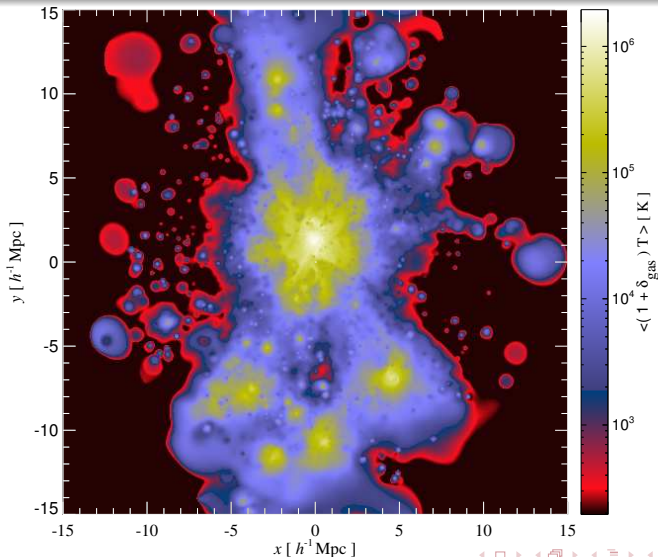
- reionization epoch at $z_{\text{reion}} = 10$ suppresses efficiently strong shocks at $z < z_{\text{reion}}$ due to jump in sound velocity
- cosmological constant causes structure formation to cease



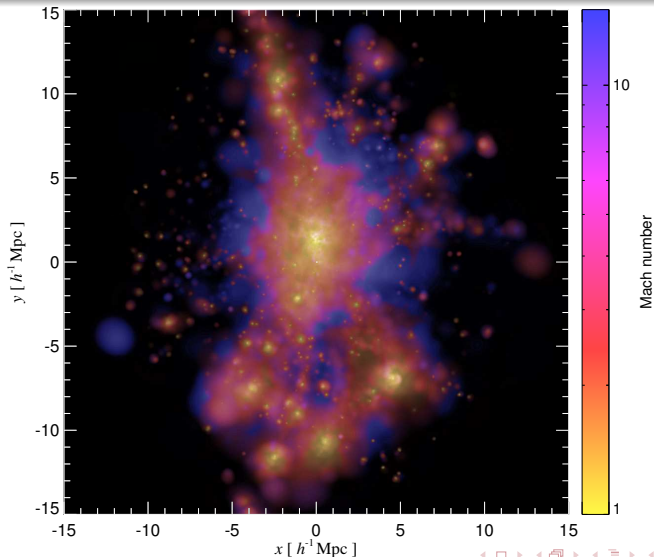
Adiabatic cluster simulation: gas density



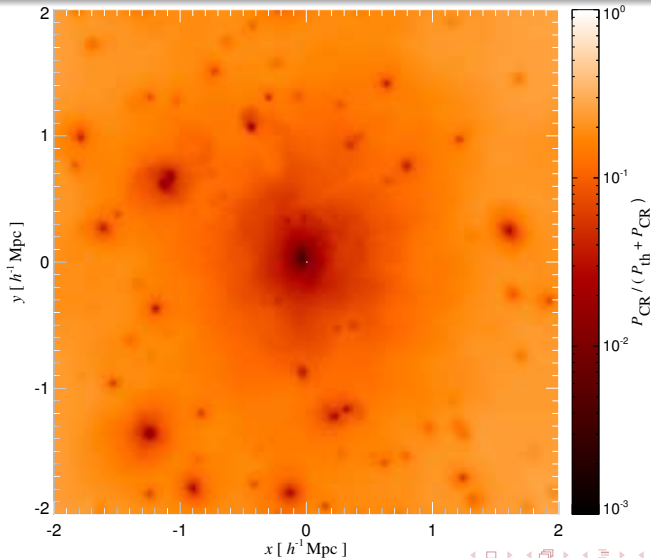
Mass weighted temperature



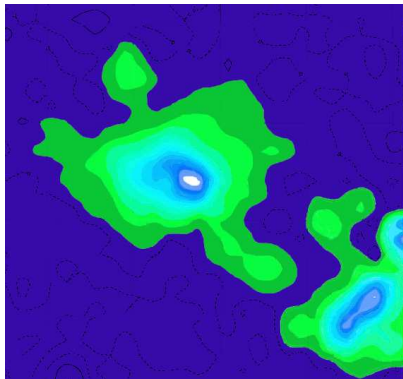
Mach number distribution weighted by ϵ_{diss}



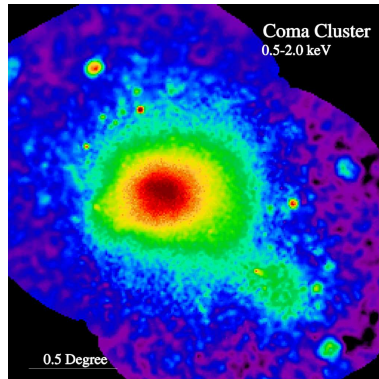
Relative CR pressure $P_{\text{CR}}/P_{\text{total}}$



Radio halos as window for non-equilibrium processes



Coma radio halo, $\nu = 1.4$ GHz,
largest emission diameter ~ 3 Mpc
($2.5^\circ \times 2.0^\circ$, credit: Deiss/Effelsberg)



Coma thermal X-ray emission,
($2.7^\circ \times 2.5^\circ$, credit: ROSAT/MPE/Snowden)

Models for radio synchrotron halos in clusters

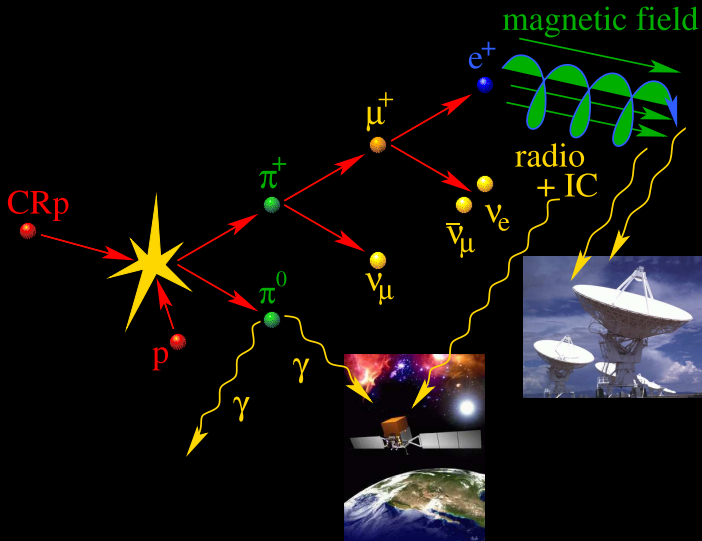
Halo characteristics: smooth unpolarized radio emission at scales of 3 Mpc.

Different CR electron populations:

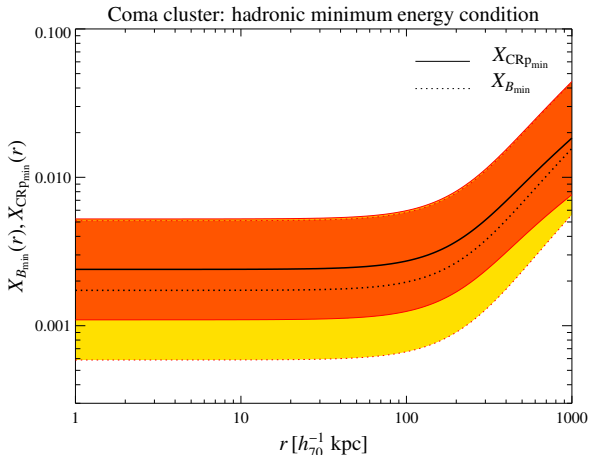
- **Primary accelerated CR electrons**: synchrotron/IC cooling times too short to account for extended diffuse emission
- **Re-accelerated CR electrons** through resonant interaction with turbulent Alfvén waves: possibly too inefficient, no first principle calculations (Jaffe 1977, Schlickeiser 1987, Brunetti 2001)
- **Hadronically produced CR electrons** in inelastic collisions of CR protons with the ambient gas (Dennison 1980, Vestrad 1982, Miniati 2001, Pfrommer 2004)



Hadronic cosmic ray proton interaction



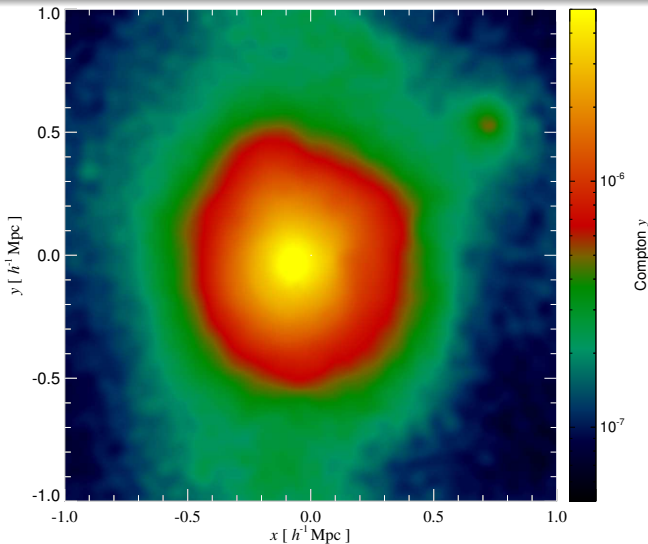
Energetically preferred CR pressure profiles



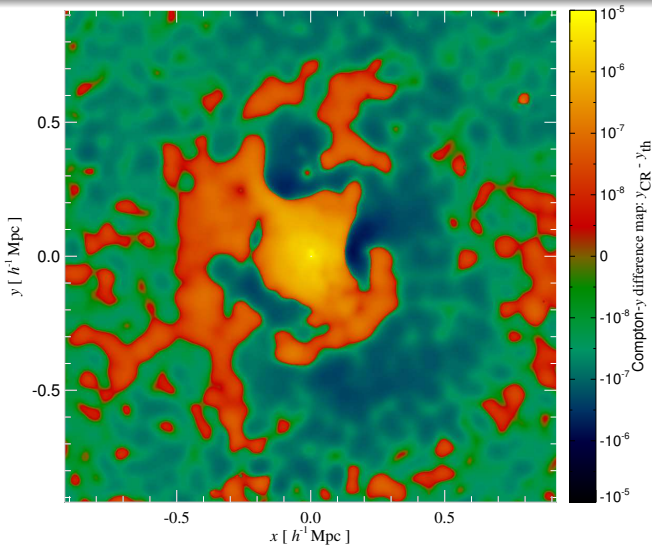
$$X_{CRp}(r) = \frac{\varepsilon_{CRp}}{\varepsilon_{th}}(r), \quad X_B(r) = \frac{\varepsilon_B}{\varepsilon_{th}}(r) \quad \rightarrow \quad B_{\text{Coma, min}}(0) = 2.4^{+1.7}_{-1.0} \mu\text{G}$$



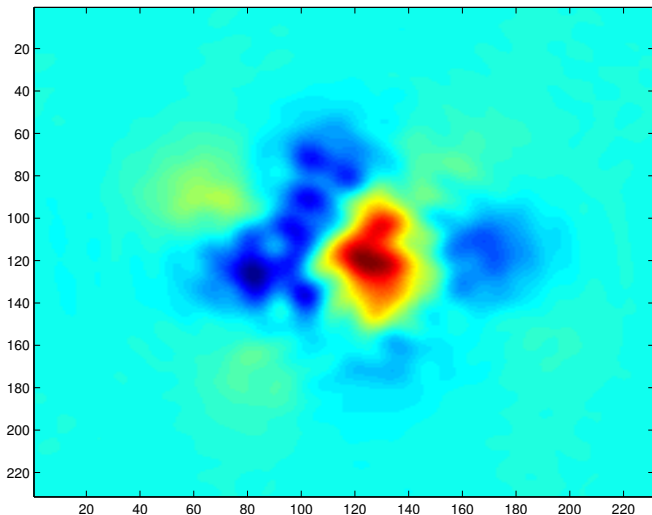
Compton y parameter in radiative cluster simulation



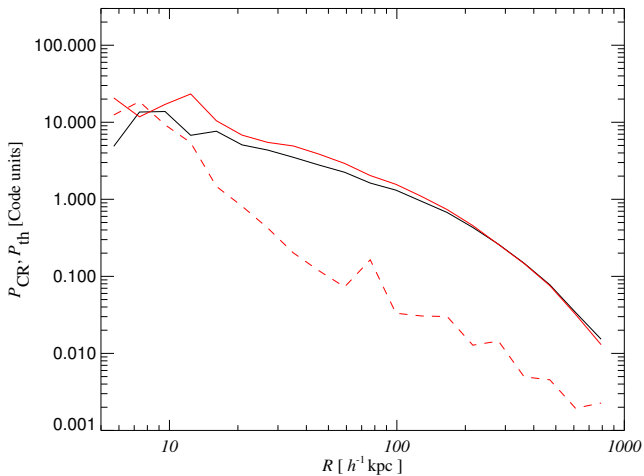
Compton y difference map: $y_{\text{CR}} - y_{\text{th}}$



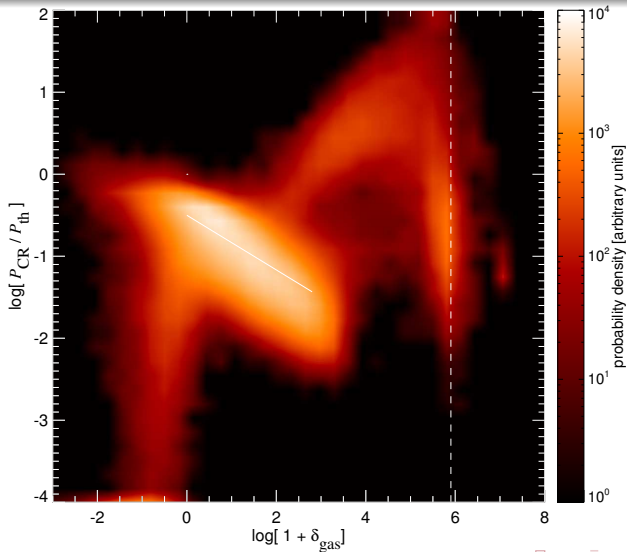
Simulated CBI observation of $y_{\text{CR}} - y_{\text{th}}$ (with Sievers & Bond)



Pressure profiles with and without CRs



Phase-space diagram of radiative cluster simulation



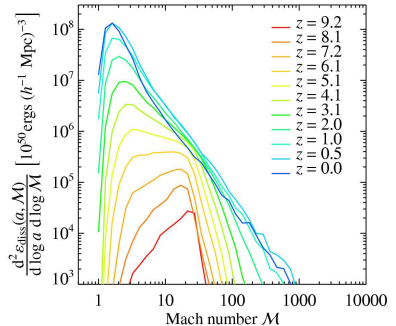
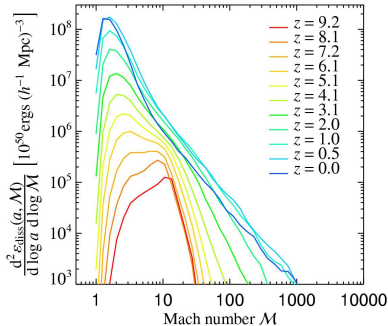
Summary

- Understanding **non-thermal processes** is crucial for using clusters as cosmological probes (high- z scaling relations).
- **Radio halos** might be of hadronic origin as our simulations suggests \rightarrow tracer of structure formation
- **Dynamical CR feedback** influences Sunyaev-Zel'dovic effect
- Outlook
 - **Galaxy evolution**: influence on energetic feedback, star formation, and galactic winds
 - Huge potential and predictive power of **cosmological CR simulations/Mach number finder** \rightarrow provides detailed γ -ray/radio emission maps



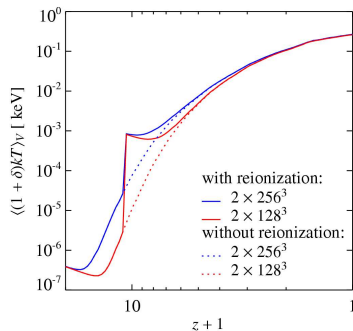
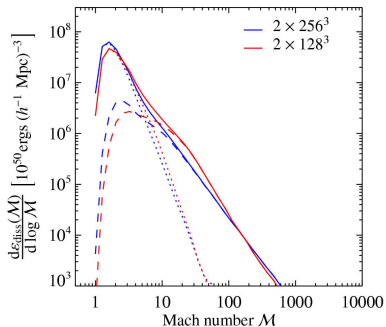
Cosmological statistics: resolution study

Differential distributions: 2×256^3 versus 2×128^3



- more energy is dissipated at later times
- mean Mach number decreases with time
- differential Mach number distributions are converged for $z < 3$

Cosmological statistics: resolution study



- in higher resolution simulations structure forms earlier
- more energy is dissipated in shocks internal to collapsed structures than in external shocks of pristine gas
- integrated Mach number distribution converged

Idea of the Mach number finder in SPH

- SPH shock is broadened to a scale of the order of the smoothing length h , i.e. $f_h h$, and $f_h \sim 2$
- approximate instantaneous particle velocity by pre-shock velocity (denoted by $v_1 = \mathcal{M}_1 c_1$)

Using the **entropy conserving formalism** of Springel & Hernquist 2002 ($A(s) = P\rho^{-\gamma}$ is the entropic function):

$$\frac{A_2}{A_1} = \frac{A_1 + dA_1}{A_1} = 1 + \frac{f_h h}{\mathcal{M}_1 c_1 A_1} \frac{dA_1}{dt} = \frac{P_2}{P_1} \left(\frac{\rho_1}{\rho_2} \right)^\gamma$$

$$\frac{\rho_2}{\rho_1} = \frac{(\gamma + 1)\mathcal{M}_1^2}{(\gamma - 1)\mathcal{M}_1^2 + 2}$$

$$\frac{P_2}{P_1} = \frac{2\gamma\mathcal{M}_1^2 - (\gamma - 1)}{\gamma + 1}$$

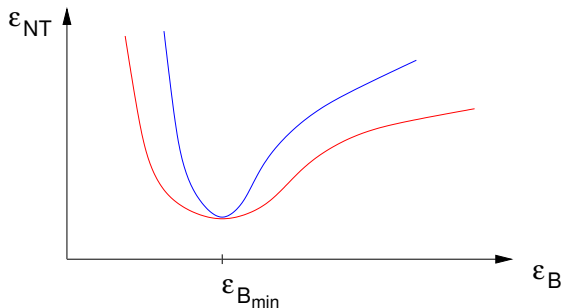


Minimum energy criterion (MEC): the idea

- What is the energetically least expensive distribution of non-thermal energy density ε_{NT} given the observed synchrotron emissivity?

- $\varepsilon_{\text{NT}} = \varepsilon_B + \varepsilon_{\text{CRp}} + \varepsilon_{\text{CRe}}$

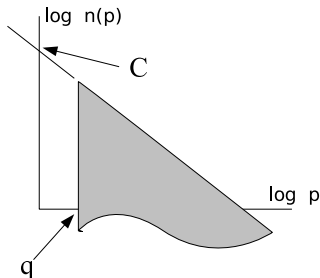
→ minimum energy criterion: $\left. \frac{\partial \varepsilon_{\text{NT}}}{\partial \varepsilon_B} \right|_{j_\nu} \stackrel{!}{=} 0$



defining tolerance levels: deviation from minimum by one e-fold

Philosophy and description

- CRs are coupled to the thermal gas by magnetic fields.
- We assume a single power-law CR spectrum: momentum cutoff q , normalization C , spectral index α (constant).
→ determines CR energy density and pressure uniquely



The CR spectrum can be expressed by three adiabatic invariants, which scale only with the gas density. Non-adiabatic processes are mapped into changes of the adiabatic constants using mass, energy and momentum conservation.

