

CBI pol to Apr'05

Acbar to Jan'06

Bicep  
QUaD

CBI2 to Apr'07

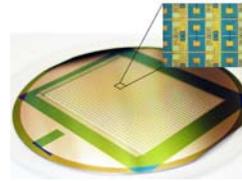
Quiet2

(1000 HEMTs)

Quiet1 Chile

SZA  
(Interferometer)  
California

APEX  
(~400 bolometers)  
Chile



SCUBA2  
(12000 bolometers)

JCMT, Hawaii

ACT

(3000 bolometers)  
Chile

Spider

(2312 bolometer LDB)

Clover

2017  
CMBpol

Boom03

2003

2005

2007

2004

2006

2008

WMAP ongoing to 2009

SPT

(1000 bolometers)  
South Pole

ALMA  
(Interferometer)  
Chile

DASI

Polarbear

(300 bolometers)  
California

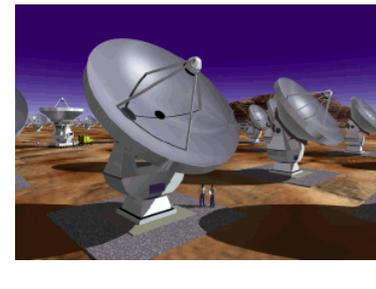
CAPMAP

AMI

(84 bolometers)  
HEMTs L2

GBT

Planck



# Standard Parameters of Cosmic Structure Formation

$$\theta \sim \ell_s^{-1}, \text{ cf. } \Omega_\Lambda$$

Period of inflationary expansion,  
quantum noise  $\rightarrow$  metric perturbations

$r < 0.6$  or  $< 0.25$  95% CL

$$\Omega_k$$

$$\Omega_b h^2$$

$$\Omega_{dm} h^2$$

$$\Omega_\Lambda$$

$$\tau_c$$

$$n_s$$

$$n_t$$

$$\ln A_s \sim \ln \sigma_8$$

$$r = A_t / A_s$$

- Inflation predicts nearly scale invariant and background of gravitational waves
  - Passive/adiabatic/coherent/gaussian
  - Nice linear regime
  - Boltzman equation + Einstein equations
- What predicts nearly scale invariant and background of gravitational waves?
- Density interactions
- Spacetime curvature
- Optical Depth to Last Scattering Surface
- When did stars reionize the universe?
- Amplitude
- Amplitude

$$\Omega_k > 0$$

$$\Omega_k = 0$$

$$\Omega_k < 0$$

flat

open

# New Parameters of Cosmic Structure Formation

$\Omega_k$

$\Omega_b h^2$

$\Omega_{dm} h^2$

$\tau_c$

$\theta \sim \ell_s^{-1}, \text{ cf. } \Omega_\Lambda$

$\ln \mathcal{P}_s(k)$

**scalar spectrum**  
use order N Chebyshev  
expansion in  $\ln k$ ,  
N-1 parameters  
amplitude(1), tilt(2),  
running(3), ...  
(or N-1 nodal point k-  
localized values)

$\ln \mathcal{P}_t(k)$

**tensor (GW) spectrum**  
use order M Chebyshev  
expansion in  $\ln k$ ,  
M-1 parameters  
amplitude(1), tilt(2), running(3),...

Dual Chebyshev expansion in  $\ln k$ :

Standard 6 is Cheb=2

Standard 7 is Cheb=2, **Cheb=1**

Run is Cheb=3

Run & tensor is Cheb=3, **Cheb=1**

Low order N,M power law but high  
order Chebyshev is Fourier-like

# New Parameters of Cosmic Structure Formation

$$\Omega_k$$

$$\Omega_b h^2$$

$$\Omega_{dm} h^2$$

$$\tau_c$$

$$\theta \sim \ell_s^{-1}, \text{ cf. } \Omega_\Lambda$$

$$\epsilon(k), \quad k \approx Ha$$

$$\ln H(k_p)$$

$$H(k_p)$$

**=1+q, the deceleration  
parameter history**

$$\mathcal{P}_s(\mathbf{k}) \propto \mathbf{H}^2 / \epsilon, \quad \mathcal{P}_t(\mathbf{k}) \propto \mathbf{H}^2$$

**order N Chebyshev  
expansion, N-1 parameters  
(e.g. nodal point values)**

**Hubble parameter at  
inflation at a pivot pt**

$$-\epsilon = d \ln \mathbf{H} / d \ln a$$

$$\frac{-\epsilon}{1-\epsilon} = \frac{d \ln \mathbf{H}}{d \ln k}$$

**Fluctuations are from stochastic kicks  $\sim H/2\pi$   
superposed on the downward drift at  $\Delta \ln k = 1$ .**

**Potential trajectory from HJ (SB 90,91):**

$$\mathbf{V} \propto \mathbf{H}^2 \left(1 - \frac{\epsilon}{3}\right); \quad \frac{d\psi_{\text{inf}}}{d \ln k} = \frac{\pm \sqrt{\epsilon}}{1-\epsilon}$$

$$\epsilon = (d \ln \mathbf{H} / d\psi_{\text{inf}})^2$$

tensor (gravity wave) power to curvature power,  $r$ , a direct measure of  $e = (q+1)$ ,  $q$ =deceleration parameter during inflation  $r \sim 16 e$

$q$  (ln Ha) may be highly complex (scanning inflation trajectories)

many inflaton potentials give the same curvature power spectrum, but the degeneracy is broken if gravity waves are measured

$(q+1) = \sim 0$  is possible - low energy scale inflation – upper limit only

Very very difficult to get at this with direct gravity wave detectors – even in our dreams

**Response of the CMB photons to the gravitational wave background leads to a unique signature within the CMB at large angular scales of these GW and at a detectable level. Detecting these B-modes is the new “holy grail” of CMB science.**

**Inflation prior: on  $e$  only 0 to 1 restriction,  $< 0$  supercritical possible**

**GW/scalar curvature:** current from CMB+LSS:  $r < 0.6$  or  $< 0.25$  (**.28**) 95%; good shot at **0.02** 95% CL with **BB polarization** (+- **.02 PL2.5+Spider**), **.01 target** BUT **foregrounds/systematics??** But **r-spectrum. But low energy inflation**

forecast  
Planck2.5

100&143

Spider10d

95&150

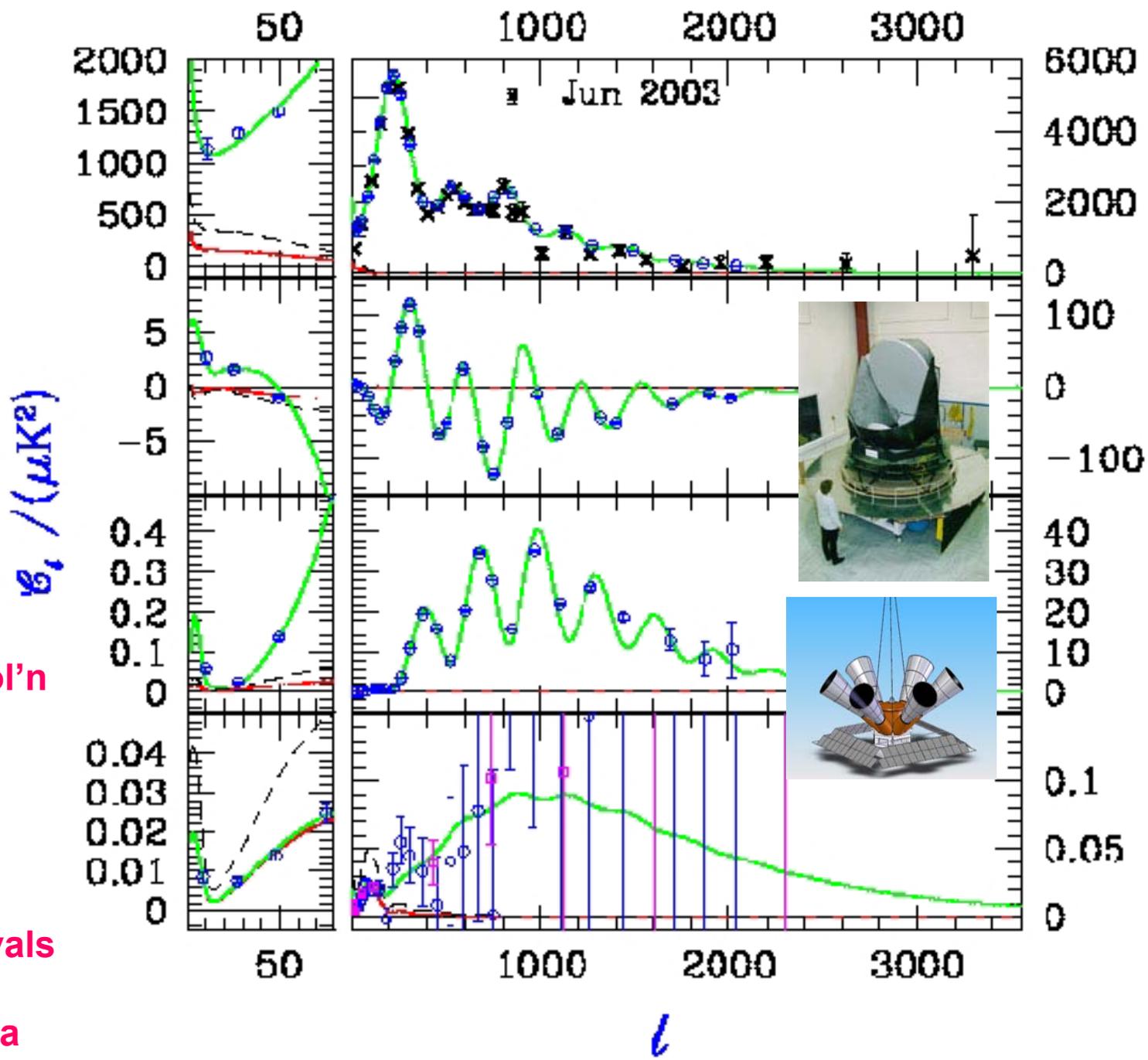
Synchrotron pol'n

< .004 ??

Dust pol'n

< 0.1 ??

Template removals  
from multi-  
frequency data



forecast

Planck2.5

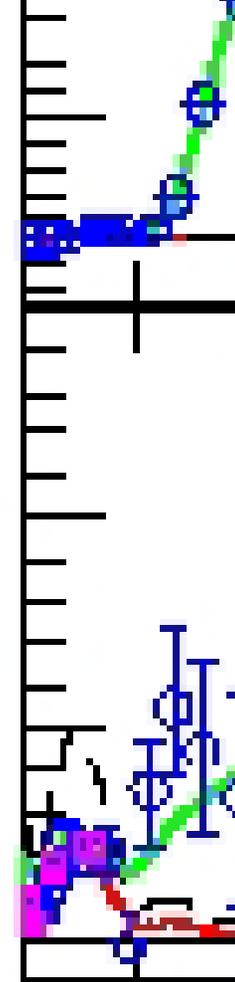
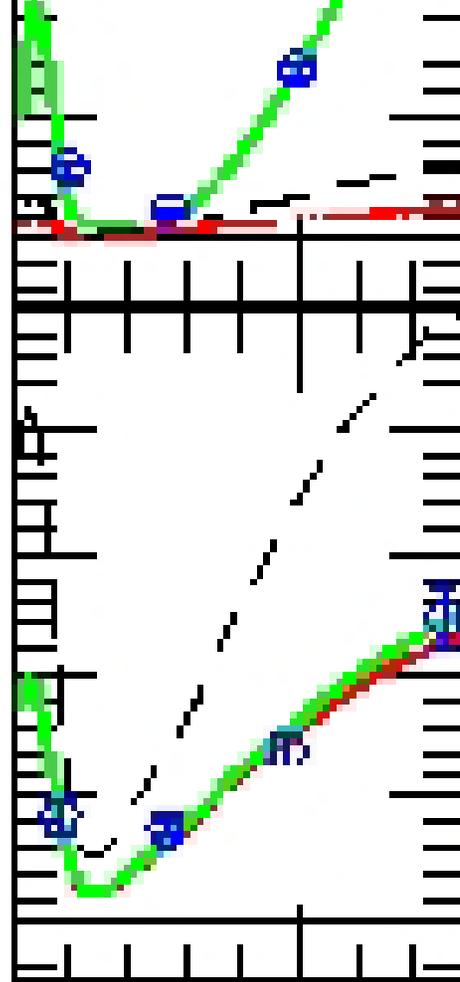
100&143

Spider10d

95&150

0.2  
0.1  
0

0.04  
0.03  
0.02  
0.01  
0



50

**GW/scalar curvature:** current from CMB+LSS:  $r < 0.6$  or  $< 0.25$  95% CL;  
 good shot at **0.02** 95% CL with **BB polarization** (+- .02 PL2.5+Spider Target .01)

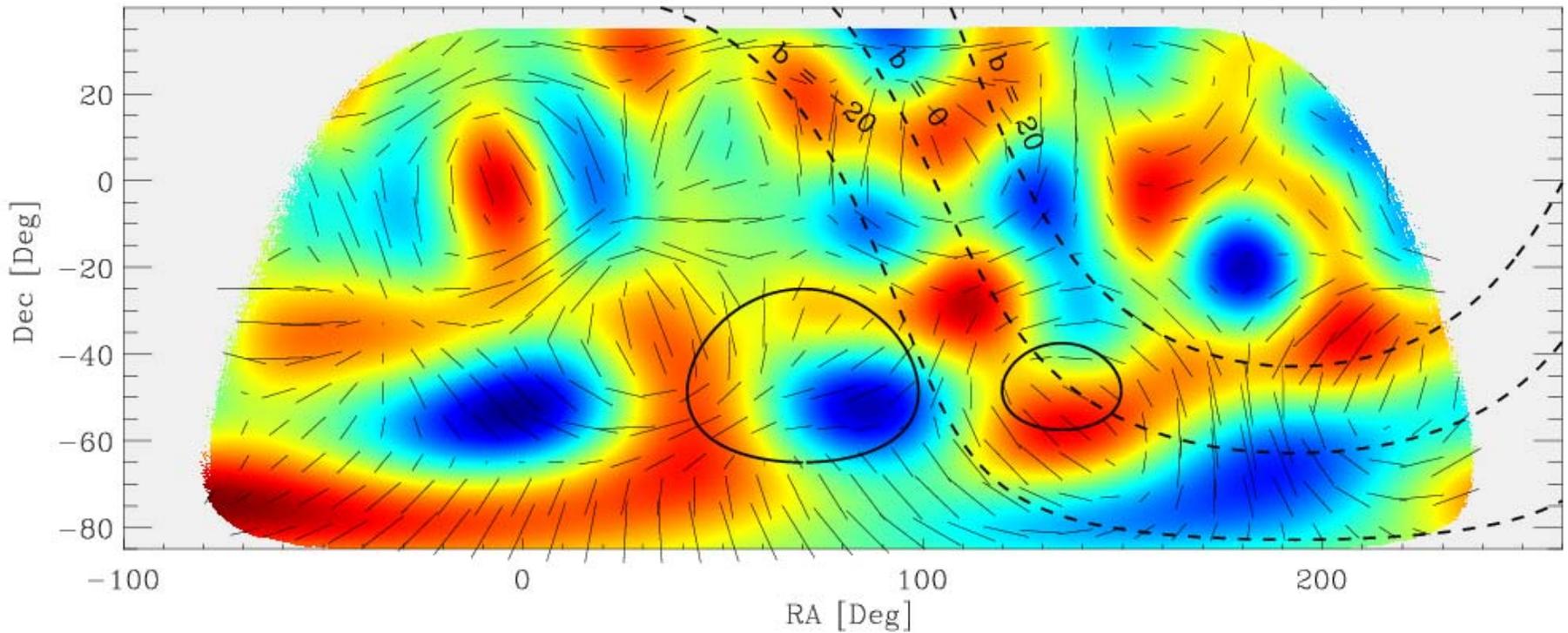
**BUT Galactic foregrounds & systematics??**

# SPIDER Tensor Signal

- Simulation of large scale polarization signal

$$\frac{A_T}{A_S} = 0.1$$

NonTensor



[http://www.astro.caltech.edu/~lgg/spider\\_front.htm](http://www.astro.caltech.edu/~lgg/spider_front.htm)

# String Theory Landscape & Inflation++ Phenomenology for CMB+LSS



14 std  
inflation  
parameters  
+ many many  
more e.g.  
“blind”  
search for  
patterns in  
the  
primordial  
power  
spectrum

any  
acceleration  
trajectory will  
do??

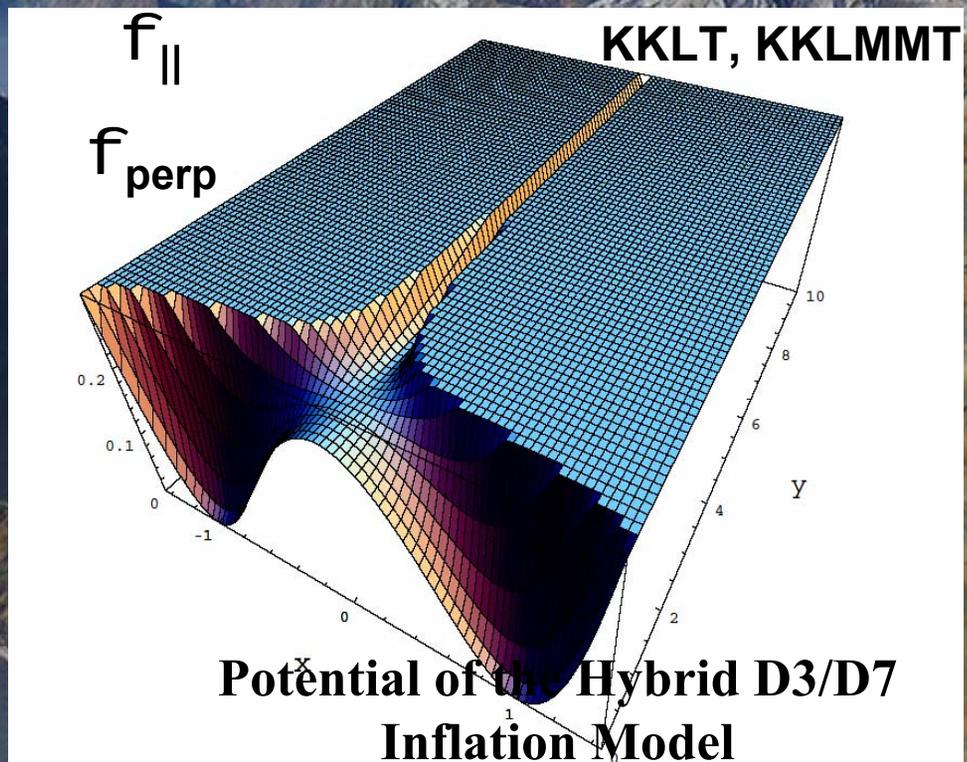
$q$  (ln  $H a$ )

$H$ (ln  $a, \dots$ )

$V(\phi, \dots)$

Measure??

anti-baroque  
prior



# Constraining Inflation Acceleration Trajectories

Bond, Contaldi, Kofman & Vaudrevange 06

“path integral” over probability landscape of theory and data, with mode-function expansions of the paths truncated by an imposed smoothness (Chebyshev-filter) criterion **[data cannot constrain high  $\ln k$  frequencies]**

$$P(\text{trajectory}|\text{data}, \text{th}) \sim P(\ln H_{p, \epsilon_k} | \text{data}, \text{th})$$

$$\sim P(\text{data} | \ln H_{p, \epsilon_k}) P(\ln H_{p, \epsilon_k} | \text{th}) \quad / P(\text{data} | \text{th})$$

Likelihood

theory prior

/ evidence

Data:

CMBall

(WMAP3, B03, CBI, ACBAR,

DASI, VSA, MAXIMA)

+

LSS (2dF, SDSS,  $\sigma_8$ [lens])

Theory prior

uniform in  $\ln H_{p, \epsilon_k}$

(equal a-prior probability hypothesis)

Nodal points cf. Chebyshev coefficients  
(linear combinations)

monotonic in  $\epsilon_k$

The theory prior matters alot

We have tried many theory priors

**Old view: Theory prior = delta function of THE correct one and only theory**

**New view: Theory prior = probability distribution on an energy landscape whose features are at best only glimpsed, huge number of potential minima, inflation the late stage flow in the low energy structure toward these minima. Critical role of collective geometrical coordinates (moduli fields) and of brane and antibrane “moduli” (D3,D7).**

**Ensemble of Kahler Moduli/Axion Inflations**

**Bond, Kofman, Prokushkin & Vaudrevange 06**

**A Theory prior in a class of inflation theories that seem to work**

**Low energy landscape dominated by the last few (complex) moduli fields  $T_1 T_2 T_3 \dots U_1 U_2 U_3 \dots$  associated with the settling down of the compactification of extra dims (complex) Kahler modulus associated with a 4-cycle volume in 6 dimensional Calabi Yau compactifications in Type IIB string theory. Real & imaginary parts are both important.**

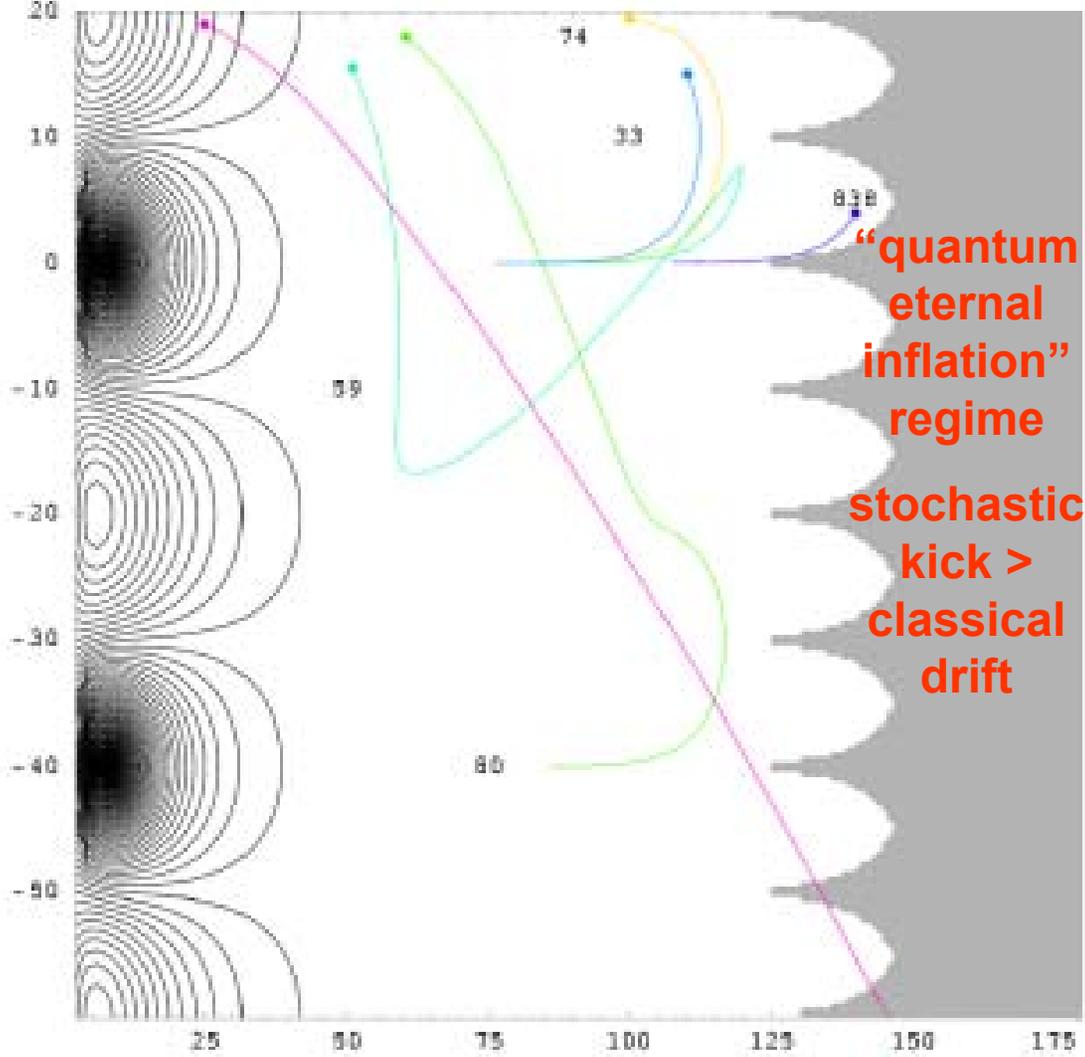
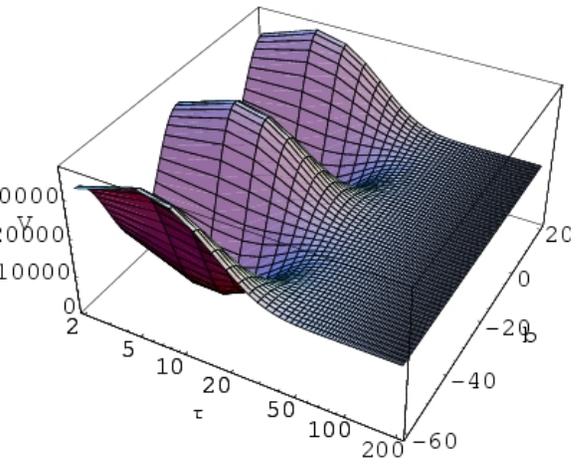
**Builds on the influential KKLT, KKLMNT moduli-stabilization ideas for stringy inflation and the Conlon and Quevada focus on 4-cycles. As motivated and protected as any inflation model. Inflation: there are so many possibilities: Theory prior ~ probability of trajectories given potential parameters of the collective coordinates X probability of the potential parameters X probability of initial collective field conditions**

**Sample trajectories  
in a Kahler  
modulus potential**

$\tau_2$  vs  $\theta_2$

$T_2 = \tau_2 + i\theta_2$

**Fixed  $\tau_1, \theta_1$**



**“quantum  
eternal  
inflation”  
regime**

**stochastic  
kick >  
classical  
drift**

**Sample Kahler modulus potential**

$$V(\tau, \theta) = \frac{8(a_2 A_2)^2 \sqrt{\tau} e^{-2a_2 \tau}}{3\alpha \lambda_2 V_s} + \frac{4W_0 a_2 A_2 \tau e^{-a_2 \tau} \cos(a_2 \theta)}{V_s^2} + \frac{3W_0^2 \xi}{4V_s^3} + V_{\text{uplift}}$$

# Beyond $P(k)$ : Inflationary trajectories

dynamical trajectory

$$u(\mathcal{I}) = \sum_{\beta} \phi_{\mathcal{I}\beta} q_{\beta} + r(\mathcal{I})$$

The mode amplitudes  $q_{\beta}$  are generalized bandpowers and the mode functions  $\phi_{\mathcal{I}\beta}$  are generalized splines or

$\beta$  as pairs  $(XP)$

- Economic way to scan the space of observables
- Increasing the order of Chebyshev expansion  
→ opening up the space of observables
- Huge degeneracy of  $V(\phi)$  without data for tensor modes

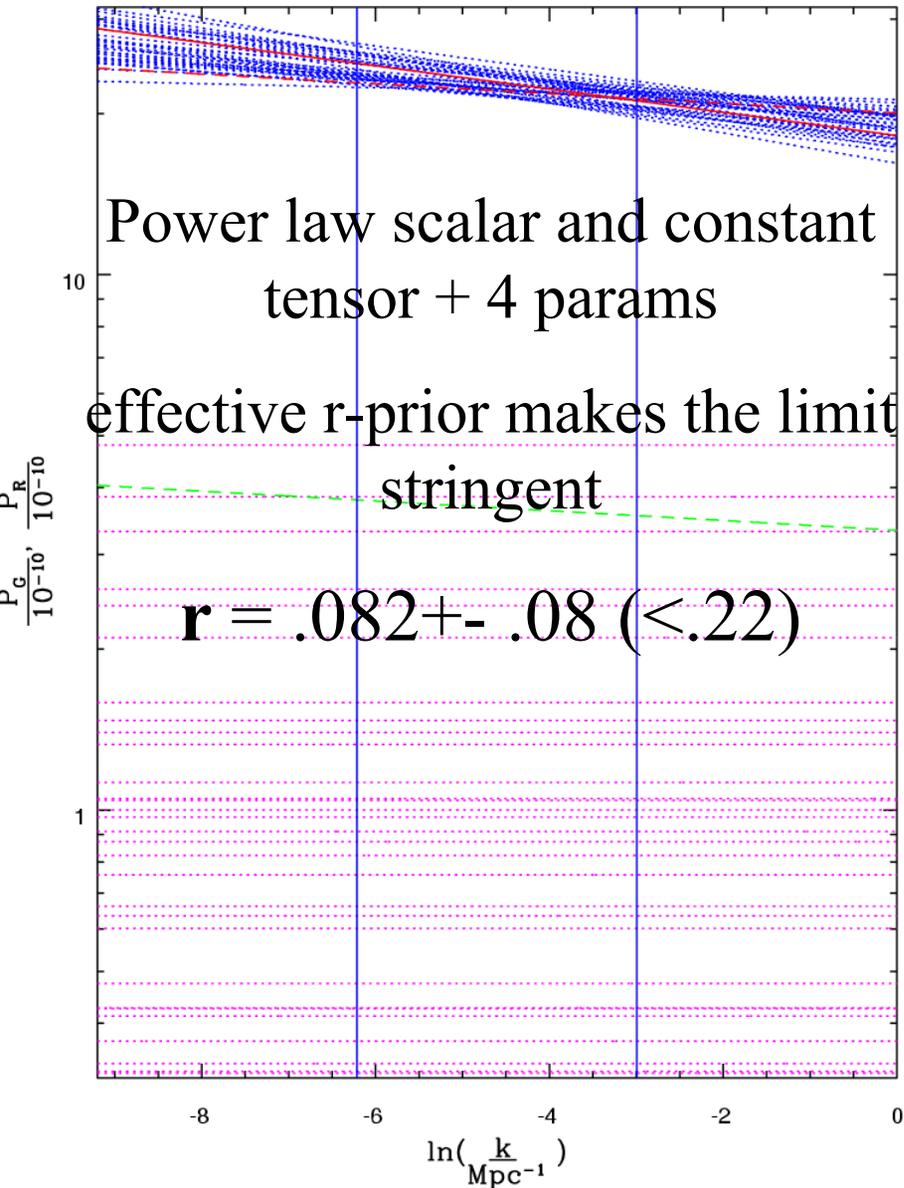
$H(\ln Ha)$

HJ + expand about uniform acceleration,  $1+q$ ,  $V$  and power spectra are derived

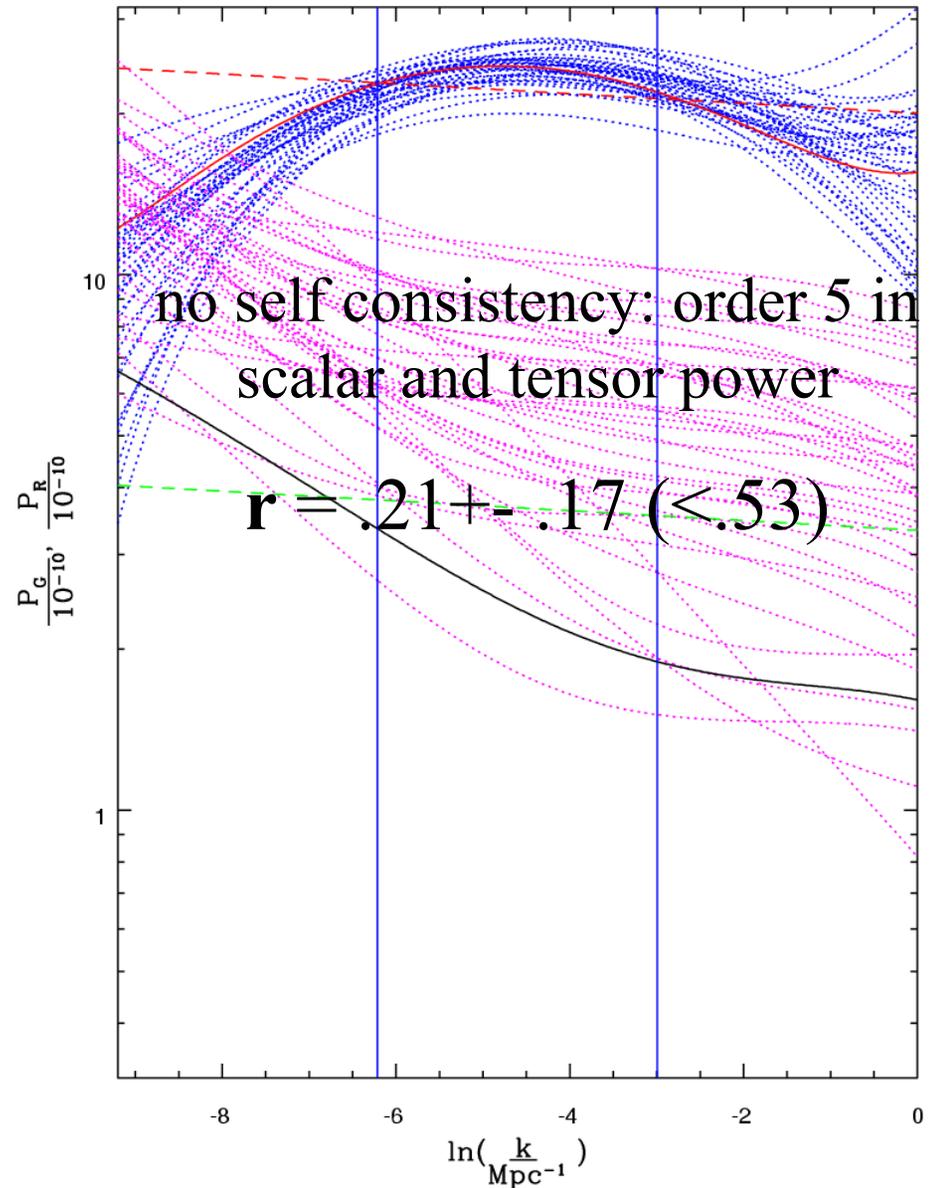
$$u_1 = \mathcal{P}_s / \mathcal{P}_s^{(s)} \quad u_2 = \mathcal{P}_t / \mathcal{P}_t^{(s)} \quad \ln k$$

**$\ln P_s P_t$  (nodal 2 and 1) + 4 params of  $P_s P_t$  (nodal 5 and 5) + 4 params**  
reconstructed from CMB+LSS data using Chebyshev nodal point expansion & MCMC

lnPR2\_1\_all\_paramsb.powerspectrum.likestats

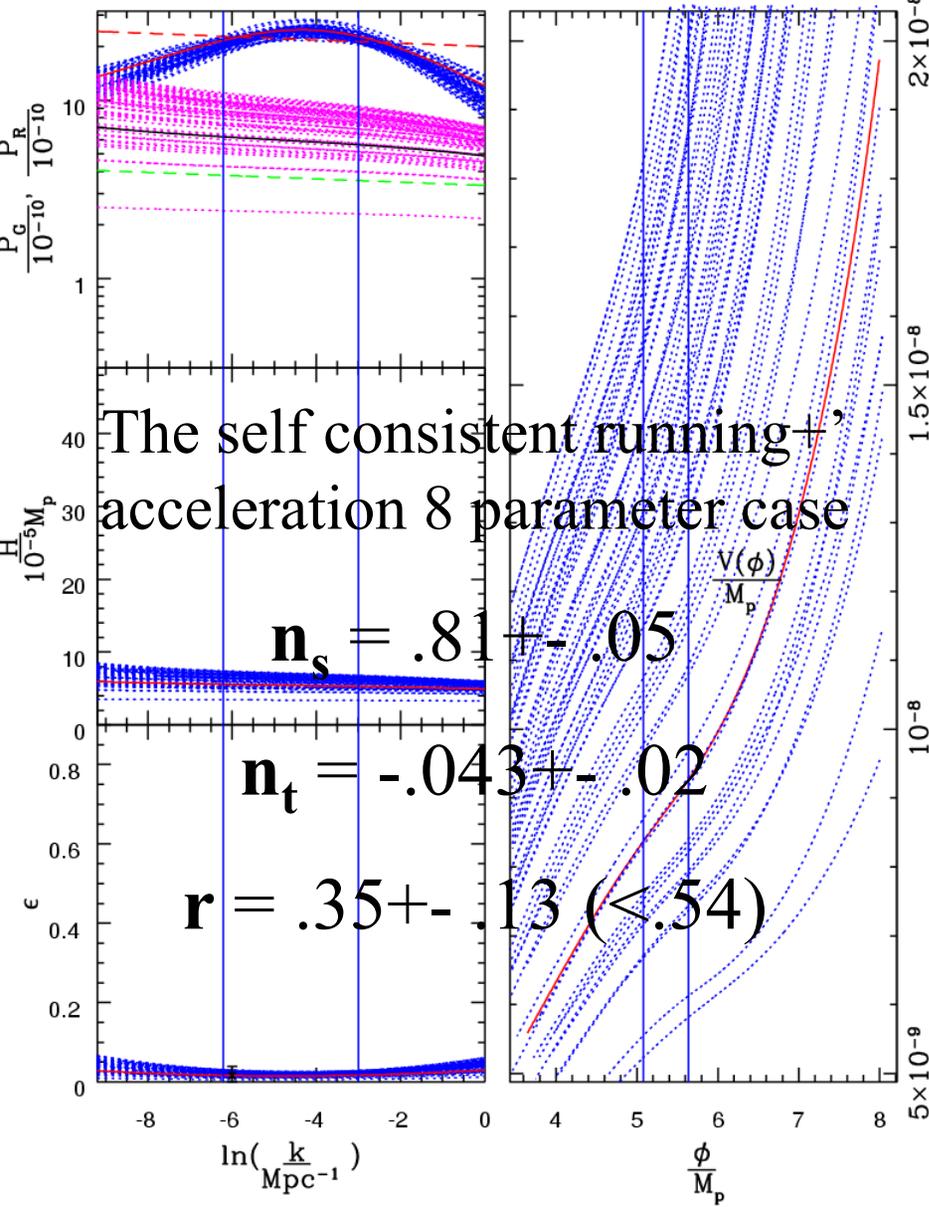


PR\_nodal5\_5\_all\_params\_cont.powerspectrum.likestats

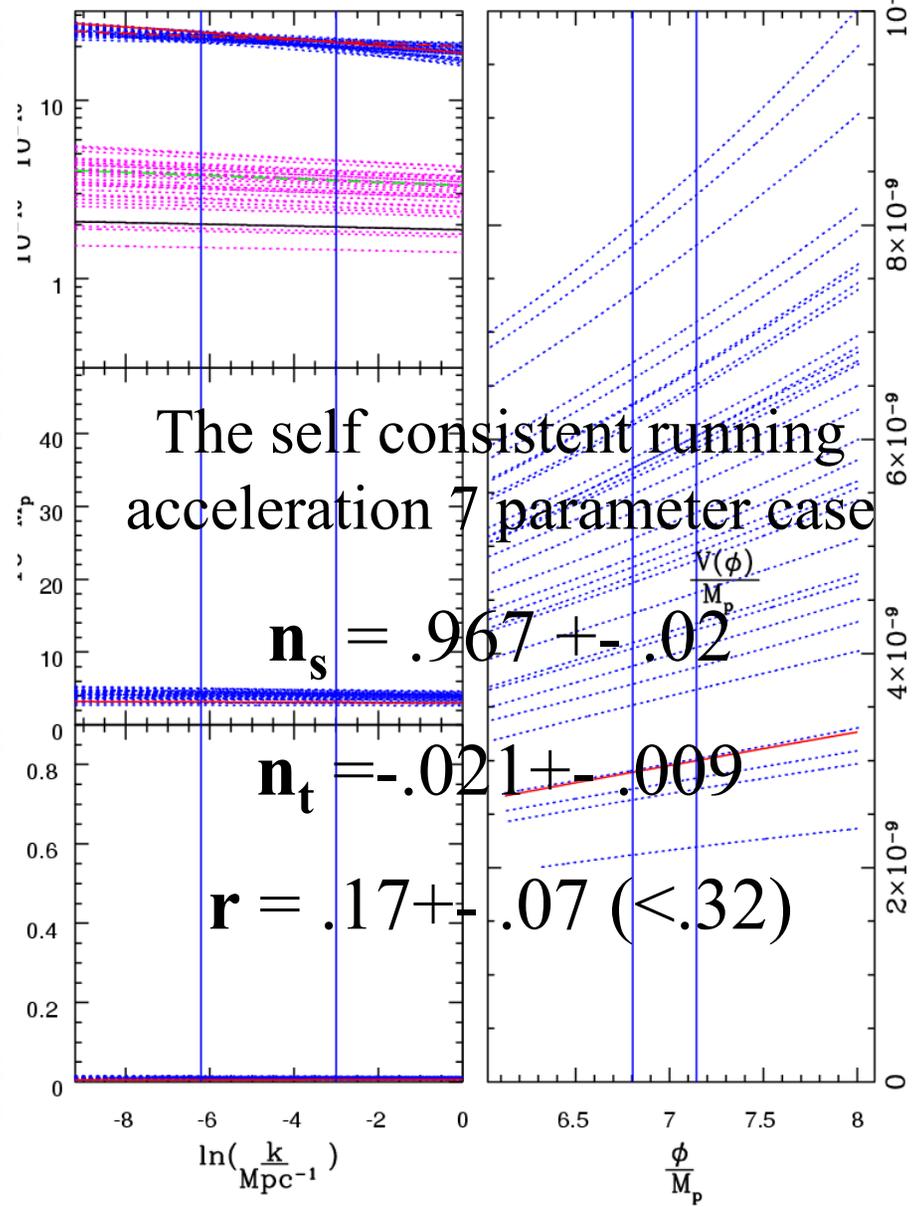


**$\epsilon$  (In Ha) order 3 + amp + 4 params cf. order 2 reconstructed from CMB+LSS data using Chebyshev nodal point expansion & MCMC**

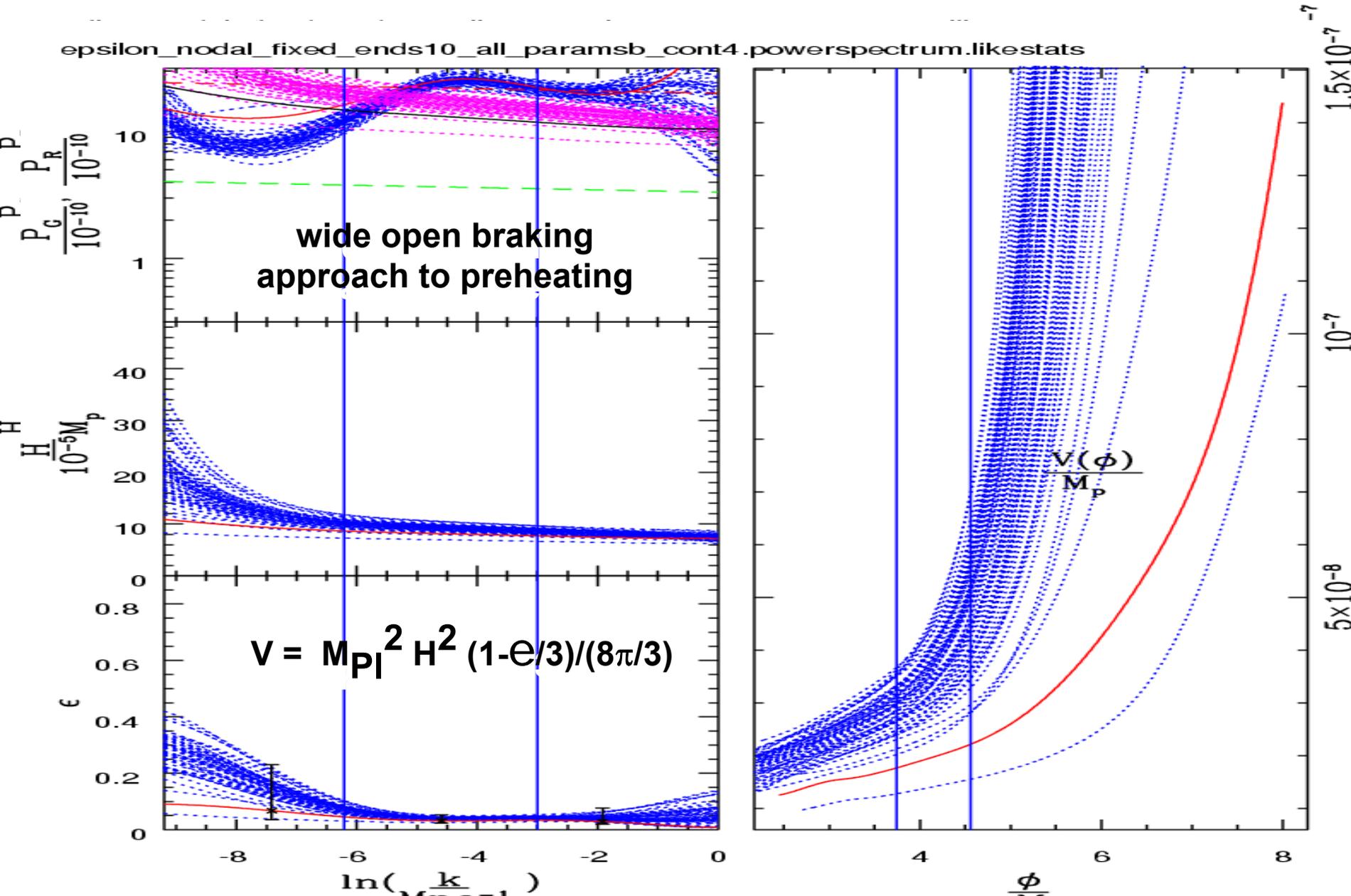
epsilon\_nodal3\_all\_paramsg\_cont5.powerspectrum.likestats



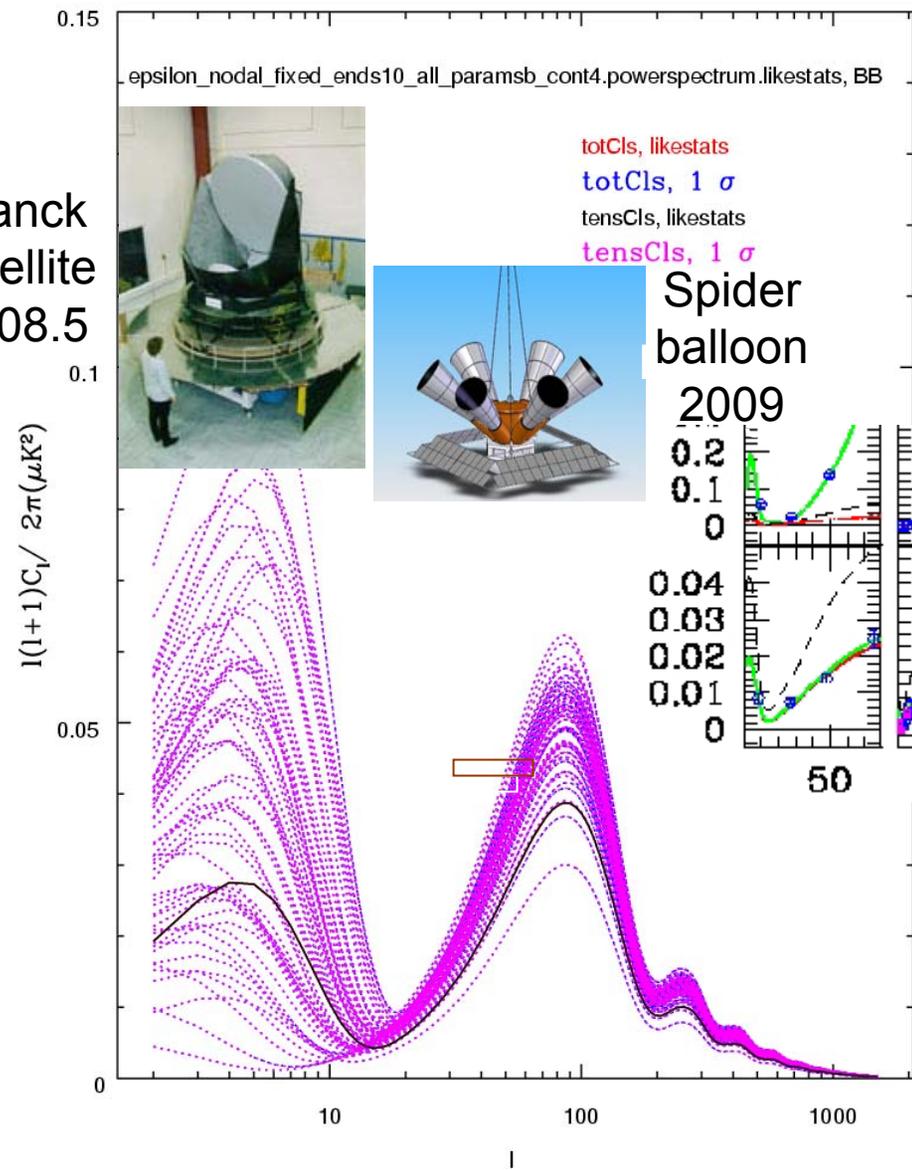
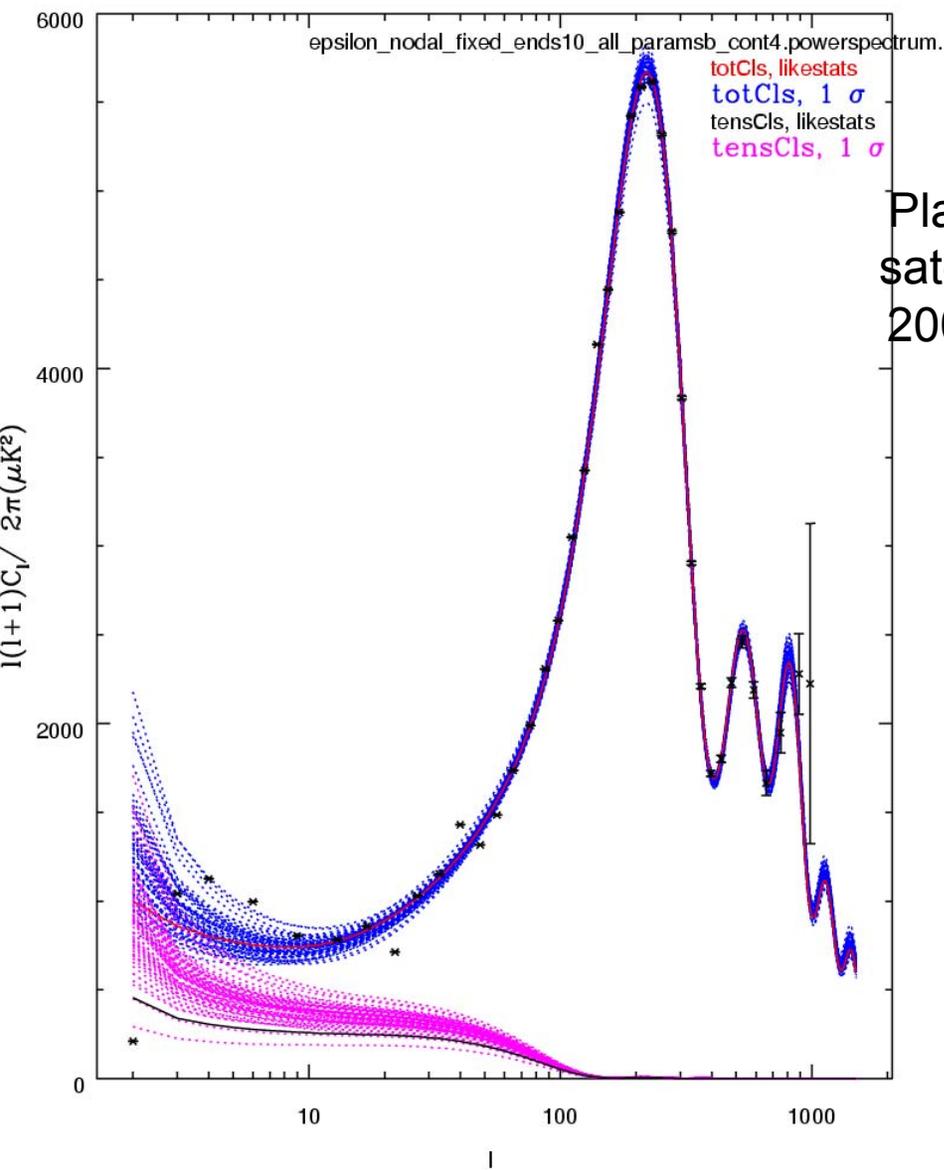
psilon\_nodal2\_cont.powerspectrum.likestats



$\epsilon$  (In Ha) order 10 + amp + 4 params reconstructed from CMB+LSS data using Chebyshev nodal point expansion & MCMC

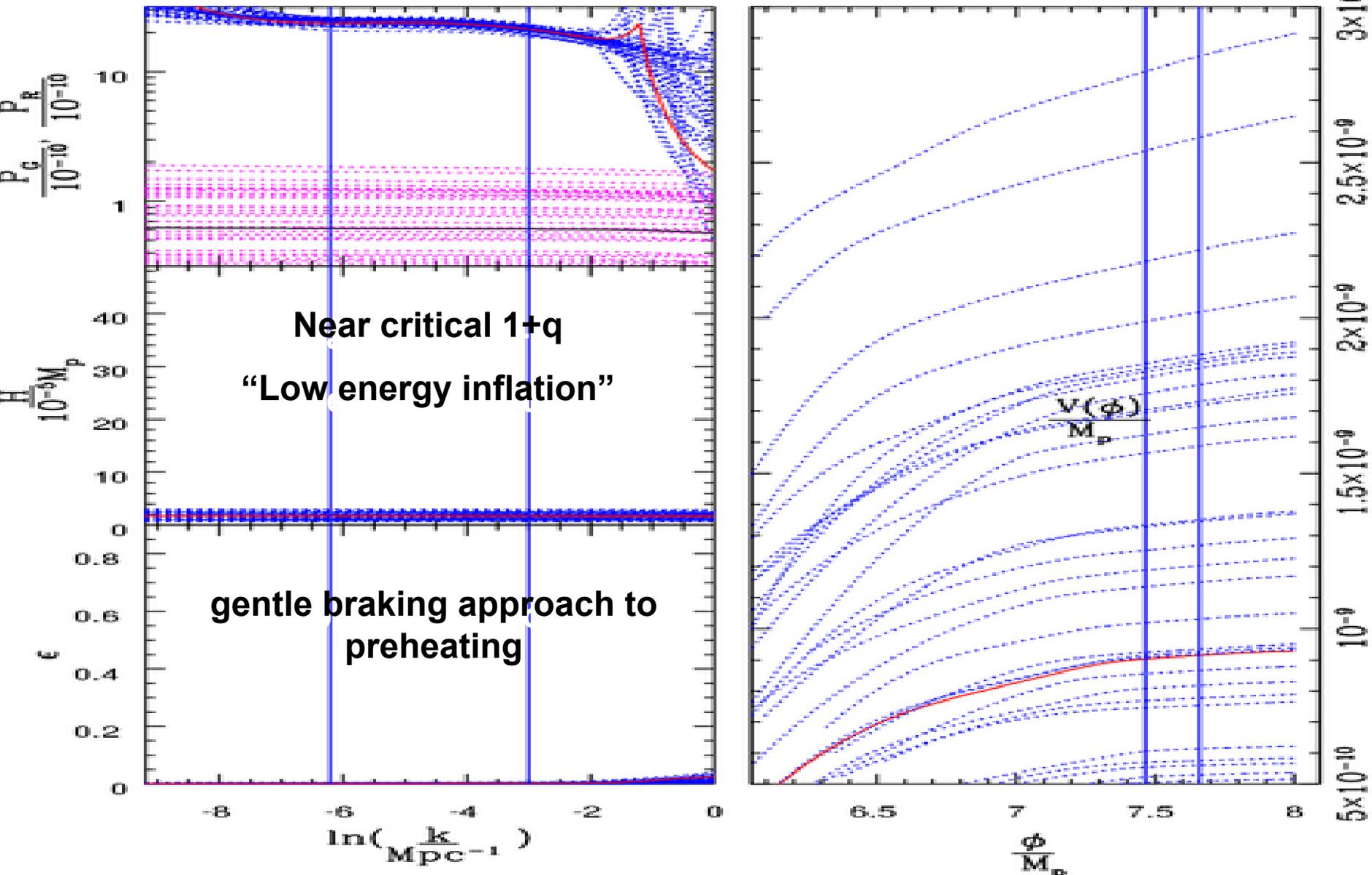


# $C_L$ TT BB for $\varepsilon$ (ln Ha) inflation trajectories reconstructed from CMB+LSS data using Chebyshev nodal point expansion (order 10) & MCMC

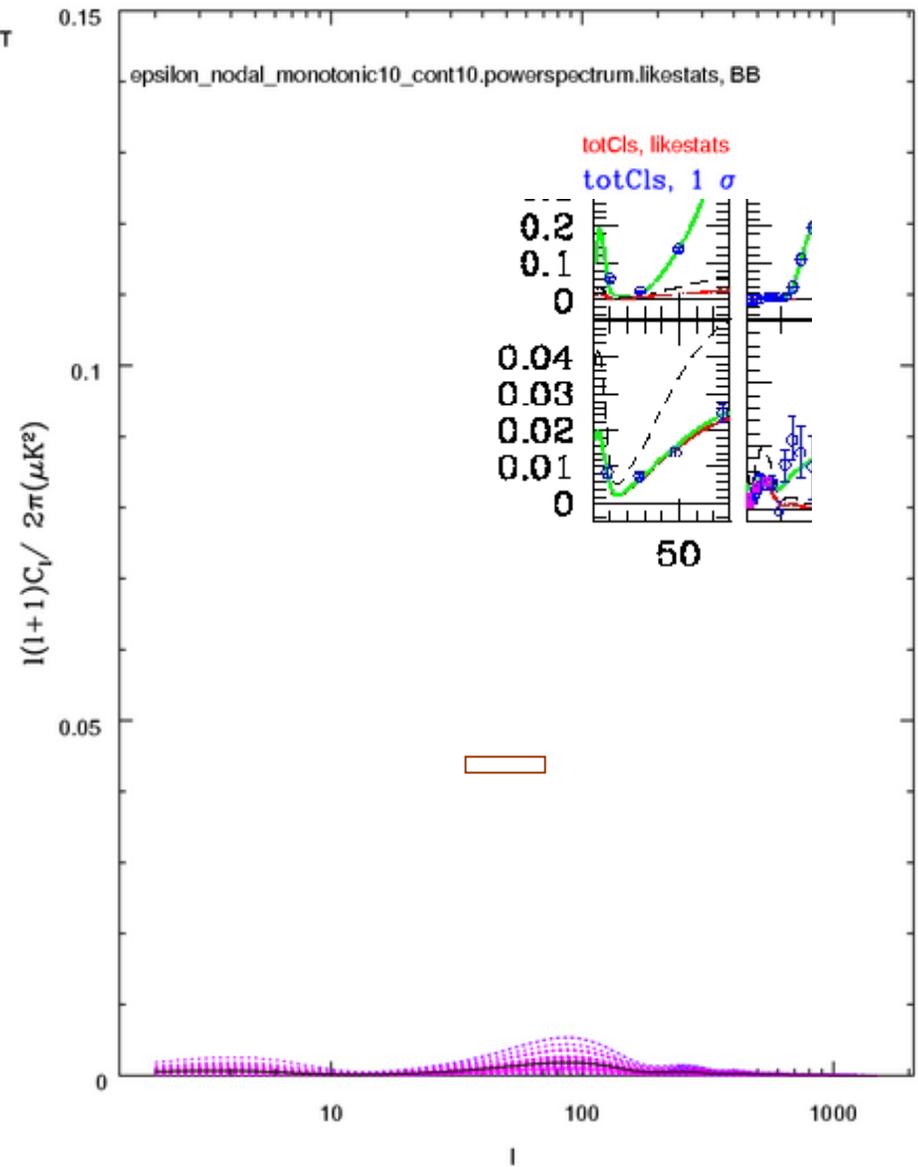
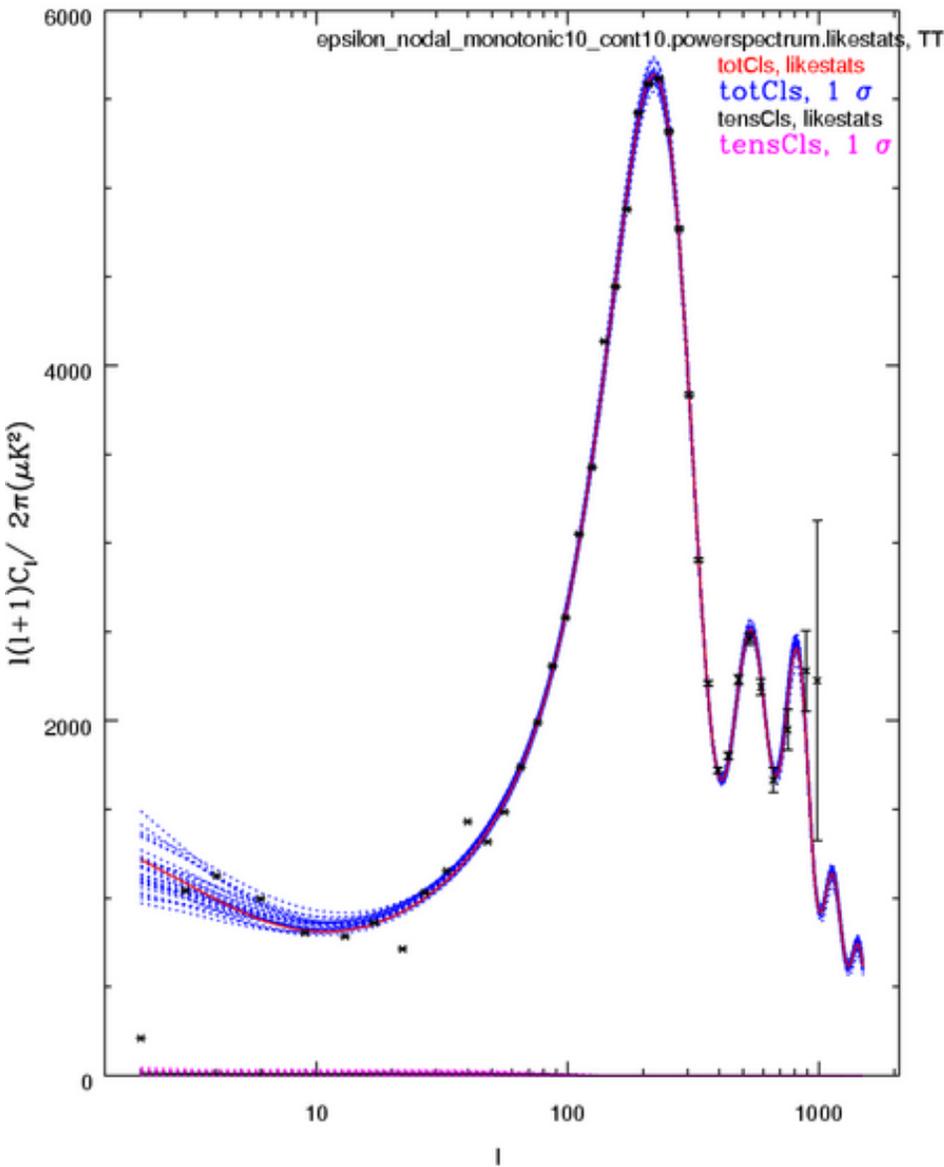


**$\epsilon$  (In Ha) order 10 monotonic + amp + 4 params reconstructed from CMB+LSS data using Chebyshev nodal point expansion & MCMC**

epsilon\_nodal\_monotonic10\_cont10.powerspectrum\_likestats



# $C_L$ TT BB for $\varepsilon$ (In Ha) monotonic inflation trajectories reconstructed from CMB+LSS data using Chebyshev nodal point expansion (order 10) & MCMC

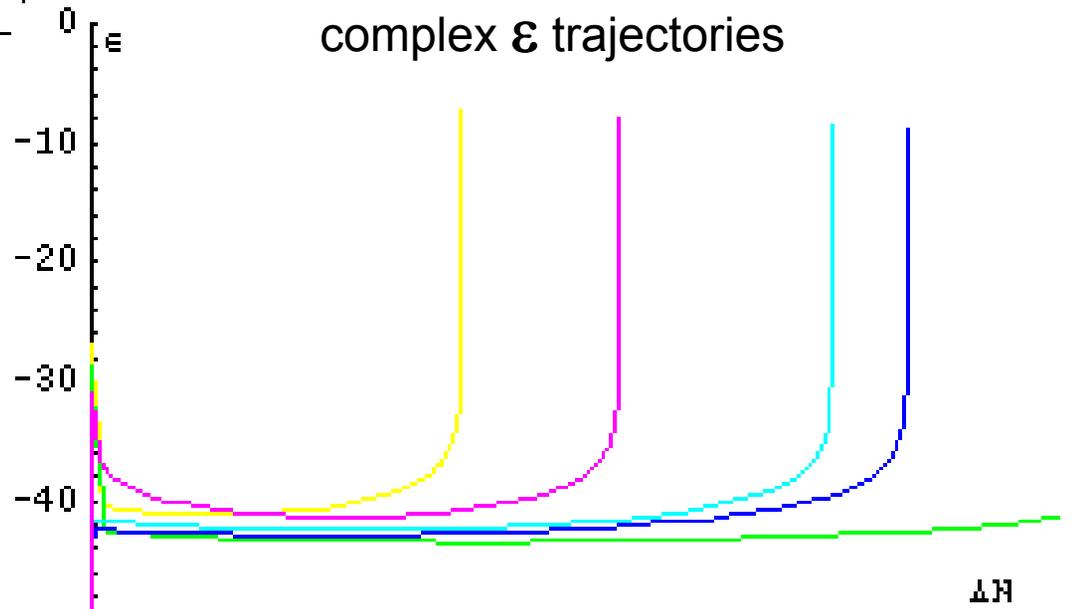
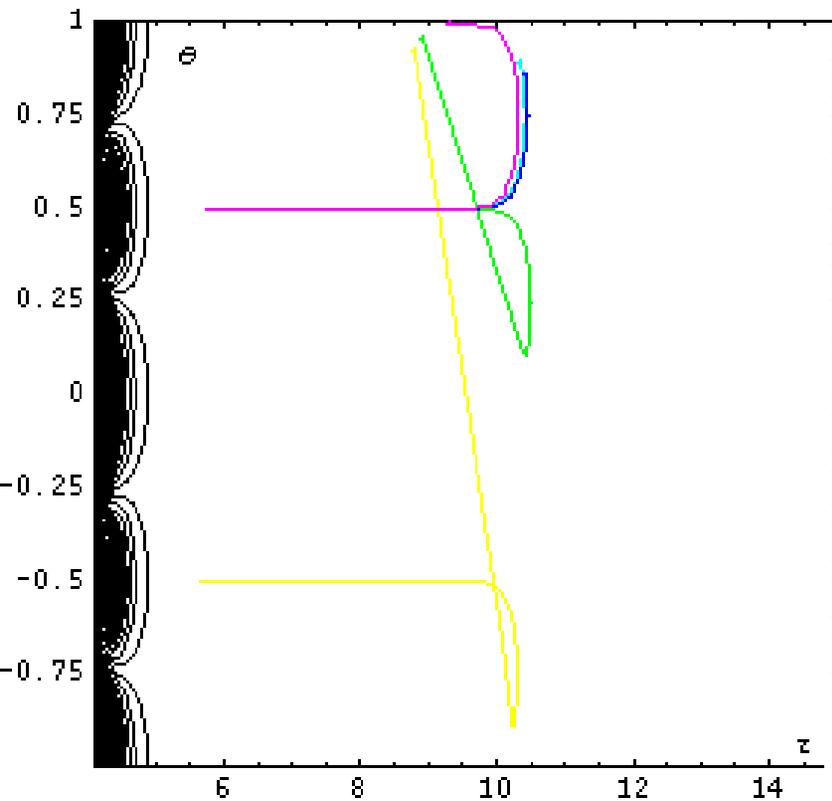


# $\varepsilon$ (ln a) trajectories in Kahler potentials

Paths that follow the downward  $\tau$ -minimum trough tend to have low  $\varepsilon$ , hence very low gravity waves (as in KKLMMT)

Some trajectories do not give enough e-foldings of inflation ( $\sim 70$  needed)

Angular direction trajectories give more complex  $\varepsilon$  trajectories



## summary

the basic 6 parameter model with no GW allowed fits all of the data OK

Usual GW limits come from adding  $r$  with a fixed GW spectrum and no consistency criterion (7 params)

Adding minimal consistency does not make that much difference (7 params)

$r$  constraints come from relating high  $k$  region of  $\sigma_8$  to low  $k$  region of GW  $C_L$

Prior probabilities on the inflation trajectories are crucial and cannot be decided at this time. Philosophy here is to be as wide open and least prejudiced about inflation as possible

**Complexity of trajectories could come out of many moduli string models.  
Example: 4-cycle complex Kahler moduli in Type IIB string theory**

Uniform priors in  $\epsilon$  nodal-point-Chebyshev-coefficients +  $H_p$  & std Cheb-coefficients give similar results: the scalar power downturns at low  $L$  if there is freedom in the mode expansion to do this. Adds GW to compensate, **breaks old  $r$  limits.**

**Monotonic uniform prior in  $\epsilon$  drives us to low energy inflation and low gravity wave content.**

***Even with low energy inflation, the prospects are good with Spider and even Planck to detect the GW-induced B-mode of polarization. Both experiments have strong Canadian roles (CSA).***

**end**