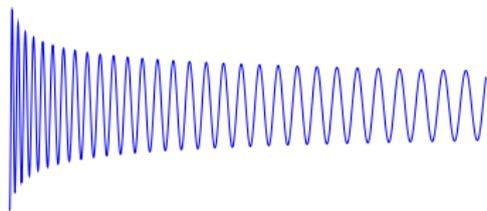
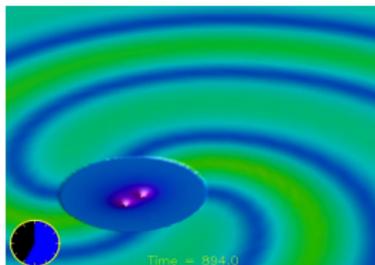


Binary Black Hole Simulations

Harald Pfeiffer

California Institute of Technology

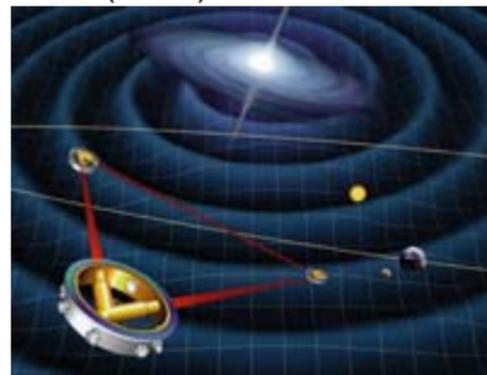
Mike Boyle, Duncan Brown, Lee Lindblom, Geoffrey Lovelace,
Larry Kidder, Mark Scheel, Saul Teukolsky



APS April Meeting, Jacksonville, Apr 17, 2007

Gravitational wave detectors

LISA (201x)



LIGO (2 sites)



GEO 600

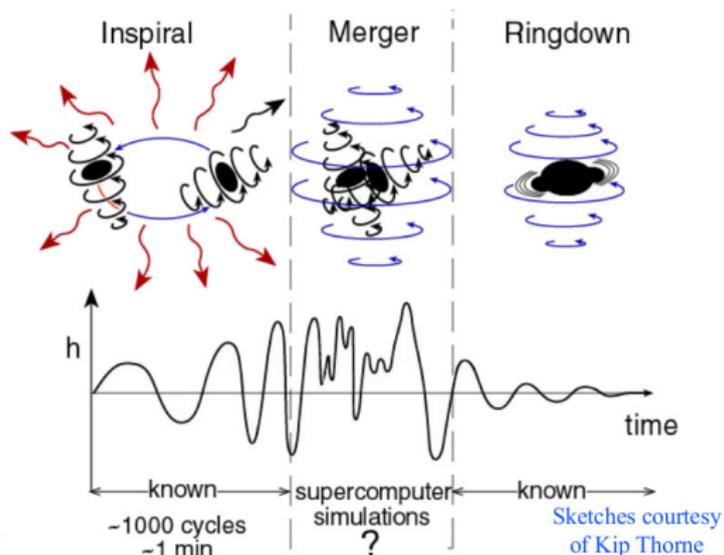


VIRGO



- Among prime targets: Binary black hole systems

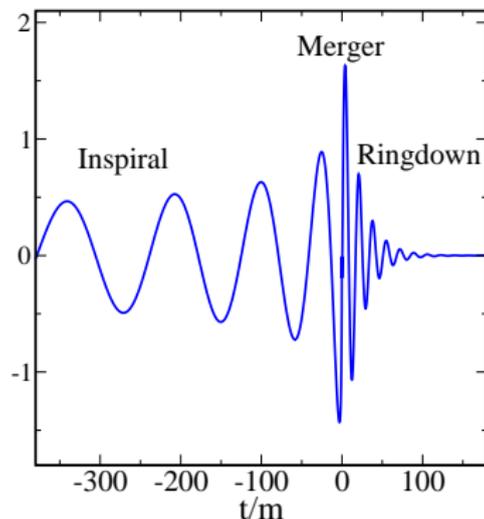
Stages of binary black hole evolution



- Knowledge of waveform allows to
 - ▶ Enhance detector sensitivity
 - ▶ Test general relativity
 - ▶ Extract information about source (→ astrophysics)

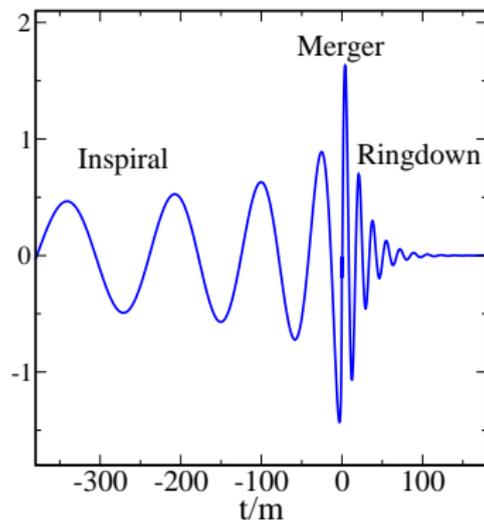
Tools for computing the waveform

- **Inspiral**
 - $v \ll c$: perturbative expansion in v/c (post-Newtonian expansion)
 - v/c large: Numerical relativity
- **Merger**
 - Numerical relativity
- **Ringdown**
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 - BH perturbation theory
 - Numerical relativity
- **Tasks for Numerical relativity:**
 - simulate “late” inspiral and merger.
 - determine what “late” means.



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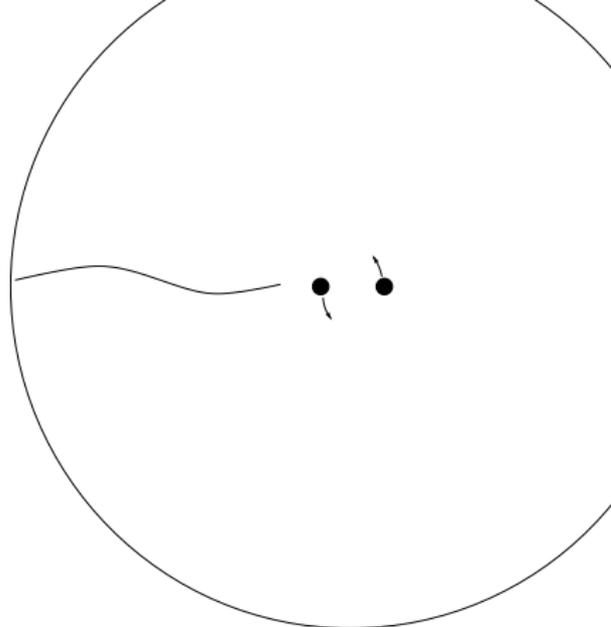
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Overview

Problem characteristics

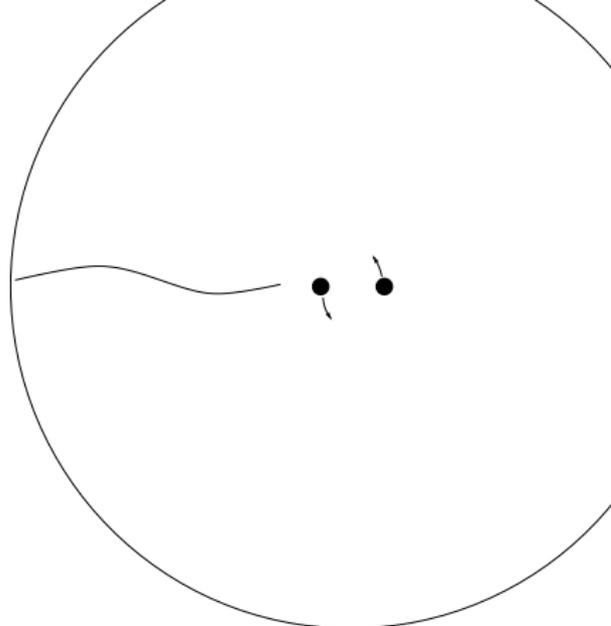
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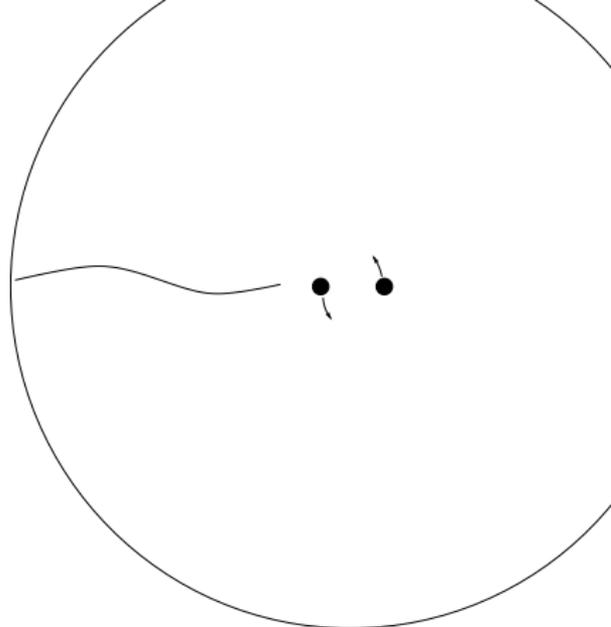
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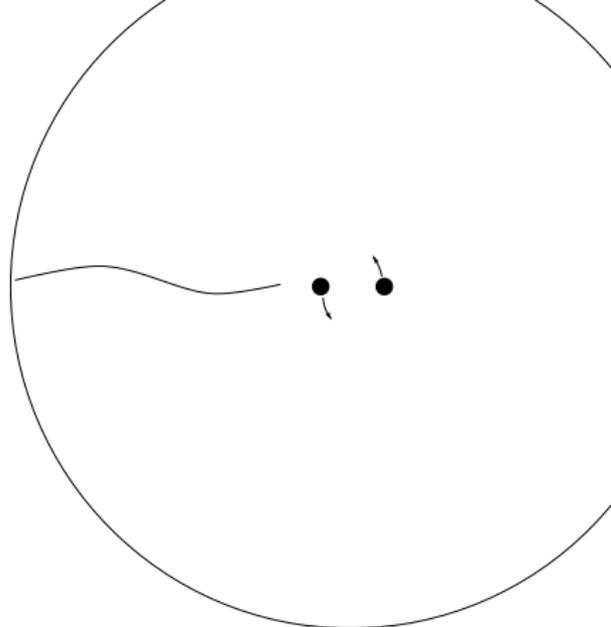
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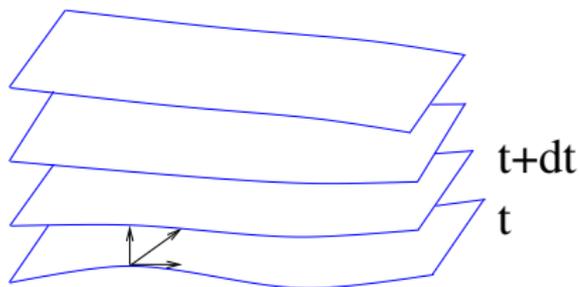
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 - ▶ Cornell/Caltech
 - ▶ Impressive long inspiral simulations
 - ▶ Merger difficult

Solving Einstein's equations – basic idea

- Task: Find space-time metric g_{ab} such that $R_{ab}[g_{ab}] = 0$

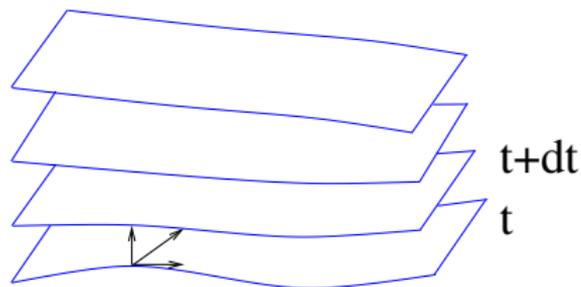
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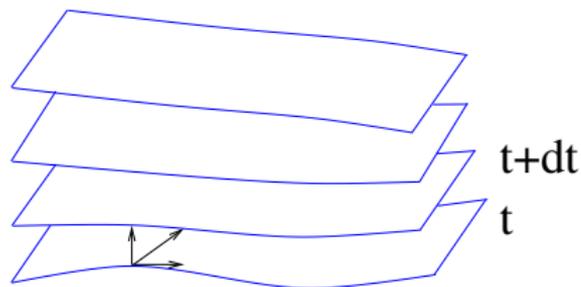
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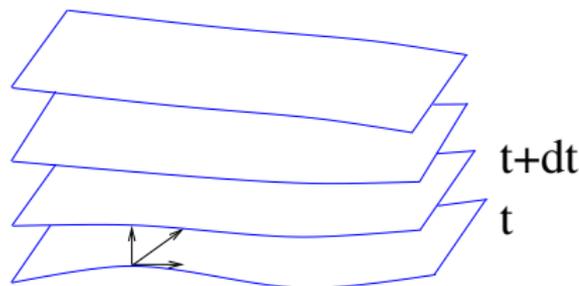
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cf. Maxwell equations

$$\partial_t \vec{E} = \nabla \times \vec{B}$$

$$\partial_t \vec{B} = -\nabla \times \vec{E}$$

$$\nabla \cdot \vec{E} = 0$$

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$$0 = R_{ab}[g_{ab}] = -\frac{1}{2}\square g_{ab} + \nabla_{(a}\Gamma_{b)} + \text{lower order terms} \quad \Gamma_a = -g_{ab}\square x^b.$$

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Constraint damping (Gundlach, et al., Pretorius, 2005)

$$0 = -\frac{1}{2}\square g_{ab} + \nabla_{(a}C_{b)} + \gamma \left[t_{(a}C_{b)} - \frac{1}{2}g_{ab}t^c C_c \right] + \text{l. o.}$$

$$\partial_t C_a \sim -\gamma C_a.$$

Spectral Evolution code

- Rewrite as **first order symmetric hyperbolic system** (Lindblom et al. 2005)

$$\partial_t u + A(u)^k \partial_k u = F(u).$$

- Approximate solution by **truncated series**

$$u(x, t) \approx u^{(N)}(x, t) \equiv \sum_{k=0}^{N-1} \tilde{u}_k(t) \Phi_k(x),$$

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- Evolve $u(x_j)$ by **method of lines**

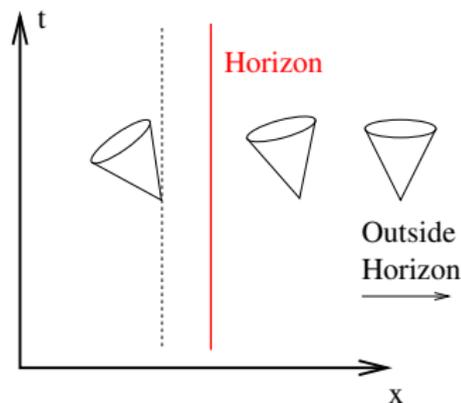
$$\partial_t u(x_j) = \left[F - A(u)^k \partial_k u \right]_{x=x_j}.$$

Black hole singularity excision

- Boundary conditions
 - ▶ Find characteristic fields & speeds
 - ▶ Impose BCs on incoming fields only

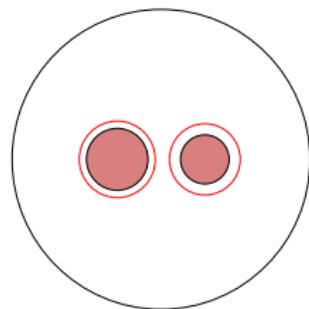
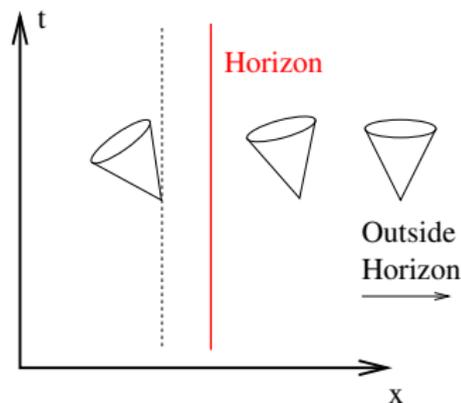
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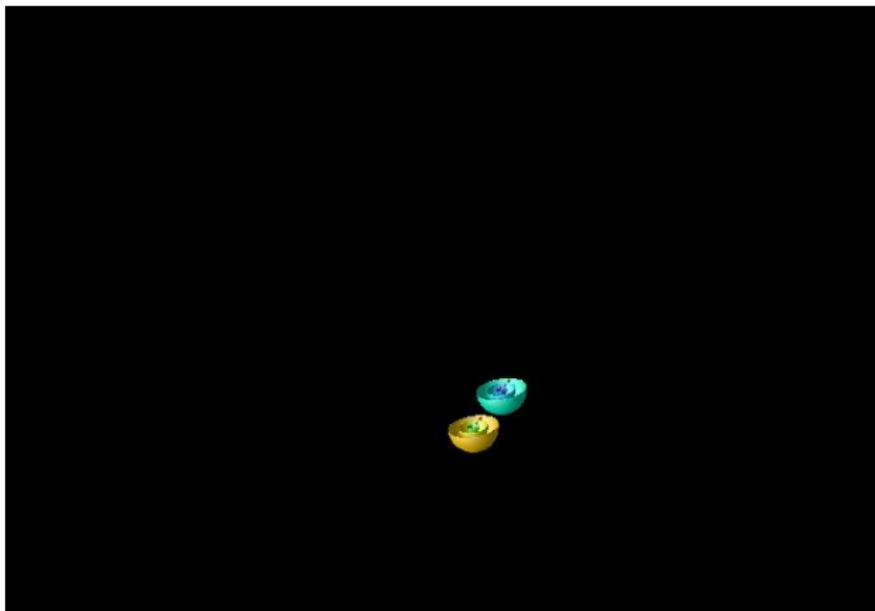


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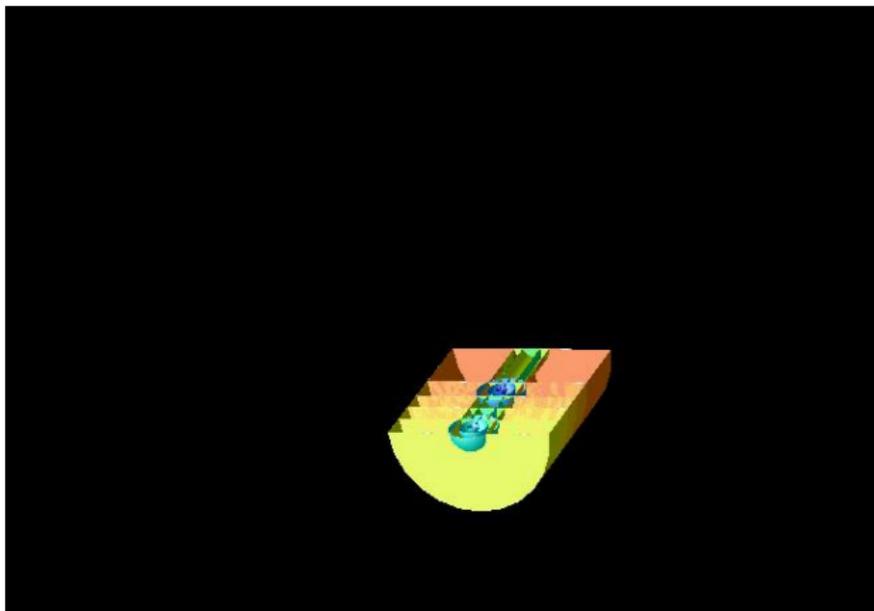
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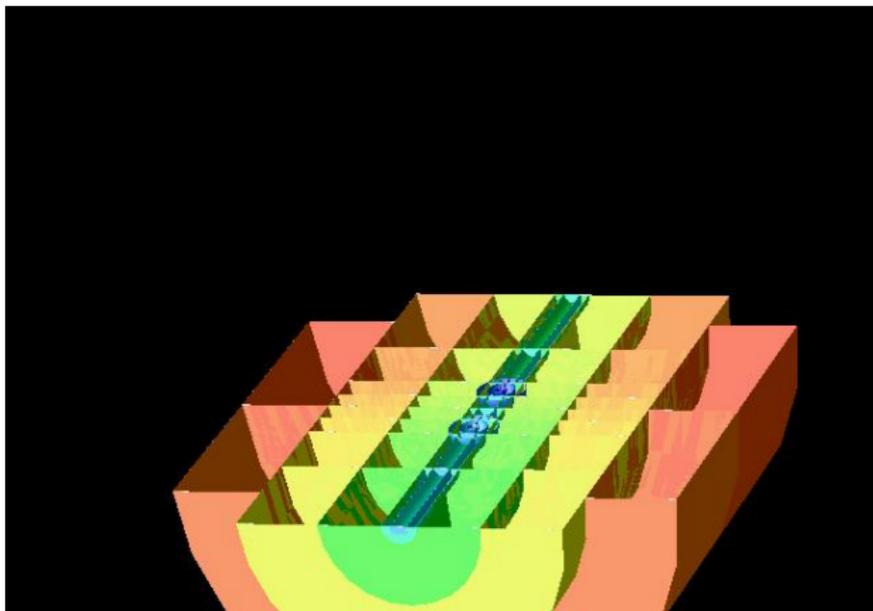
Domain-decomposition



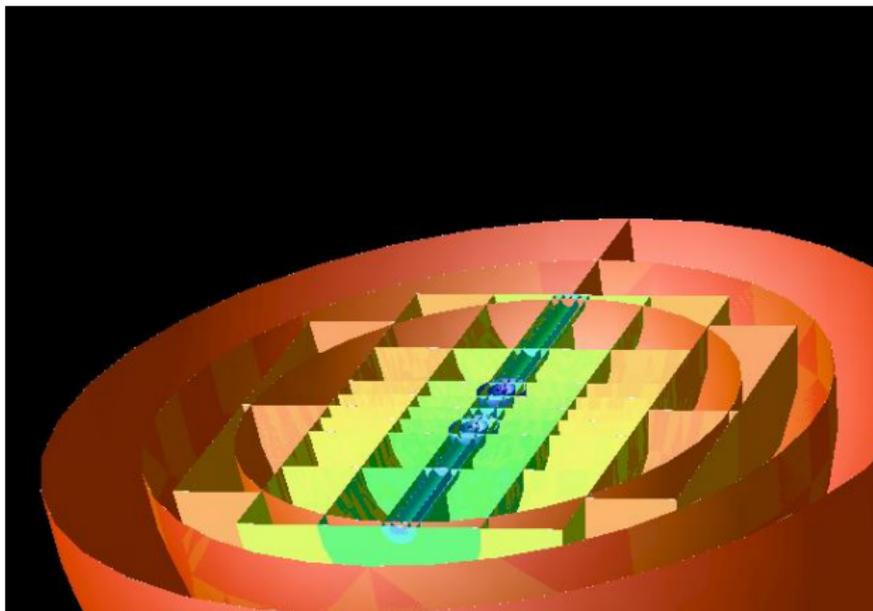
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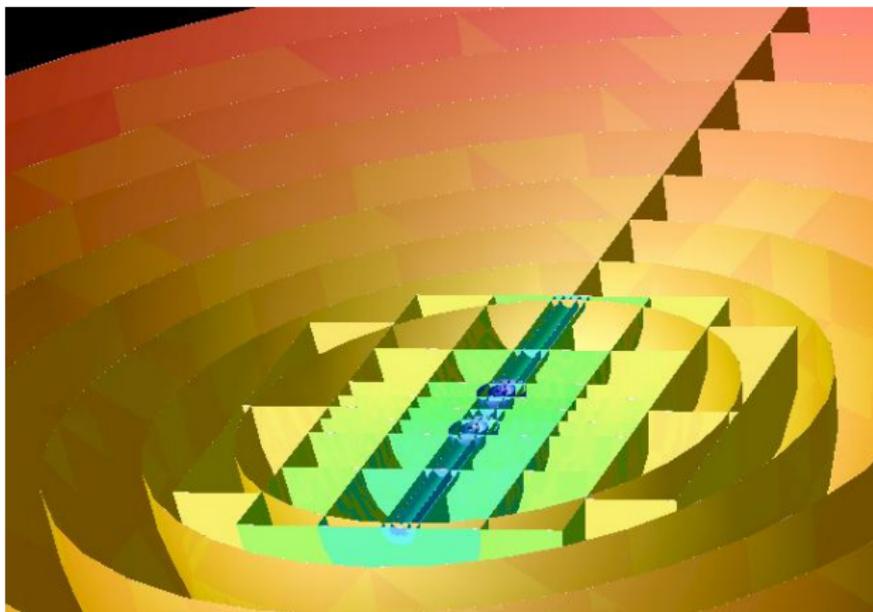
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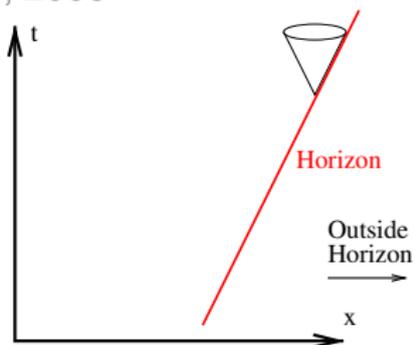
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- Outer shells have fixed angular resolution. Cost increases **linearly** with radius of outer boundary.

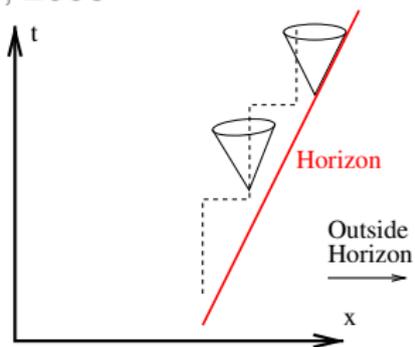
Moving black holes – Dual frame method

Scheel et al., 2006



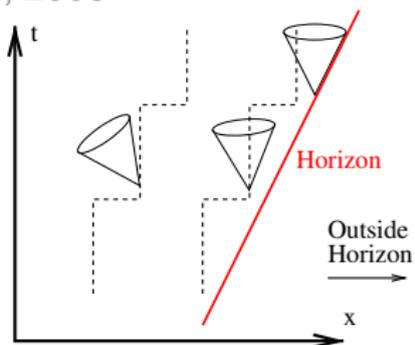
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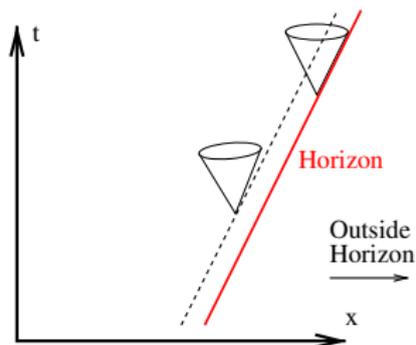
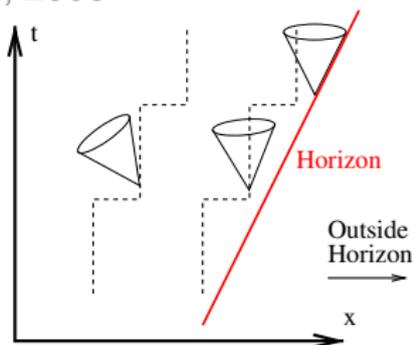
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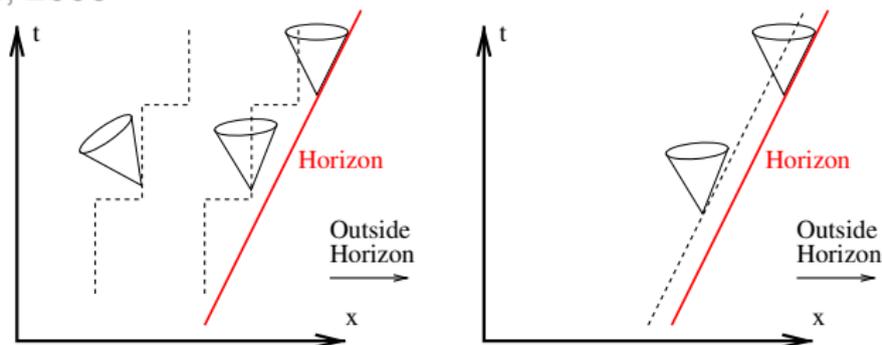
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- Map between “moving” and “inertial” coordinates:

$$\vec{x}_{\text{inertial}} = a(t)R(t)\vec{x}_{\text{moving}}$$

$R(t)$ rotation, $a(t)$ radial scaling.

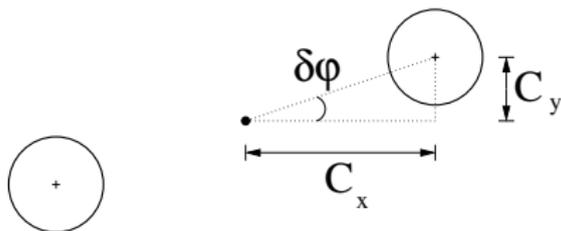
Dynamic feedback control

Example: Control of ϕ

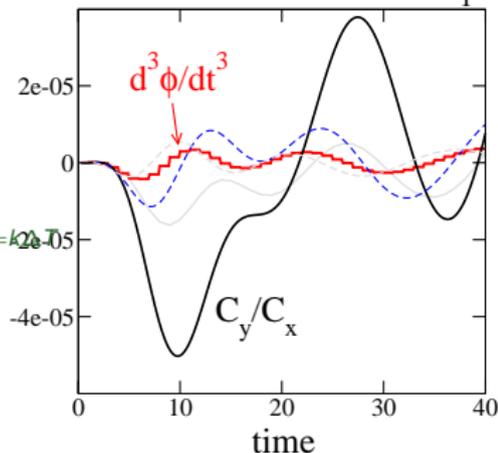
- Measure BH location $C_x(t), C_y(t)$
- Update ϕ periodically s.t. $C_y \rightarrow 0$:

$$\frac{d^3\phi}{dt^3}(t) = \ddot{\phi}_k, \quad k\Delta T \leq t < (k+1)\Delta T$$

$$\ddot{\phi}_k = - \left[\frac{1}{\tau^3} \frac{C_y}{C_x} + \frac{3}{\tau^2} \frac{d}{dt} \left(\frac{C_y}{C_x} \right) + \frac{3}{\tau} \frac{d^2}{dt^2} \left(\frac{C_y}{C_x} \right) \right]_{t=k\Delta T}$$



ϕ -Control in action ($\tau=5M_1$)



Outer boundary conditions

- Must prevent influx of constraint violations

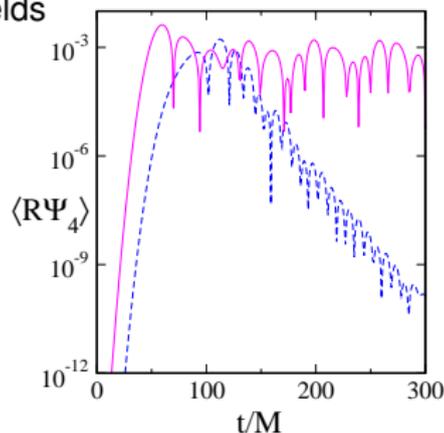
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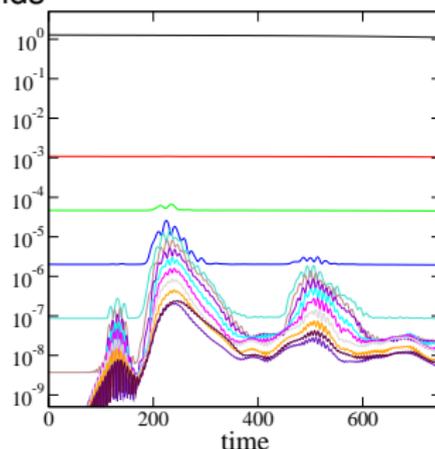
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Lindblom et al., 2006



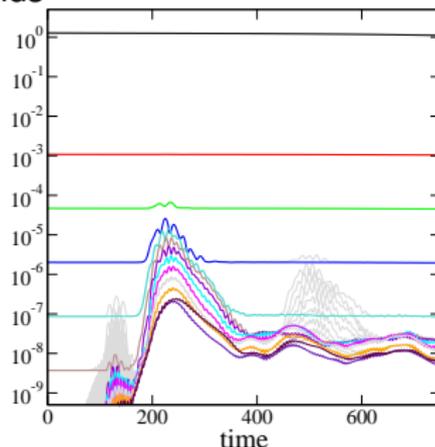
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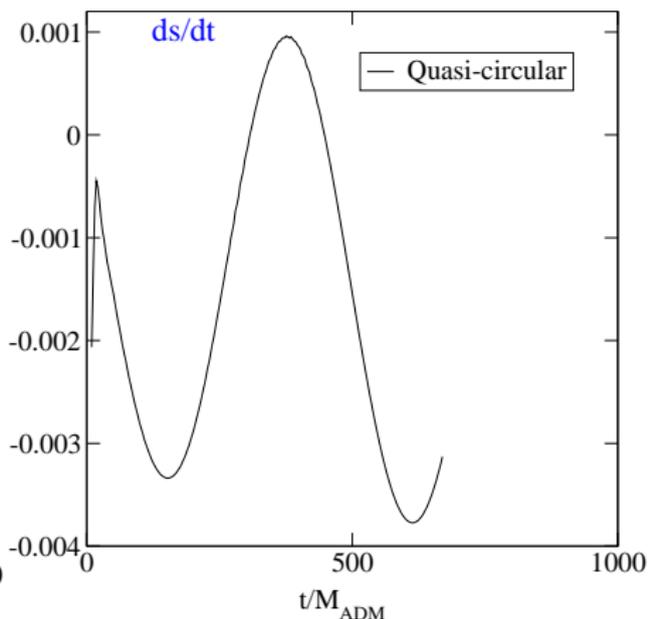
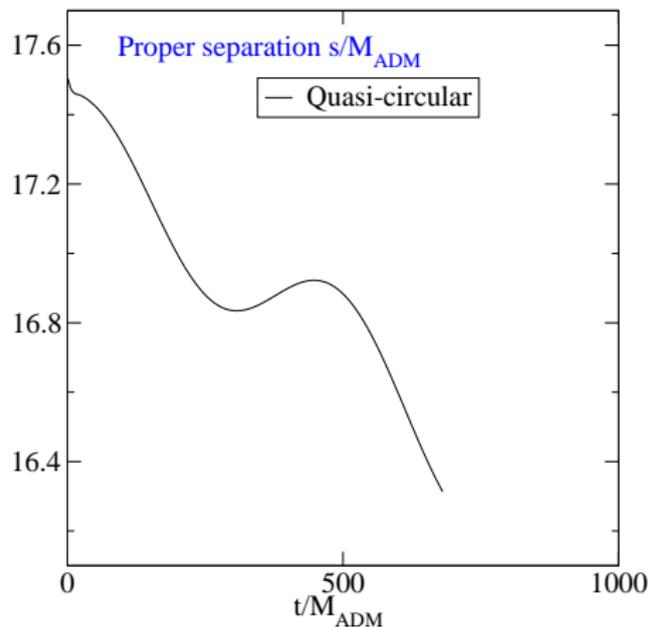
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Improve initial data through evolutions

HP et al., 2007

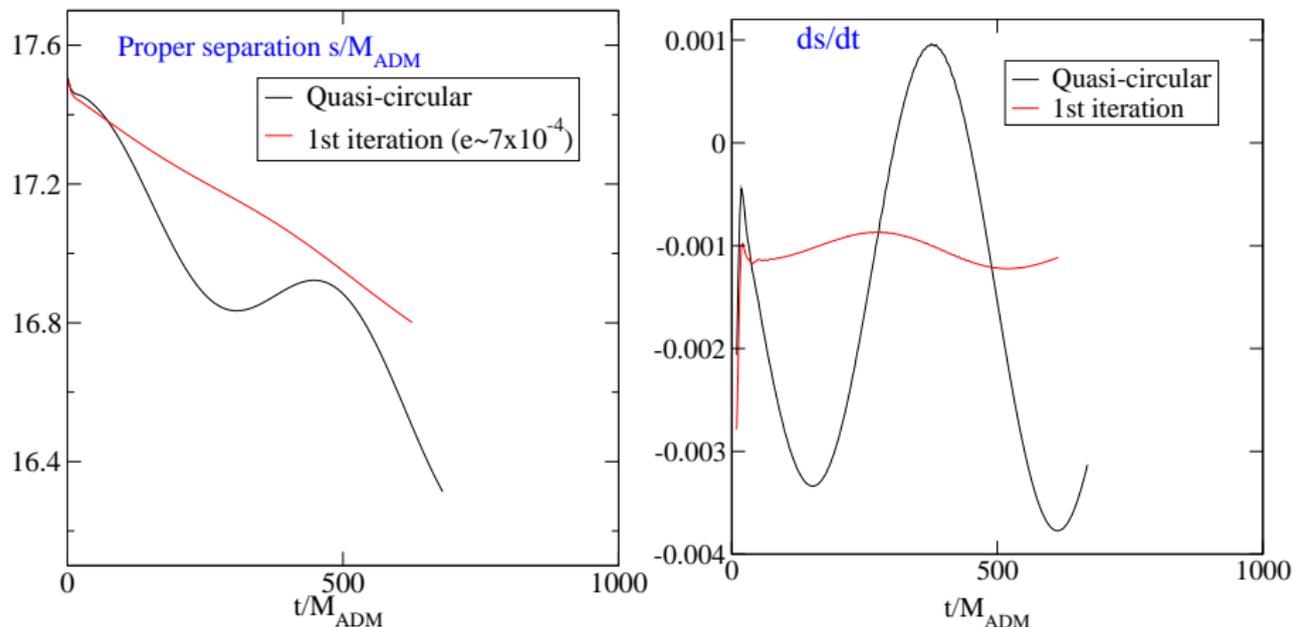
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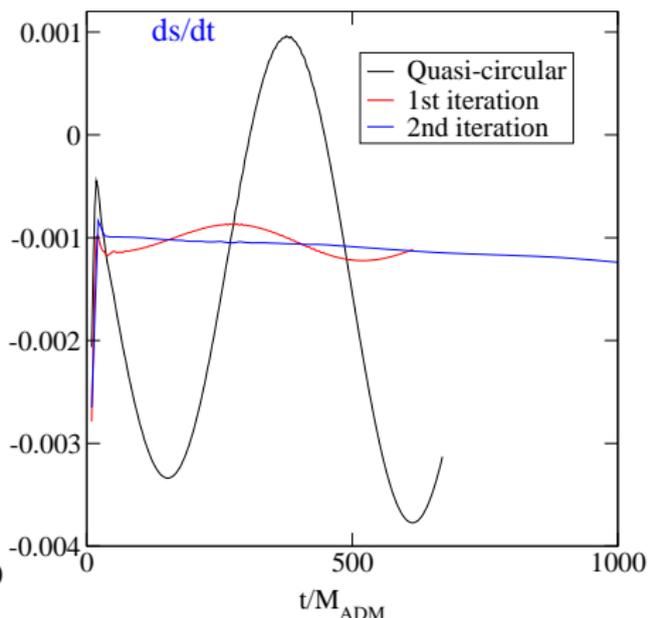
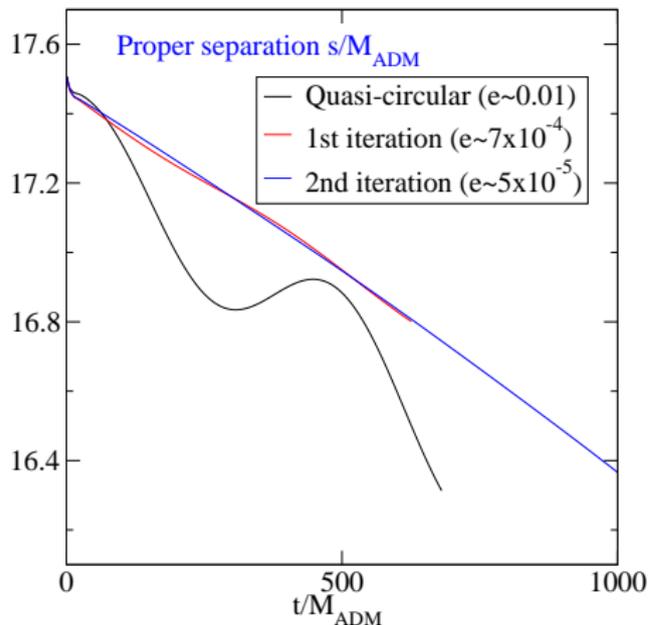
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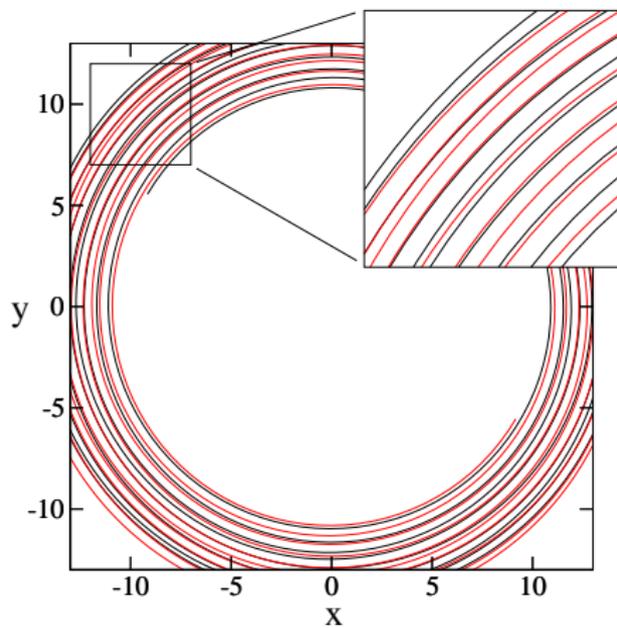
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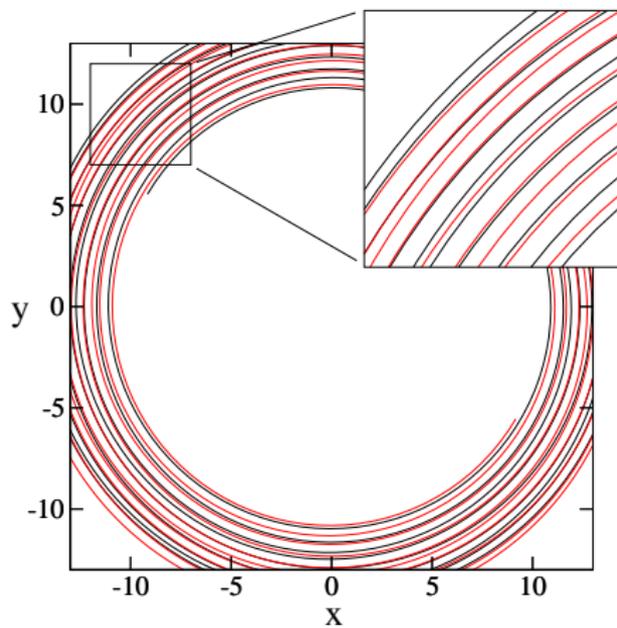
Orbital trajectory

Quasi-circular

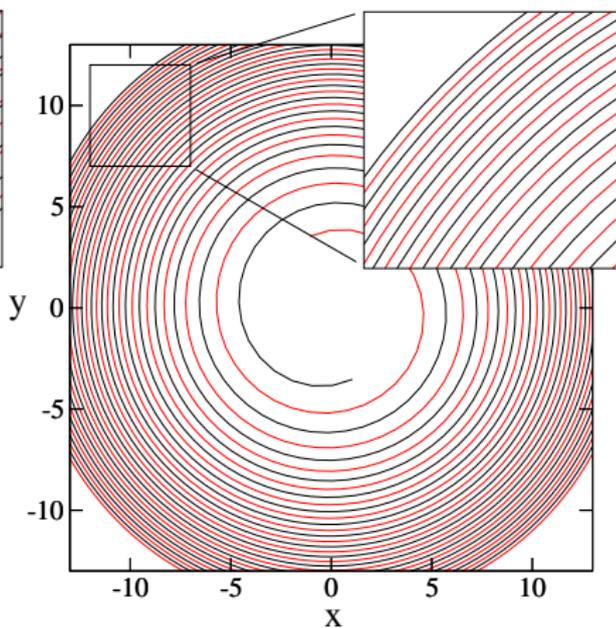


Orbital trajectory

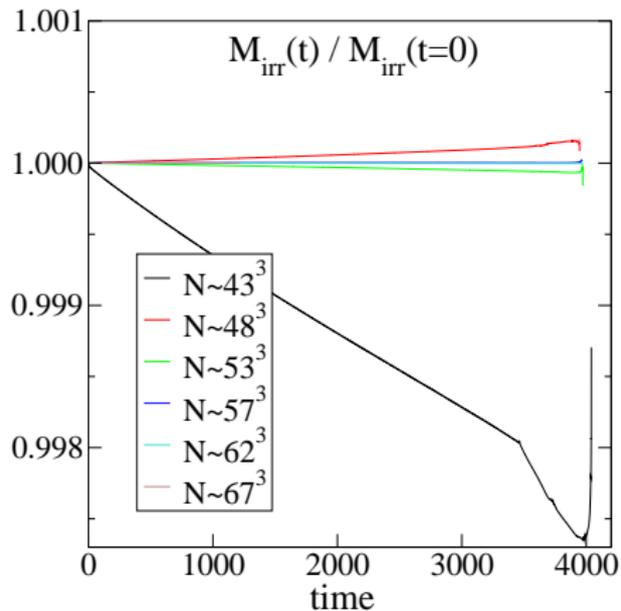
Quasi-circular



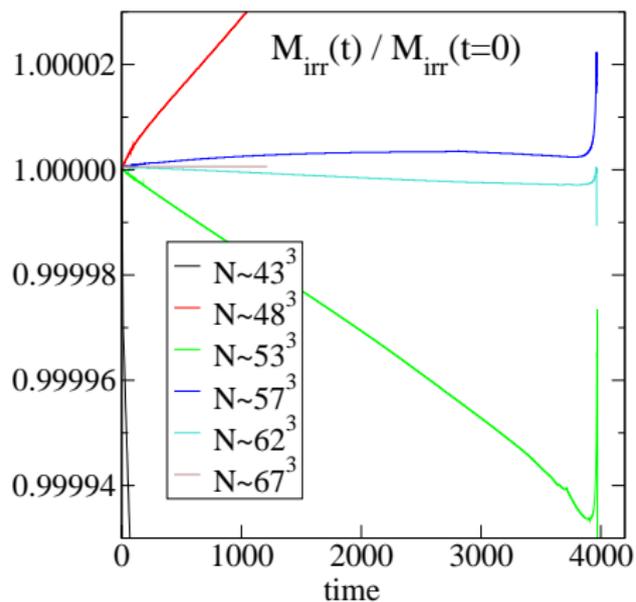
Eccentricity reduced



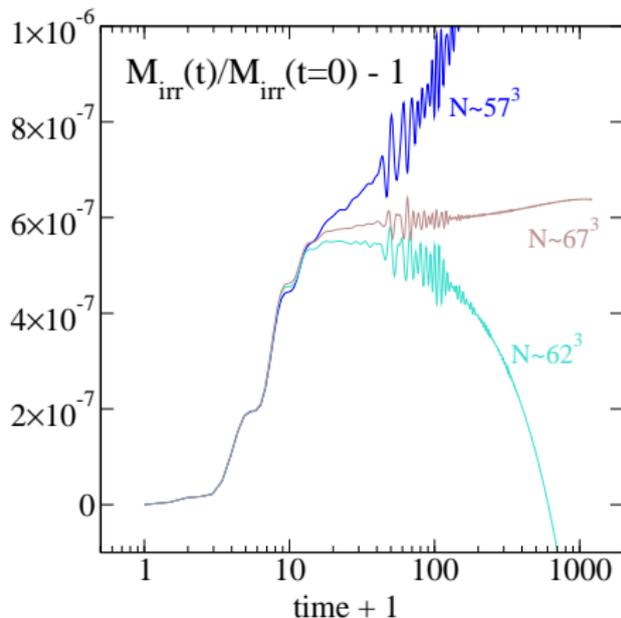
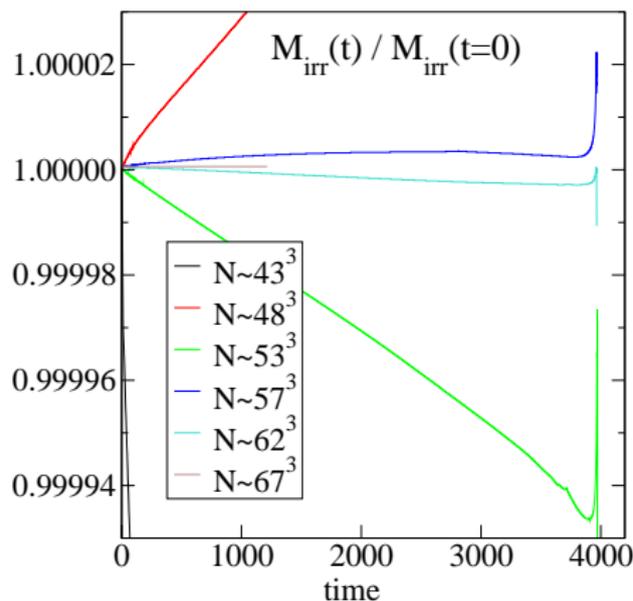
Irreducible Mass



Irreducible Mass

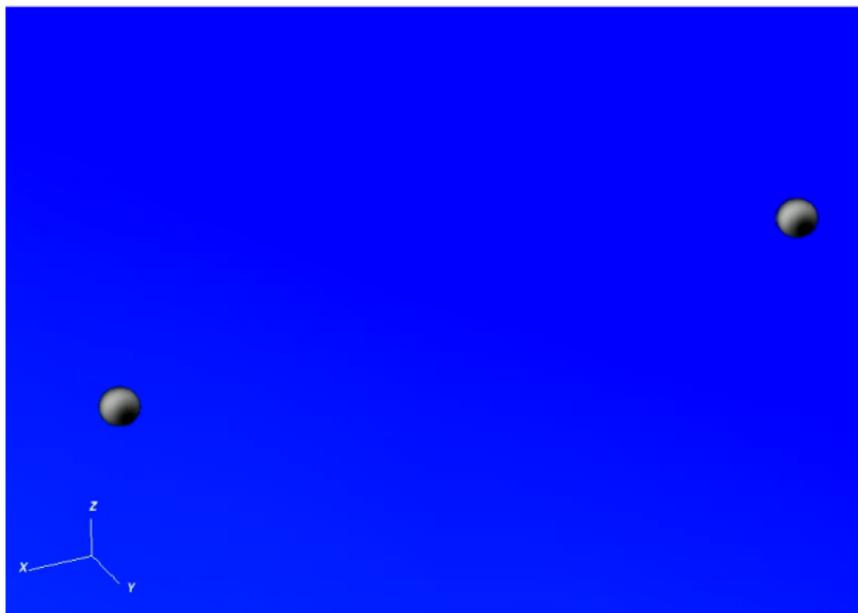


Irreducible Mass

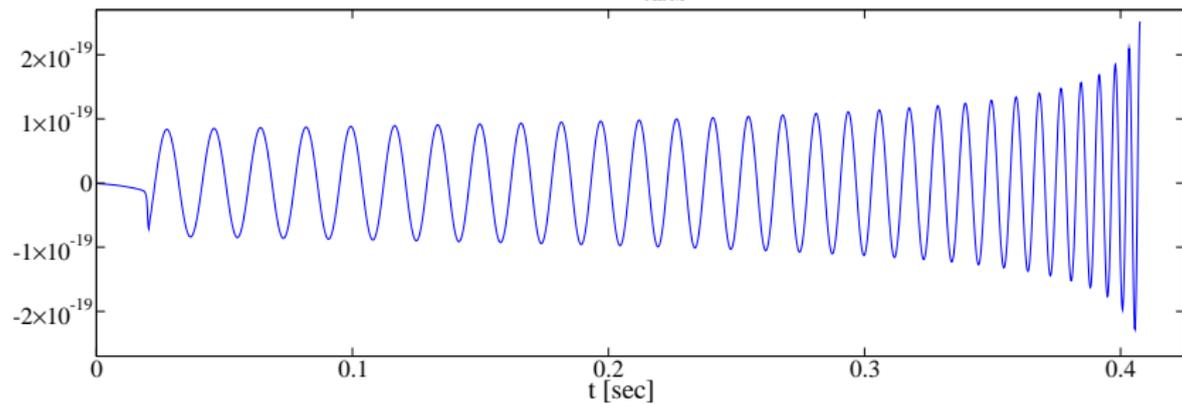
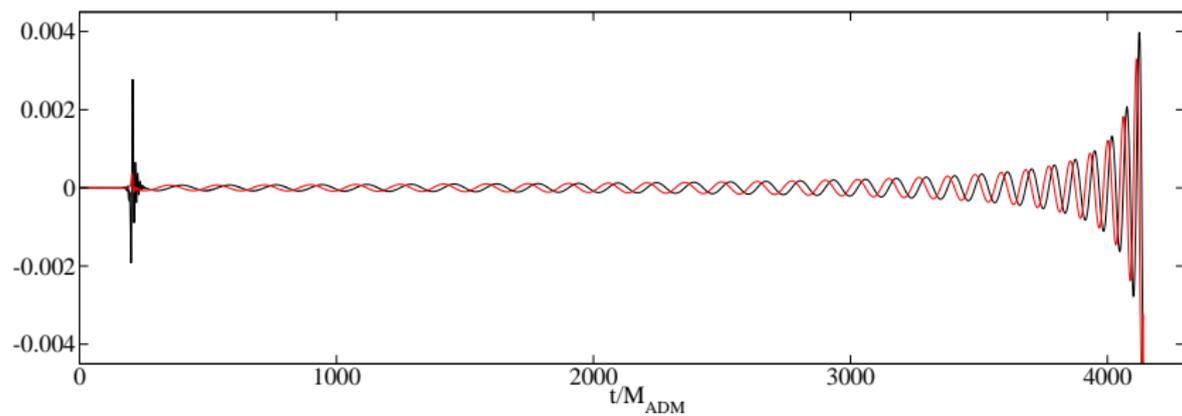


- M_{irr} initially increases by $6 \cdot 10^{-7}$ – “Junk” radiation falling into BH
- $|\delta M_{\text{irr}}| < 10^{-7}$ in the next 3.5 orbits – Limit on tidal heating

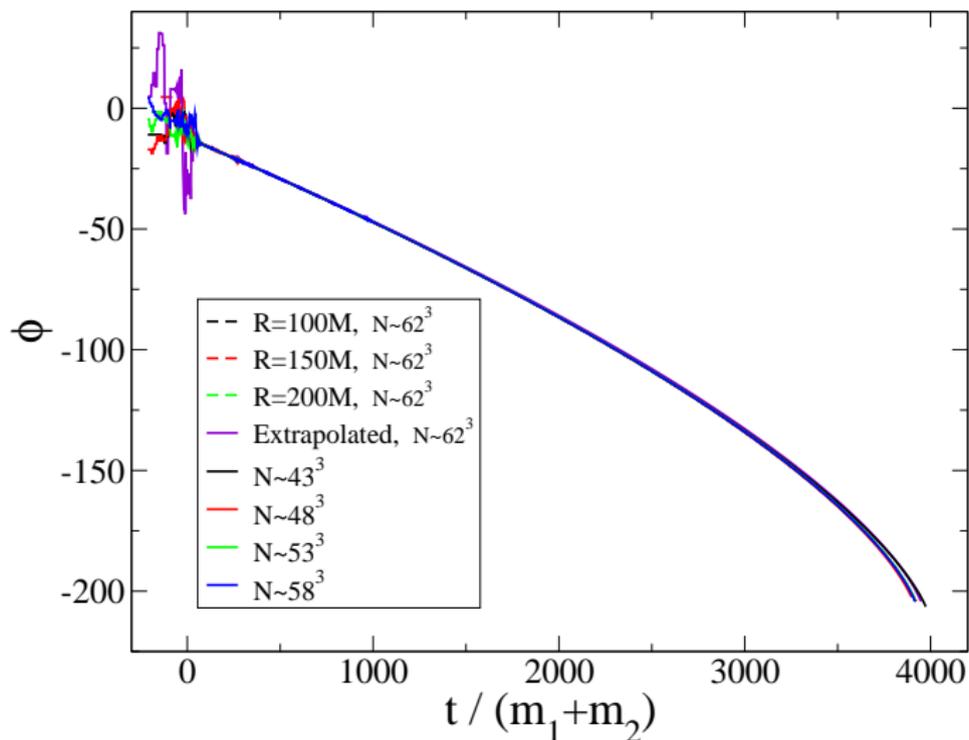
Movie



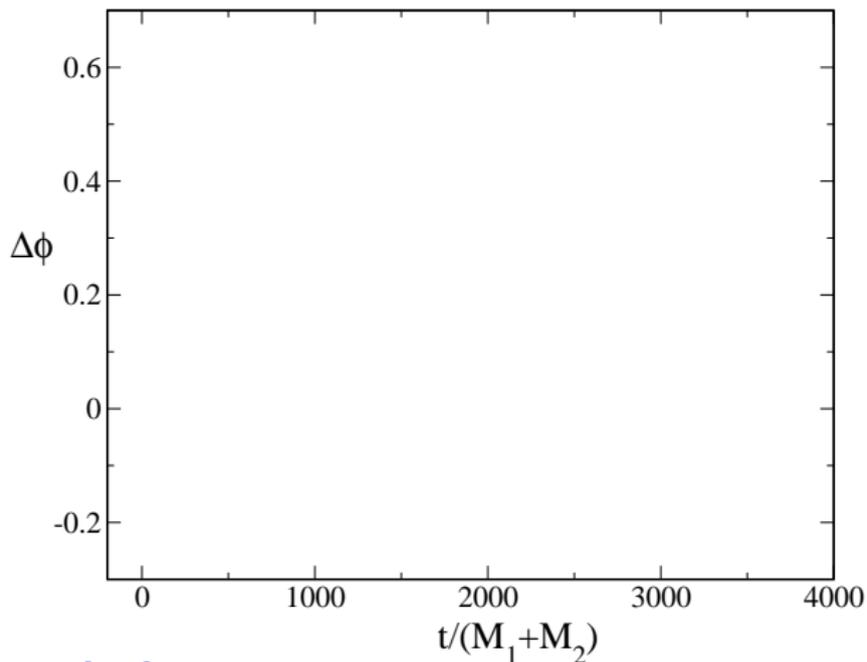
Waveforms



Gravitational wave phase

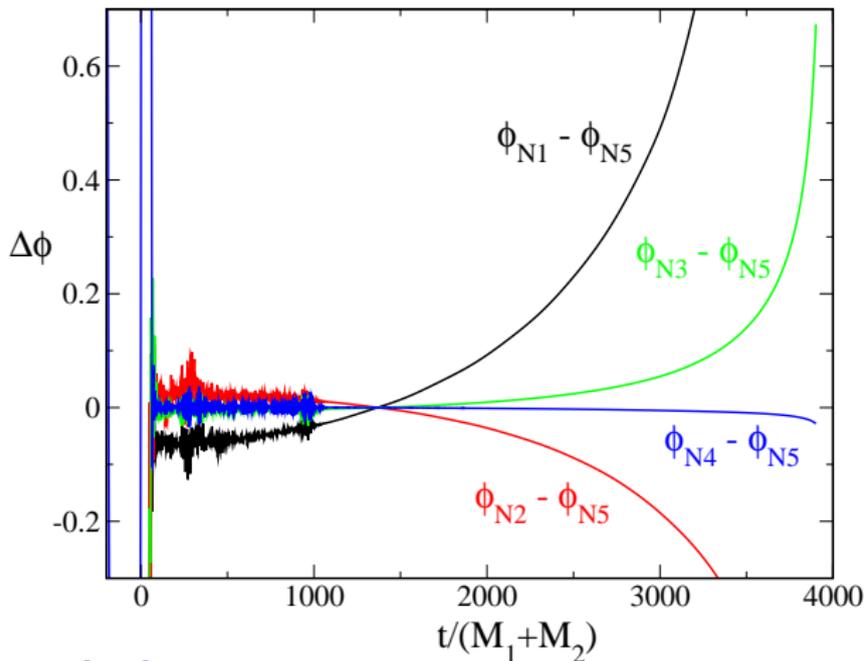


Comparison with post-Newtonian (GW-phase)



Numerical error budget

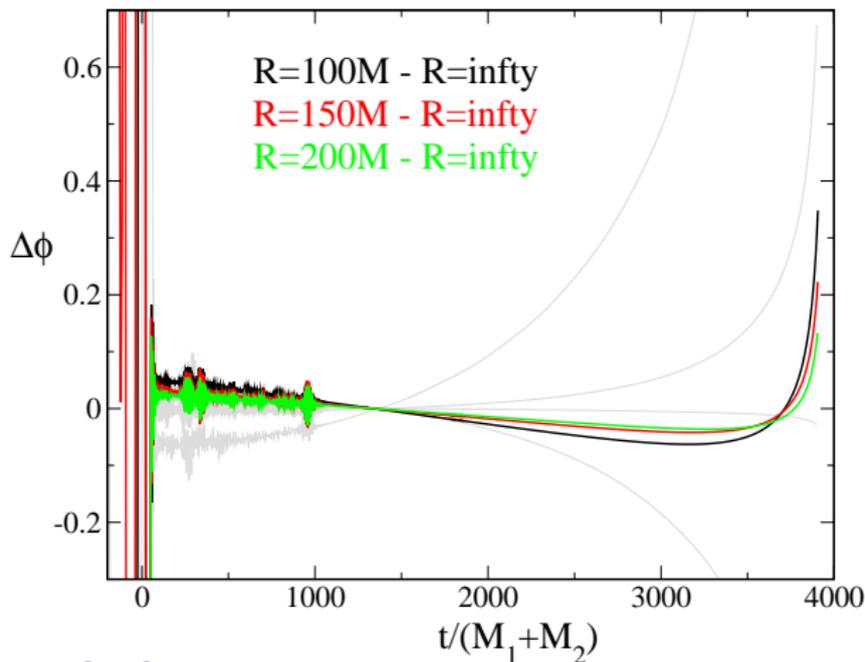
Comparison with post-Newtonian (GW-phase)



Numerical error budget

- Truncation error $\sim 0.03\text{rad}$

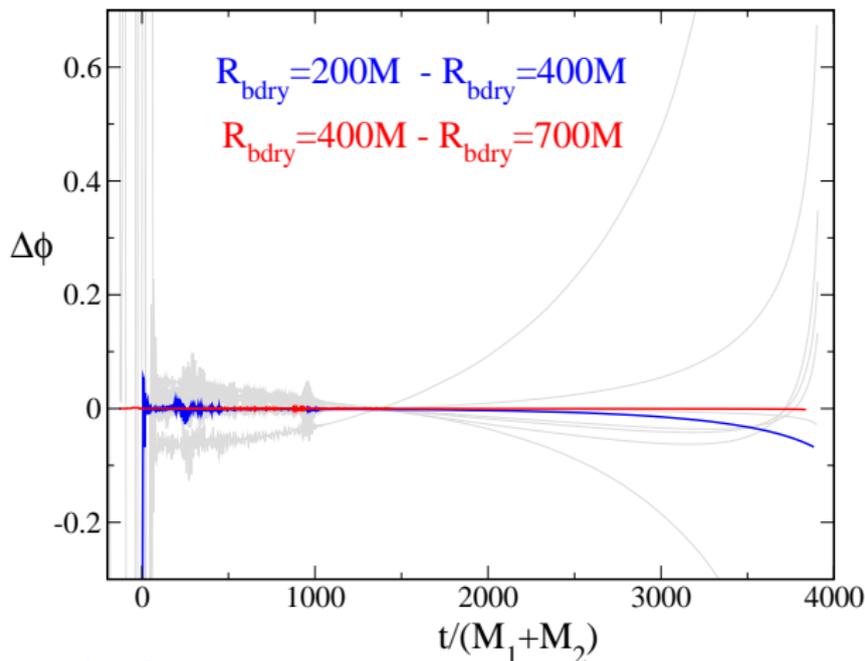
Comparison with post-Newtonian (GW-phase)



Numerical error budget

- Truncation error $\sim 0.03\text{rad}$
- Extraction radius $\sim 0.1\text{rad}$

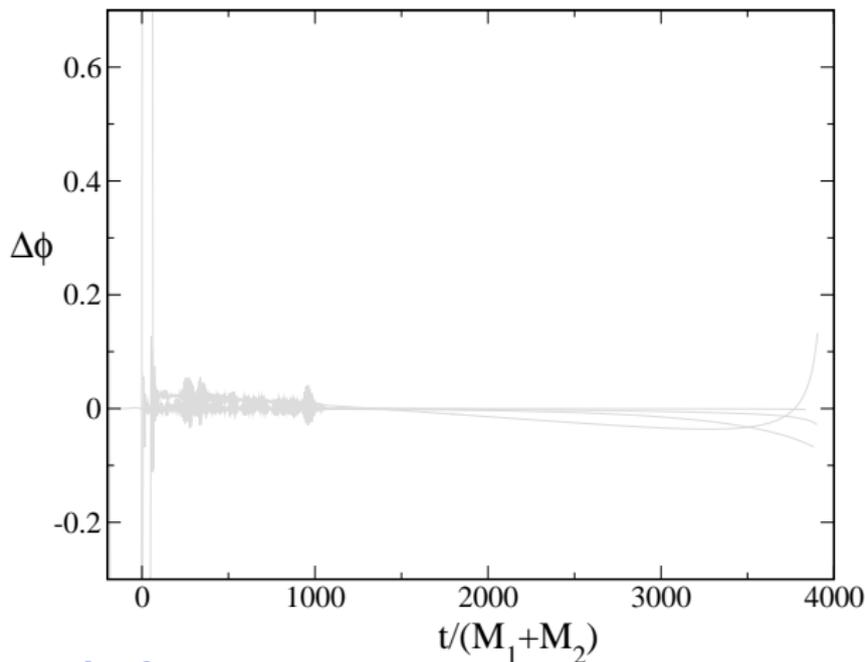
Comparison with post-Newtonian (GW-phase)



Numerical error budget

- Truncation error $\sim 0.03\text{rad}$
- Extraction radius $\sim 0.1\text{rad}$
- Outer boundary $\sim 0.07\text{rad}$

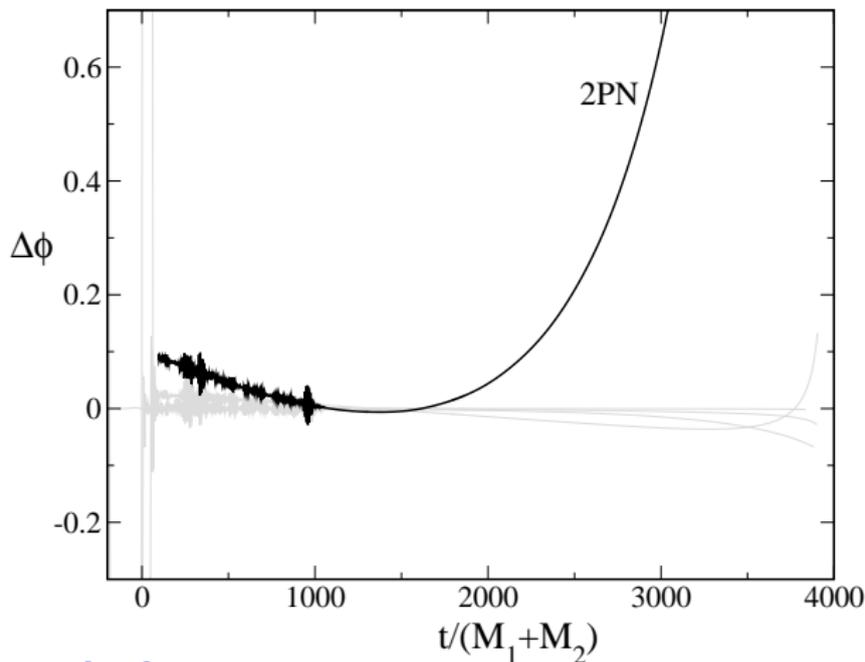
Comparison with post-Newtonian (GW-phase)



Numerical error budget

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- Extraction radius $\sim 0.1\text{rad}$
- Outer boundary $\sim 0.07\text{rad}$

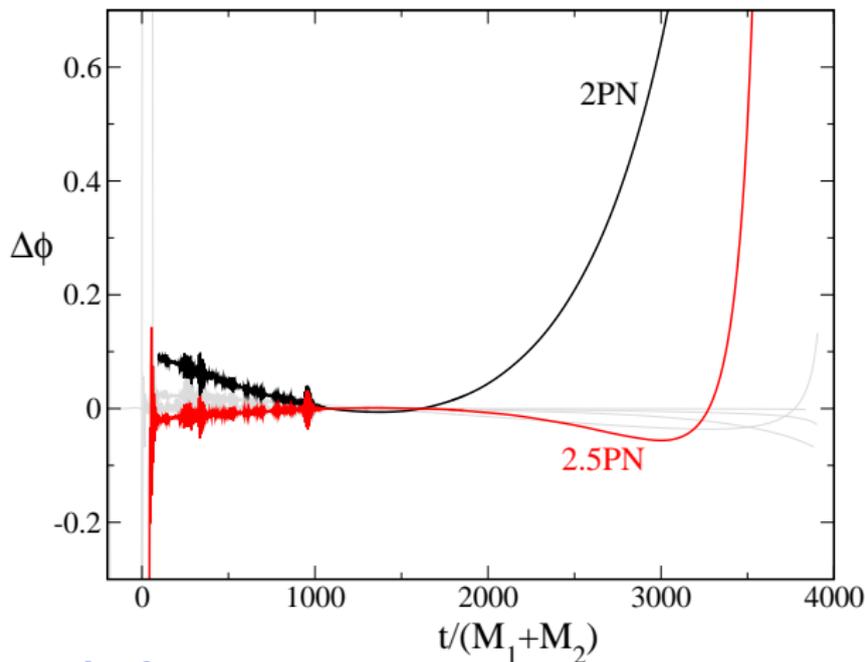
Comparison with post-Newtonian (GW-phase)



Numerical error budget

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- Extraction radius $\sim 0.1\text{rad}$
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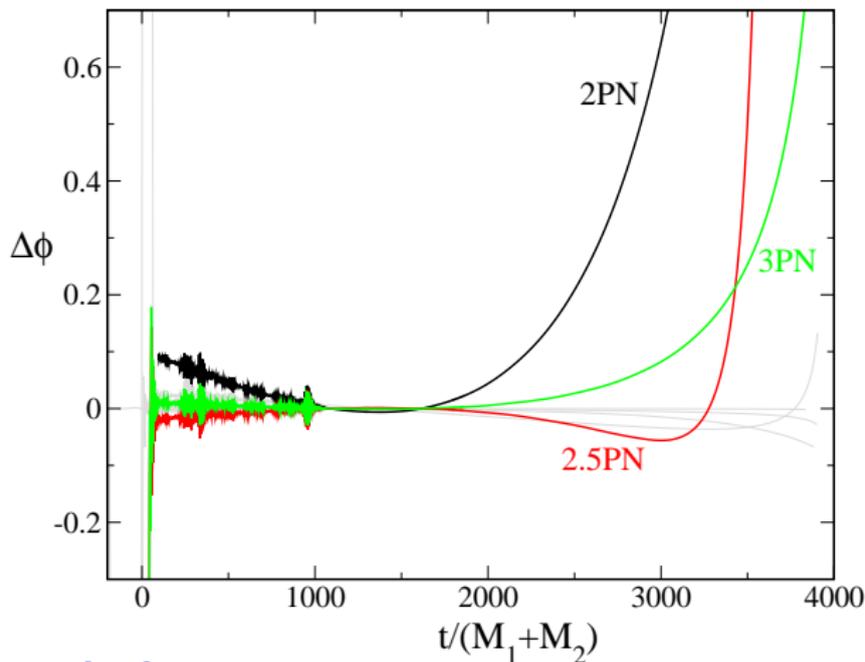
Comparison with post-Newtonian (GW-phase)



Numerical error budget

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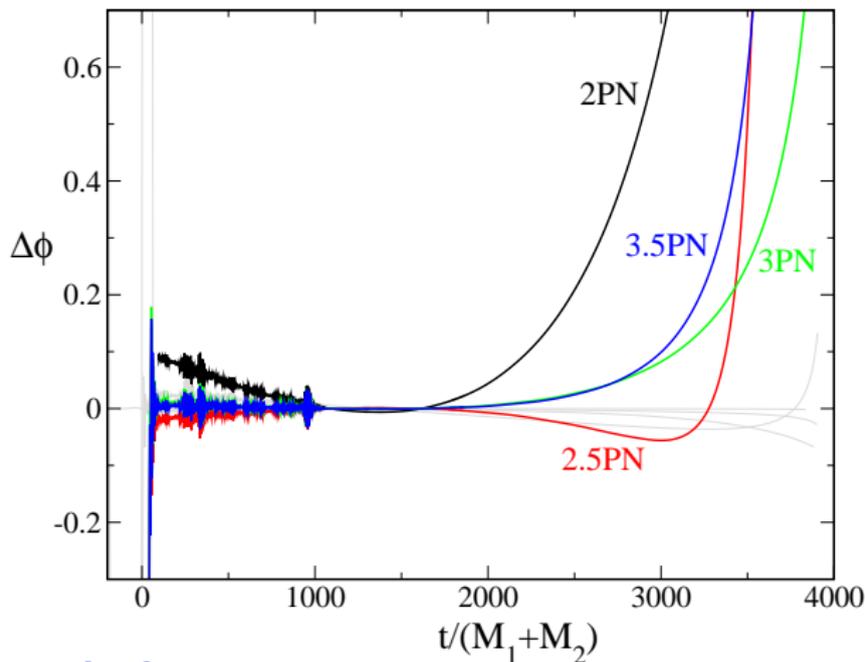
Comparison with post-Newtonian (GW-phase)



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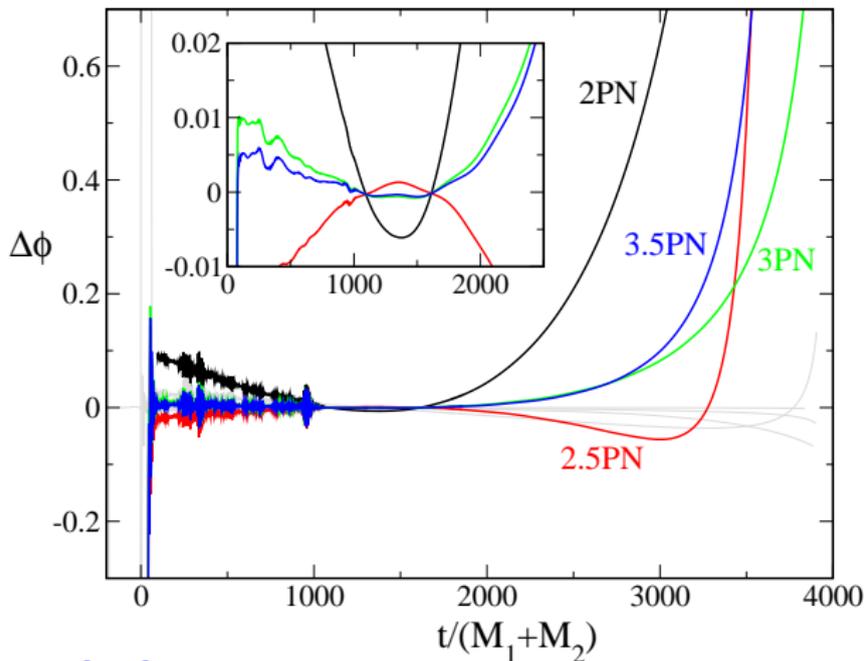
Comparison with post-Newtonian (GW-phase)



Numerical error budget

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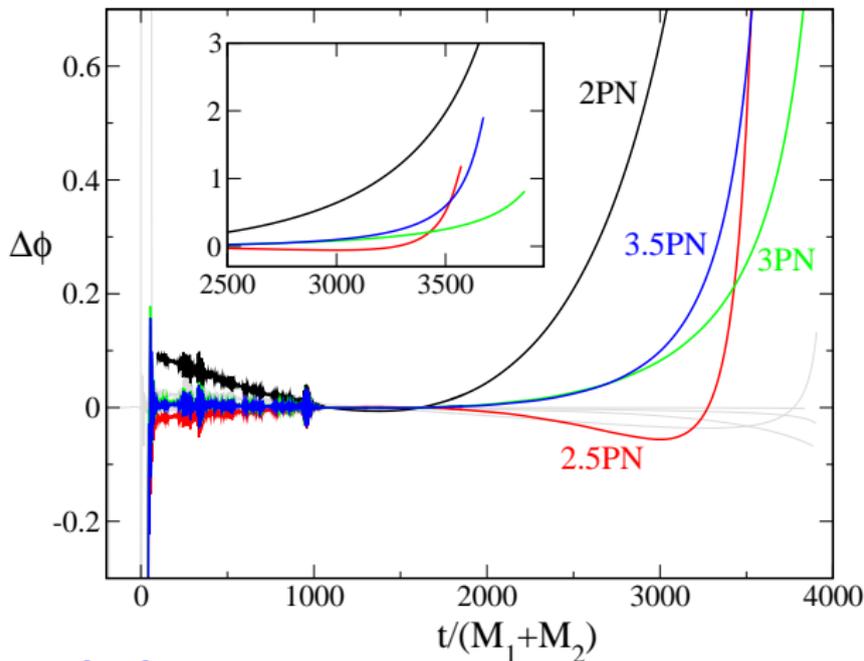
Comparison with post-Newtonian (GW-phase)



Numerical error budget

- Truncation error $\sim 0.03\text{rad}$
- Extraction radius $\sim 0.1\text{rad}$
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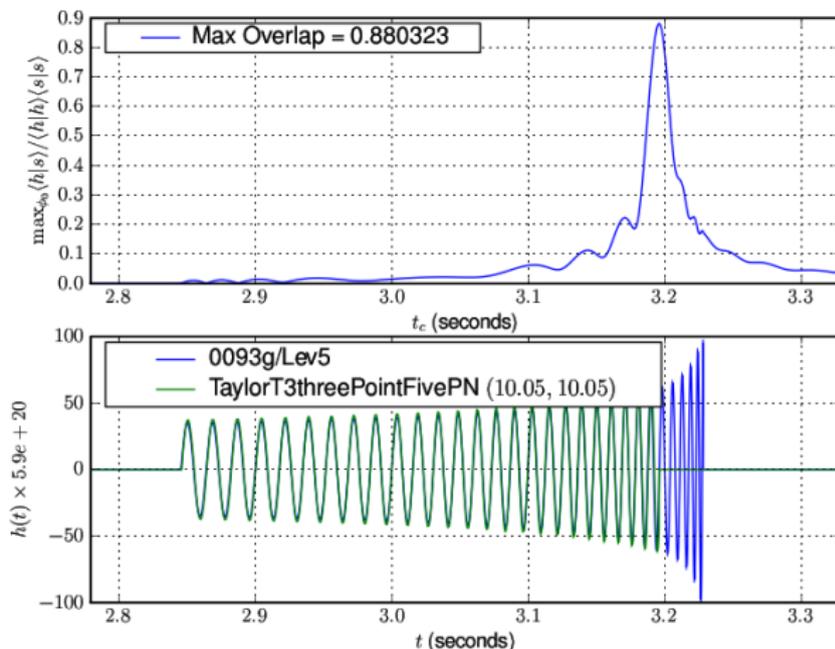
Comparison with post-Newtonian (GW-phase)



Numerical error budget

- Truncation error $\sim 0.03\text{rad}$
- Extraction radius $\sim 0.1\text{rad}$
- Outer boundary $\sim 0.07\text{rad}$

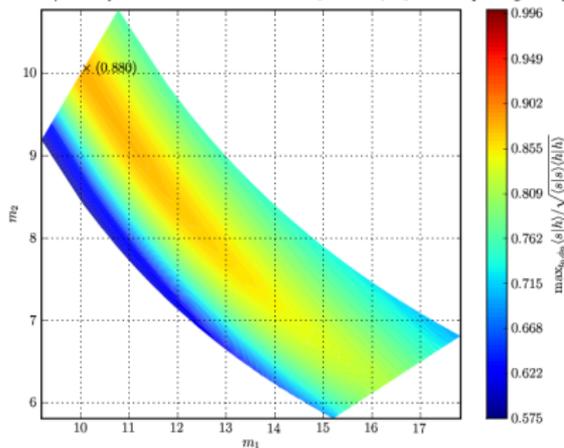
Testing LIGO detection templates



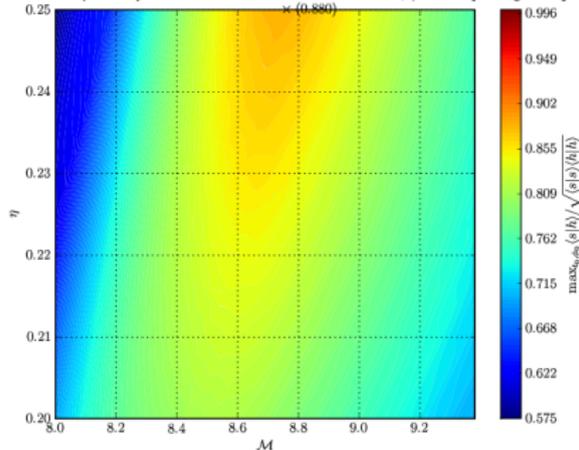
- In collaboration with D. Brown
- Good overlap!

Parameter estimation

Maximum overlap for TaylorT3threePointFivePN at $m_1 = 10.05$, $m_2 = 10.05$ [0093g/Lev5]



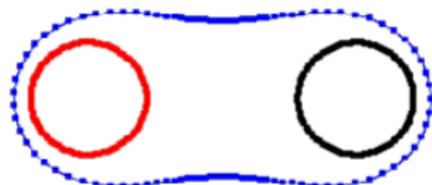
Maximum overlap for TaylorT3threePointFivePN at $\mathcal{M} = 8.75$, $\eta = 0.25$ [0093g/Lev5]



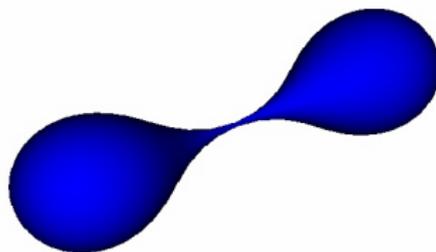
- Masses in simulation: 10.07Msun. Recovered 10.05Msun.

Head-on Merger

- Evolve to common horizon, regrid, continue.
- Apparent Horizons



Event horizon (M. Cohen, in prep.)



- Orbiting BHs have reached separation $2.2M$
(common horizon forms at $\sim 2M$.)

Summary

- Numerical relativity is in its golden age
- Spectral methods achieve stunning accuracy
- Comparison to PN in progress
 - ▶ Equal-mass non-spinning BHs
 - Full agreement NR - PN up to ~ 15 cycles before merger
 - Sufficient accuracy to identify deviations at all available PN-orders
- Future
 - ▶ Merger!
 - ▶ PN comparisons for spinning, non-equal mass binaries
 - Blanchet: PN converges exceptionally fast for q
 - For $q > 1$, $T_{\text{inspiral}} \propto q$