

# Spin down (?) of protostars through gravitational torques

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# Plan

- Problem
- Numerical modelling
- Simulation 1: hindered spin down
- Simulation 2: spin down
- Discussion and implications
- Problems

# Problem

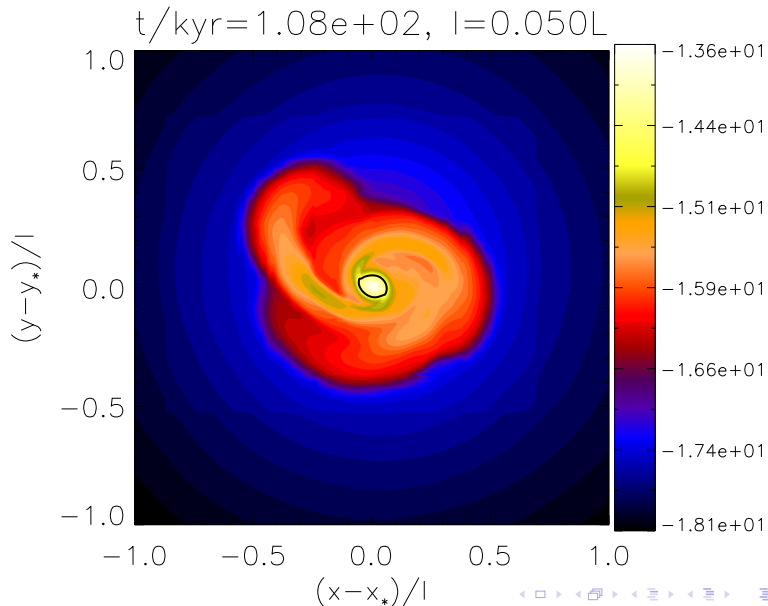
Angular momentum problem in star formation (Bodenheimer, 1995).

Specific angular momenta (cgs) of:

- Molecular cloud core,  $j \sim 10^{21}$
- T Tauri,  $j \sim 10^{17}$
- Sun,  $j \sim 10^{15}$

Need to transport angular momentum out of central region, otherwise star spins too fast. How?

# Problem



# Governing equations

Inviscid, non-magnetised and self-gravitating fluid with customised equation of state:

$$\begin{aligned}\frac{D\rho}{Dt} &= -\rho\nabla\cdot\mathbf{v} \\ \frac{D\mathbf{v}}{Dt} &= -\frac{1}{\rho}\nabla P - \nabla\Phi \\ \rho\frac{De}{Dt} &= -\rho\mathbf{v}\cdot\nabla\Phi - \nabla\cdot(P\mathbf{v}) \\ \nabla^2\Phi &= 4\pi G\rho, \\ P &= c^2\rho^{\gamma_1}\left[1 + \left(\frac{\rho}{\rho_*}\right)^{\gamma_2-\gamma_1}\right].\end{aligned}$$

$c \simeq 266\text{ms}^{-1}$  is isothermal sound speed of  $\mu = 2.33$  gas at 20K,  $\gamma_1 = 1$  and  $\gamma_2 = 5/3$ . Here,  $e$  is the sum of internal and kinetic energy densities.

# Angular momentum transport

Conservation equation

$$\frac{\partial}{\partial t} (\rho R v_\phi) + \nabla \cdot (\rho R v_\phi \mathbf{v}) = -\rho \frac{\partial \Phi}{\partial \phi} - \frac{\partial P}{\partial \phi},$$

integrate over volume  $V$ , use  $\rho = \nabla^2 \Phi / 4\pi G$  on RHS to get:

$$\frac{\partial J}{\partial t} + \oint \mathbf{F} \cdot d\mathbf{S} = 0$$

with

$$\begin{aligned} \mathbf{F} &= \mathbf{F}_A + \mathbf{F}_G \\ \mathbf{F}_G &= \frac{1}{4\pi G} \frac{\partial \Phi}{\partial \phi} \nabla \Phi. \end{aligned}$$

Also have the Reynolds stress,  $\rho R \delta v_\phi \delta \mathbf{v}$ .

## Collapse into star-disc system

Kratter et al. (2010) description of collapse of spherical, rotating cloud into a disc (mass  $M_d$ ) with central object (mass  $M_*$ ). Call

$$M_* + M_d = M_{\text{sys}}.$$

- Infall parameter

$$\xi = \frac{GM\dot{M}}{c^3}$$

- Rotation parameter

$$\Gamma = \frac{\dot{M}}{M_{\text{sys}}\Omega_k}$$

$\Omega_k$  is Keplerian frequency, due to  $M_{\text{sys}}$ , of material joining the system from the cloud, assumed to occur at cylindrical radius  $R_k$ . Can show disc aspect-ratio

$$h = \left(\frac{\Gamma}{\xi}\right)^{1/3},$$

and  $R_k = h^2 \xi ct$ .

# Initial conditions & numerical method

- Start with spherical cloud of radius  $r_c$  of density profile

$$\rho(r) = \frac{Ac^2}{4\pi Gr^2}$$

Shu (1977)  $\rightarrow$  self-similar collapse.

- Designate a central region  $r \leq r_* \equiv qr_c$  to be the 'star'. Set  $\rho_* = \rho(r_*)$ .
- Set  $v_r = v_\theta = 0$  and azimuthal velocity

$$v_\phi = 2Ach \times \begin{cases} R/r_* & R \leq r_* \\ 1 & R > r_* \end{cases}$$

- Need  $2h\sqrt{A} < 1$  for below break-up speed.

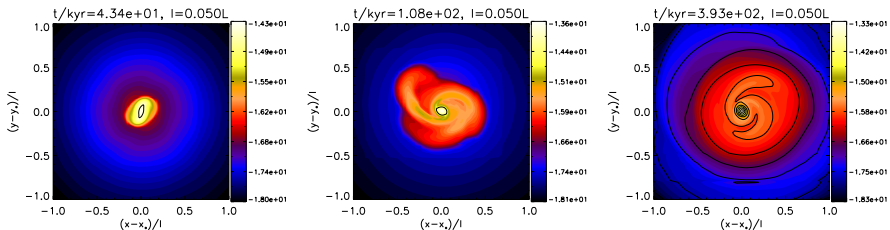


# Initial conditions & numerical method

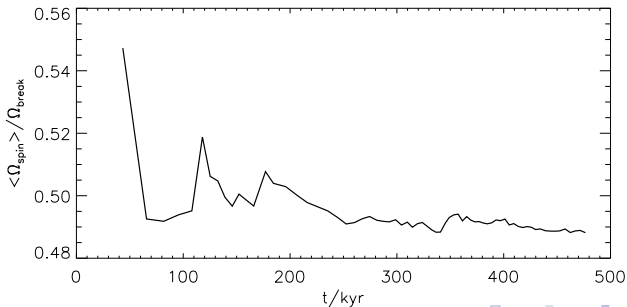
- Solve in Cartesian box of length  $L = 4r_c$ .
- ORION: Godunov-type code with adaptive mesh refinement.
- Base grid  $128^3$ , 6 refinement levels (effective highest resolution  $8192^3$ ).

# Case 1: $\xi = 5.58$ , $h = 0.1$ , $q = 0.005$

Density slices:

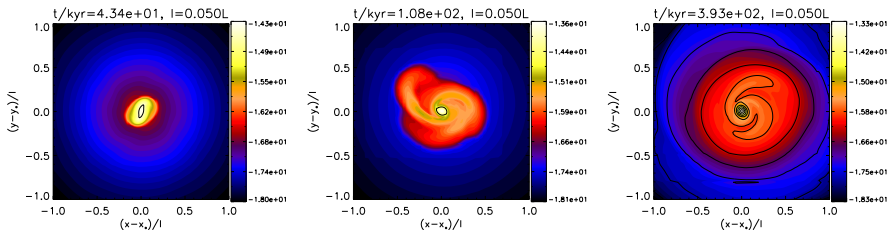


Star spin

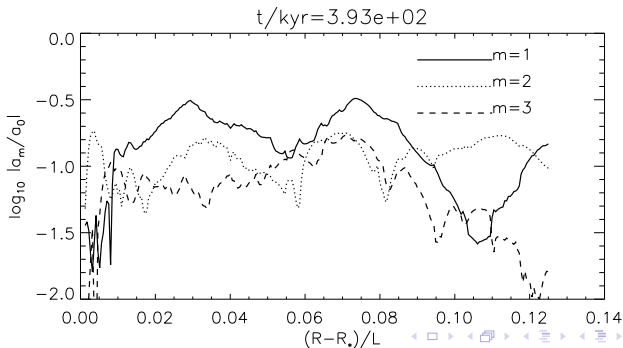


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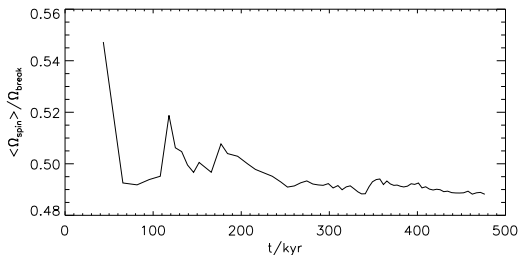


Mode amplitudes

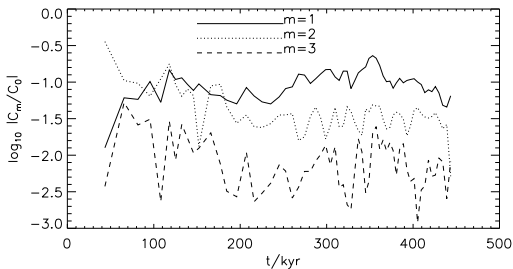


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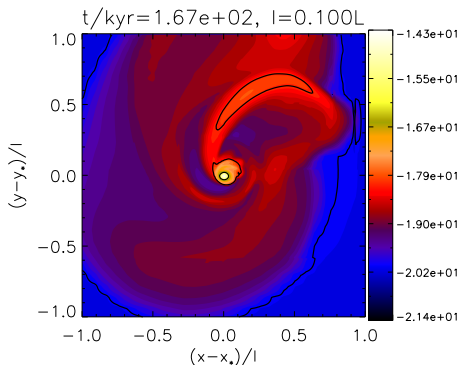
Star spin



Modes evolution

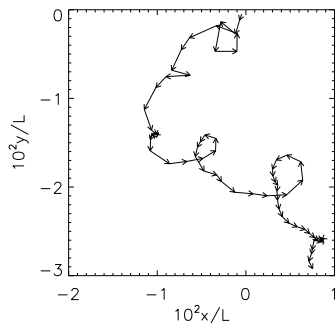


# Influence of $m = 1$ modes



- Theoretical studies: Adams et al. (1989); Heemskerk et al. (1992).
- **$m = 1$  displaces star from COM (of box)**
- Require sufficient disc mass.
- Exchange of *orbital* angular momentum.

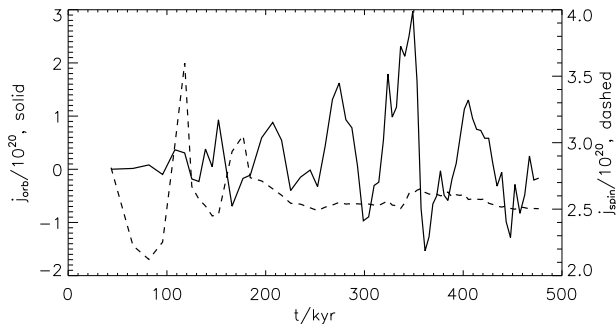
# Case 1: stellar motion & angular momenta



## Estimates:

- $\Omega_{\text{spin}} \sim 1.5 \times 10^{-10}$
- $\Omega_{\text{orb}} \sim 2 \times 10^{-13}$
- $\Omega_{\text{disc}}(R_k) \sim 6 \times 10^{-12}$
- $\Omega_{\text{patt}} \sim (3-5) \times 10^{-12}$

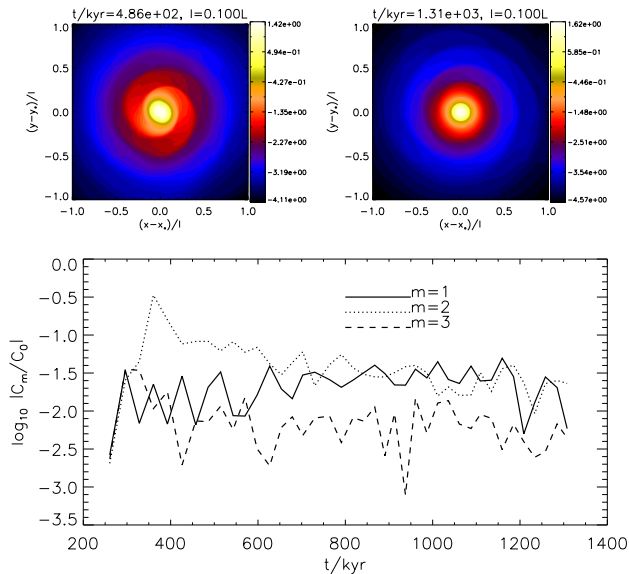
# Case 1: stellar motion & angular momenta



## Estimates:

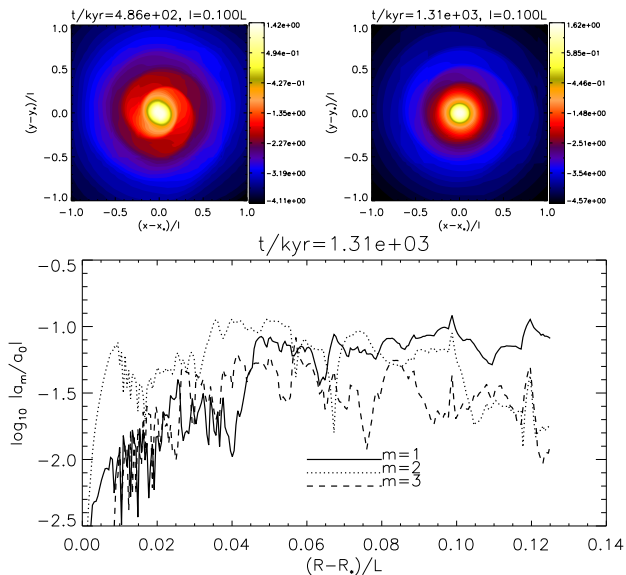
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## Case 2: $\xi = 2.74$ , $h = 0.05$ , $q = 0.01$

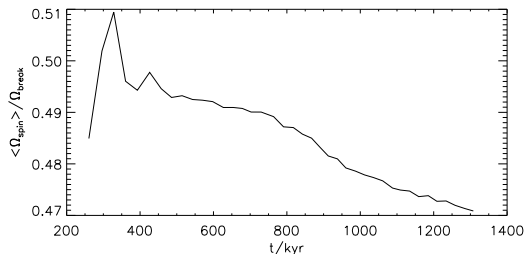




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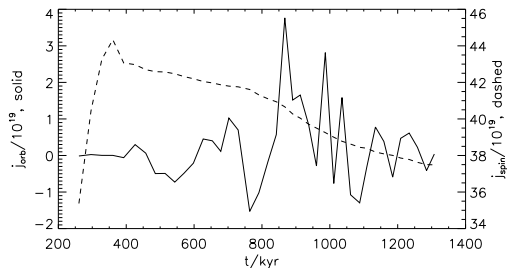


## Case 2: $\xi = 2.74$ , $h = 0.05$ , $q = 0.01$



- Spin down around  $t = 500\text{kyr}$  with visible  $m = 2$ .
- Spin down near the end but no visible  $m = 2$ , although FT  $\rightarrow m = 2$  still the main non-axisymmetry near star.
- FT  $\rightarrow m = 1$  becomes important in outer region but limited orbital motion.

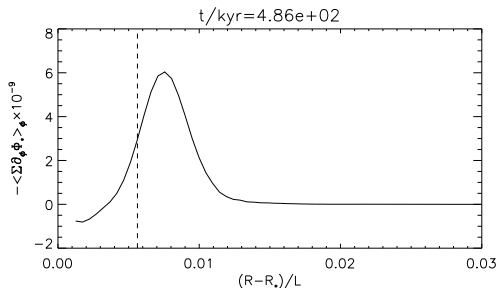
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# Star torque & gravity flux

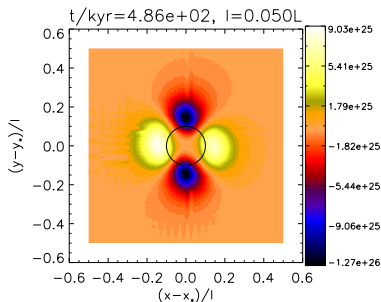
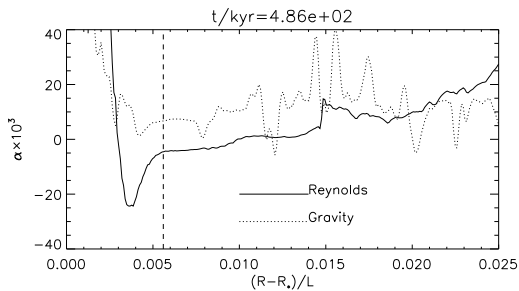
$\nabla^2 \phi_* = 4\pi G \rho_{\text{star}} \rightarrow$  get star torque per unit area:



consistent with disc-on-star torques, but...

# Star torque & gravity flux

$\partial_t J + \oint \mathbf{F} \cdot d\mathbf{S} = 0 \rightarrow$  look at radial gravity flux near star

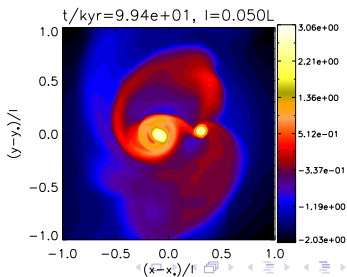
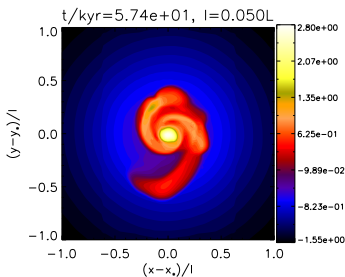
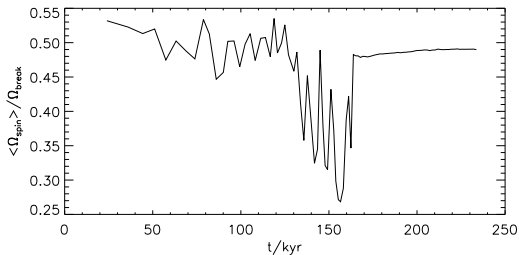


Solid and dashed line: characteristic star size.

- $\alpha < O(10^{-2})$  also reported in Kratter et al. (2010) but is SMALL compared to numerical  $\alpha$ !
- Numerical spin down? But why ineffective in Case 1?

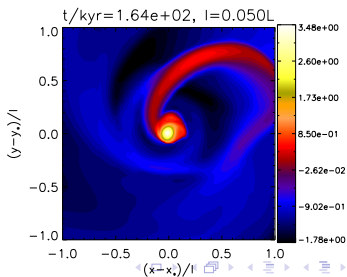
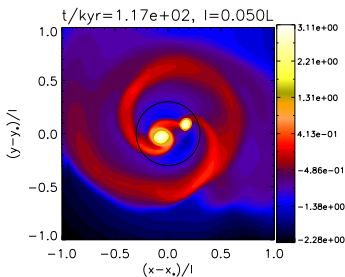
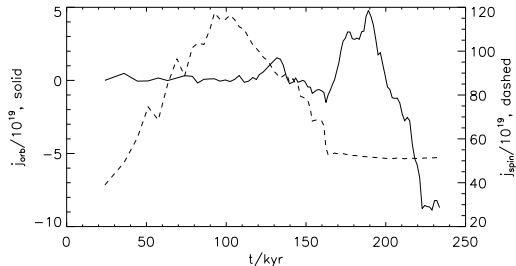
# Case 3: binary spin down

Theoretical work: Boss (1984).

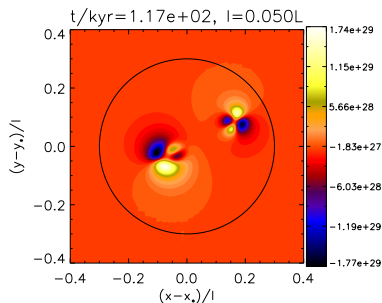
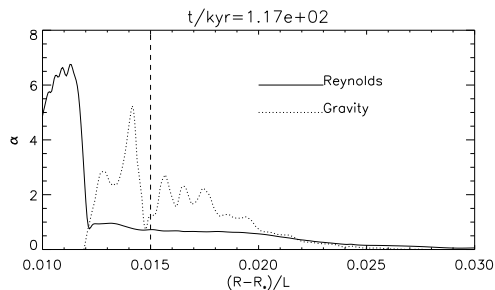


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# Gravity flux into binary

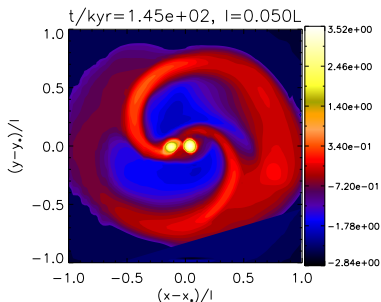


- Gravity  $\alpha \sim 2$  from binary  $\rightarrow$  torque down
- Numerical  $\alpha < 0.5$  in this region
- $\Omega_{\text{spin}} > \Omega_p$  on circle



# Problems

- Spin down using gravity flux? Maybe, but cannot have influences from  $m = 1$ .
- Need better experiment designs to overcome numerical spin down.
- Binary results are less unconvincing.



Thank you

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# References

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